# The identification problem in SPECT: Uniqueness, non-uniqueness and stability 

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Numerical results obtained with Sonting Luo and Jianliang Qian

## The Identification Problem in SPECT

This problem arises in SPECT. We measure the radiation emitted by radioactive markers in the patient's body, modeled by a source distribution $f(x)$, attenuated by the body, with attenuation $a(x)$. We want to recover $f$ but we know neither $f$, nor a. So the question is: can we recover both? - but we care about $f$ only.

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## Math Model

The attenuated X -ray transform

$$
X_{a} f(x, \theta)=\int e^{-B a(x+t \theta, \theta)} f(x+t \theta) \mathrm{d} t, \quad x \in \mathbf{R}^{2}, \theta \in S^{1},
$$

in the plane.
We use the notation

$$
B a(x, \theta)=\int_{0}^{\infty} a(x+t \theta) \mathrm{d} t
$$

to denote the "beam transform" of $a$, usually denoted by Da.

## Identification Problem

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- There is a hidden dynamical system (in the phase space). Generally speaking, the problem is well posed, locally near some $(a, f)$, if the perturbation $\delta a$ is supported in a set which is non-trapping w.r.t. that flow.
- If $\delta$ a is supported in a trapping set; well posedness and uniqueness (even up to a finite dimensional set) may be lost. In the radial case, at least, they are lost.
- There are various degrees of instability when the non-trapping condition fails.


## If we know a, problem solved.

When a is known, it is well known that $f$ can be reconstructed uniquely, even by means of explicit formulas: Bukhgeim, Kazantsev \& Arbuzov; Novikov; Natterer.

For this reason, some of the numerical attempts to do a reconstruction are focused on recovery, or getting a good annroximation of a first instead of treating $(a f)$ as a nair. Then they get a better approximation for $a$, etc. Sometimes this is called attenuation correction. In clinical applications, additional $X$-rays are taken to reconctruct $a(x)$ first. Fliminating or reducing those additional $X$-rays remains an important problem

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## Clinical scans with the wrong and the right attenuation



Figure : SPECT cardiac scans reconstructed assuming $a=0$ (top), and with an approximation of the actual a (bottom)

## a SPECT/CT scanner



Figure: The Siemens Symbia SPECT/CT scanner

## a CT (only) scanner



Figure : The Siemens Somatom Sensation Spirit CT scanner

No much progress in the mathematical understanding of the identification problem so far. A related but not identical problem for finding both a constant attenuation and the source in the exponential X-ray transform has been solved by Solmon and Hertle. The main result in Solmon is, roughly speaking, that specific pairs of constant $a$ and radial $f$ cannot be distinguished but all other pairs can.

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The identification problem with $f$ a finite sum of delta sources has been studied by Natterer, also Boman, but the results there do not (and cannot) imply uniqueness.

Natterer also viewed the problem as a range characterization problem: if the ranges of $X_{a_{1}}$ and $X_{a_{2}}$ happen to be the same, for example, then there cannot be uniqueness. If they intersect at the origin only, there is. Range conditions, in example, in a work by Novikov, have been viewed as a possible tool for solving the problem, both numerically, for example by Bronnikov; and analytically, as in the recent work by Jolllivet \& Bal. Numerical reconstructions have been tried, too, with variable success, by
Censor et al., Manglos et al., Welch et al., Ramlau et al., Bronnikov, Zaidi, for example. Some of them use clinical data.

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A. L. BuKhgeim recently outlined a recovery algorithm if $a$ is a priori known to be a constant multiple of the characteristic function of a star-shaped domain.

## Linearization

Linearize. Notation: $\delta X_{a, f}$ acting on $(\delta a, \delta f)$. Another notation:

$$
I_{w} f(x, \theta)=\int w(x+t \theta, \theta) f(x+t \theta) \mathrm{d} t
$$

which is the weighted X -ray transform of $f$ with weight $w(x, \theta)$. Note that $X_{a}$ is of the same type but with a more special weight: $X_{a}=I_{e^{-B a}}$.

## Linearization

$$
\delta X_{a, f}(\delta a, \delta f)=I_{w} \delta a+X_{a} \delta f
$$

where

$$
w(x, \theta)=-\int_{-\infty}^{0} e^{-B a(x+t \theta, \theta)} f(x+t \theta) d t
$$

Another way to write $w$ :

$$
w=-e^{-B a} u
$$

with $u$ solving the transport equation $\left(\theta \cdot \partial_{x}+a\right)=f, u=0$ for $\theta \cdot x \ll 1$ :

$$
u(x, \theta)=\int_{-\infty}^{0} e^{-\int_{t}^{0} a(x+s \theta) \mathrm{d} s} f(x+t \theta) \mathrm{d} t .
$$

## A more general problem: inverting a sum of two weighted X-ray transforms

In other words,

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\delta X_{a, f}(\delta a, \delta f)=I_{w_{1}} \delta a+I_{w_{2}} \delta f,
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This brings us to the more general problem:

## Inverting a sum of two weighted X-ray transforms

Given two weights $w_{1,2}(x, \theta)$, and

$$
\mathcal{I}\left(g_{1}, g_{2}\right)=I_{w_{1}} g_{1}+I_{w_{2}} g_{2}
$$

find $g_{1}$ and $g_{2}$.

The first impression is that this might be too much to ask for, but the second impression is that we seem to have two equations (integrals in the $\theta$ and the $-\theta$ directions) for two unknowns. So this might work, if a certain determinant is not zero.
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To study $\operatorname{Ker} \mathcal{I}$, take the Fourier transform of

$$
\left(I_{w_{1}} g_{1}+I_{w_{2}} g_{2}\right)(z, \pm \theta)=0, \quad z \perp \theta
$$

w.r.t. $z$, to get
$\left(w_{1}\left(x, \pm D^{\perp} /|D|\right)+\right.$ l.o.t. $) g_{1}+\left(w_{2}\left(x, \pm D^{\perp} /|D|\right)+\right.$ l.o.t. $) g_{2}=0$.
Here, "l.o.t." $=$ "lower order terms". This is actually a $2 \times 2$ system of $\Psi D O$ equations.

The determinant of the principal symbol is given by the following

## Hamiltonian

$$
p_{0}(x, \xi)=W\left(x, \xi^{\perp} /|\xi|\right)
$$

where

$$
W(x, \theta)=w_{1}(x, \theta) w_{2}(x,-\theta)-w_{1}(x,-\theta) w_{2}(x, \theta) .
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This function is of fundamental importance. Since $W$ is an odd function of $\theta$, it has zeros for any $x$ ! Therefore, $p_{0}$ cannot be elliptic in any domain.

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This function is of fundamental importance. Since $W$ is an odd function of $\theta$, it has zeros for any $x$ ! Therefore, $p_{0}$ cannot be elliptic in any domain. The Hamiltonian flow of $p_{0}$ then plays a fundamental role by the Hörmander's propagation of singularities theorem. We call the projections of the Hamiltonian curves on the $x$-space rays.

## A radial example

Choose

$$
w_{1}=\frac{1}{2} \theta \cdot x, \quad w_{2}=1
$$

Then

$$
W(x, \theta)=\theta \cdot x, \quad|\xi| p_{0}=x_{1} \xi_{2}-x_{2} \xi_{1} .
$$

Therefore,

$$
|D| p_{0}(x, D)=x_{1} D_{2}-x_{2} D_{1}=-\mathrm{i} \partial / \partial \phi
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where $\phi$ is the polar angle.

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where $\phi$ is the polar angle. Bicharacteritics:

$$
x=R(\cos t, \sin t), \quad \xi=\lambda(\sin t, \cos t), \quad R \geq 0, \lambda \neq 0
$$

The rays are the circles $|x|=R \geq 0$, including the origin.

One can easily show that

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In other words, there is an infinite dimensional kernel, and the rays here appear as level curves of the functions in the kernel.

## UDOs of real principal type

We will apply the theory of $\Psi$ DOs of real principal type. If $P$ is such an operator, singularities (points of the wave front set WF $(g)$ ) of the solution $\mathrm{Pg}=0$ occupy whole bicharacteristics. Under the a priori assumption supp $g \subset K$, we can actually recover WF $(g)$ if all bicharacteristics over $K$ leave $K$ eventually.

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## Definition 1

We call $K$ non-trapping (for $p_{0}$ ) if there is no complete bicharacteristic over $K$.

The non-trapping condition plays a fundamental role in the theory of solvability of $P u=h$ in $K$. In our case, we get

If a priori, $\operatorname{supp} g \in K$, and $K$ is non-trapping, then $\mathcal{I} g=0 \Longrightarrow g \in C^{\infty}$.
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We can make it more precise:

If $\operatorname{supp} g \in K$, and $K$ is non-trapping, then
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$\mathcal{I} g \in H^{s} \Longrightarrow g \in H^{s-3 / 2}$.
The operator $\mathcal{I}$ is of order $-1 / 2$, so there is a loss of one derivative there.

## Trapping and non-trapping sets in the "radial example"



The rays are the circles $|x|=R$.

## Non-trapping, left; trapping, right

In the first case, the problem of inverting $\mathcal{I}$ is well posed
(uniqueness and stability); in the second one - it is not (ir finite dimensional kernel)

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## Back to the linearized map $\delta X_{a, f}$

What does this mean for $\delta X_{a, f}$ ?
$u(x, \theta)=u(-x, \theta)$,
where, as before, $u$ is the solution of the transport eqn. (an
attenuated integral of $f$ ). If $a=0$, then this just means that $x$ is the midpoint of the chord below. Then $(x, \theta)$ is a zero of $W$, and $\left(x, \pm \theta^{-}\right)$are characteristic. The green curve represents a ray.


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## Main results for the linearized map $\delta X_{a, f}$

Assume $\delta$ a supported in a non-trapping compact set $K$. Then

- Knowing $\delta X_{a, f}(\delta a, \delta f)$, we can recover the singularities with a loss of one derivative; there is an estimate.


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- Ker $\delta X_{a, f}$ is finite dimensional and is in $C_{0}^{\infty}(K)$.
- If $\delta X_{a, f}$ is injective on $K$, then there is a stability estimate:

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\|\delta a\|_{H^{s}(K)}+\|\delta a\|_{H^{s}(K)} \leq C\left\|\delta X_{a, f}(\delta a, \delta f)\right\|_{H^{s+3 / 2}(Z)}
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- This estimate is preserved with a uniform $C$ under a slightly stronger condition: $K \subset \Omega$ with $\Omega$ pseudo-convex: (any compact subset is non-trapping, and $\forall$ compact $K_{1} \subset \Omega$, there exists a compact set $K_{2} \subset \Omega$ so that every ray in $\Omega$ having endpoints over $K_{1}$, lies entirely in $K_{2}$ ).


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Why $\delta$ a must be supported in a non-trapping set only? The reason is that $\delta X_{a, f}$ is elliptic where $\delta a=0$ because it reduces to $X_{a}$ then.

Explicit conditions for injectivity of $\delta X_{a, f}$ (and therefore, for stability):

- Local Condition: With $W_{0}=u(x, \theta)-u(x,-\theta)$, if

$$
W_{0}\left(x_{0}, \theta\right)=0 \quad \Longrightarrow \quad \partial_{\theta^{\perp}} W_{0}\left(x_{0}, \theta\right) \neq 0
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- If $B a$ and $u$ are analytic in $K$, then $\delta X_{a, f}$ is injective on $K$.


## Main results for the non-linear map $(a, f) \mapsto X_{a} f$

## Theorem 2 (local uniqueness and stability)

Fix $a_{0}, f_{0}$. Let
$\triangleright a_{j}-a_{0}, f_{j}-f_{0}, j=1,2$ be supported in a compact set $K \subset \Omega$

- non-trapping (a bit more is needed: pseudo-convex).
- $\delta X_{a, f}$ be injective on $K$.
${ }^{\downarrow} a_{j}, f_{j}$ are close to $a_{0}, f_{0}$ in the sense
$\left\|B\left(a_{j}-a_{0}\right)\right\|_{C^{k}\left(\bar{\Omega} \times S^{1}\right)}+\left\|u_{j}-u_{0}\right\|_{C^{k}\left(\bar{\Omega} \times S^{1}\right)} \leq \varepsilon, \quad j=1,2, \quad k \gg 1$.
Then, if $\varepsilon \ll 1, X_{a_{1}} f_{1}=X_{a_{2}} f_{2}$ implies $a_{1}=a_{2}$ and $f_{1}=f_{2}$.


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Moreover, there is Hölder stability.
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Conditions for injectivity were given before (small $K$ satisfying the local injectivity condition, or analyticity).

## A radial example

We study the linearization $\delta X$ w.r.t. ( $a, f$ ) near

$$
a=0, \quad f=\mathbf{1}_{B(0,1)}
$$

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Then

$$
W_{0}=-2 \theta \cdot x
$$

The Hamiltonian $H$, up to a constant factor, is as in the previous example. Therefore, $-2|\xi| H$ is the symbol of

$$
x_{1} D_{2}-x_{2} D_{1}=-\mathrm{i} \partial / \partial \phi,
$$

$\phi$ is the polar angle in the $x$ space. The rays are the concentric circles $|x|=R, R \geq 0$, including the degenerate case $x=0$. As before, $K \subset B(0,1)$ is non-trapping, if and only if $K$ does not contain an entire circle of that kind.

Non-trapping and trapping sets:


Left: A non-trapping (and pseudo-convex) set. Small enough pertu bations supported in a non-trapping set are recoverable. Actually, $\{\operatorname{supp} \delta a \subset$ non-trapping $\}$ only is enough.

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## Non-uniqueness for general radial $a$ and $f$

Next two theorems serve as an example that for non-trapping domains, the uniqueness may fail (by more than a finite dimensional space).

One can write an explicit integral formula to compute $\delta f$, given $\delta a$.

## Non-uniqueness for general radial $a$ and $f$

Next two theorems serve as an example that for non-trapping domains, the uniqueness may fail (by more than a finite dimensional space).

## Theorem 3 (non-uniqueness for the linearization)

Let $f \in C_{0}^{\infty}$ be radial. Then $\delta X_{0, f}$ has an infinite dimensional kernel.

One can write an explicit integral formula to compute $\delta f$, given $\delta a$.

## Theorem 4 (non-uniqueness for the non-linear problem)

Let $a \in C_{0}^{\infty}$ and $f \in C_{0}^{\infty}$ be radial. Then there exists a radial $f_{0} \in C_{0}^{\infty}$ so that

$$
X_{a} f=X_{0} f_{0}
$$

Again, one can write an explicit integral formula to compute $f_{0}$, given $a$ and $f$. A simple but a non-constructive proof is to observe that the l.h.s. is an even $C_{0}^{\infty}$ function of $r=|x|$, and therefore, in the range of $X_{0}$.

## The method (developed with Luo and Jianliang) works!



Figure : "Good guess" $a_{0}=0, f_{0}=1$. Top row: original $a$ and $f$; bottom row: the reconstructed ones with good guesses

## Non-uniqueness for radial $(a, f)$



Figure : Non-uniqueness for radial $(a, f)$. Top row: $a$, bottom row: $f$. Fist column: exact $a$ and $f$; second and third columns: computed $a, f$ with different initial guesses. The reconstructions are totally wrong.

## III-posedness for perturbed radial $(a, f)$



Figure: III-posedness for perturbed radial ( $a, f$ ). Fist column: exact a and $f$; second and third columns: computed $a, f$ with different choices of initial guesses. The reconstructions are very poor.

## III posedness with circular rays



Figure: III posedness with circular rays. Top row: a, bottom row: f. Fist column: exact $a$ and $f$; second and third columns: computed $a, f$ with different initial guesses. The artifacts are circular, dictated by the Hamiltonian flow.

## Stabilized example with circular rays



Figure: A stabilized example with circular rays. A constraint condition for supp a to be in $y<0$ was used. The initial guesses were the same as in Figure 8. Top row: $a$, bottom row: $f$. Fist column: exact $a$ and $f$; second and third columns: computed $a, f$ with different initial guesses.

In the examples above, the determinant $W(x, \xi)$ has a non-degenerate characteristic variety $W=0$. The Hamiltonian flow (projected on the $x$-space) consists of concentric circles.

It is possible to have an open set in the phase space, where $W=0$. Then those zeros are stationary and we can expect high instability.

In the examples above, the determinant $W(x, \xi)$ has a non-degenerate characteristic variety $W=0$. The Hamiltonian flow (projected on the $x$-space) consists of concentric circles.

It is possible to have an open set in the phase space, where $W=0$. Then those zeros are stationary and we can expect high instability.

## Large open set of zeros of $W$



Figure : Large open set of unstable points. Top row: a, bottom row: $f$. Fist column: exact $a$ and $f ; 2 n d$ and 3 rd column: computed with guess zero for both $a$ and $f$ with the LBFGS/AG solver. With the LBFGS solver, the recovered a takes values in the range $(-0.31,0.70)$ instead of $(0,1)$, while with the AG solver, that range is $(-0.25,0.34)$. Note the black spot in the reconstructed $f$.

## A smaller open set of zeros of $W$



Figure: The phantom was moved closer to a corner compared to the previous example. This decreases the set of the unstable points (the zeros of $W$. Top row: a, bottom row: $f$. Fist column: exact $a$ and $f$; second column: computed with guess zero for both $a$ and $f$ with the LBFGS solver; computed with guess zero for both $a$ and $f$ with the AG solver.

## A less unstable case; two walls removed



Figure: We removed two "walls". This changes the properties of $W$ dramatically. The set of zeros is much smaller.

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Figure: Back to four "walls" but the have linearly changing density. The zeros of $W$ are (almost) a set of lower dimension: two opposite $\theta$ 's at each $x$.


[^0]:    This brings us to the more general problem:

