The identification problem in SPECT: Uniqueness, non-uniqueness and stability

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Numerical results obtained with Sonting Luo and Jianliang Qian

The Identification Problem in SPECT

This problem arises in SPECT. We measure the radiation emitted by radioactive markers in the patient's body, modeled by a source distribution f(x), attenuated by the body, with attenuation a(x). We want to recover f but we know neither f, nor a. So the question is: can we recover both? — but we care about f only.

Math Model

The attenuated X-ray transform

$$X_a f(x, heta) = \int e^{-Ba(x+t heta, heta)} f(x+t heta) \,\mathrm{d}t, \quad x \in \mathbf{R}^2, \; heta \in S^1,$$

in the plane.

We use the notation

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Given $X_a f$, recover both a and f.

- Short answer: Depends.
- Sometimes we can, even in a stable way; sometimes we cannot.
- There is a hidden dynamical system (in the phase space). Generally speaking, the problem is well posed, locally near some (a, f), if the perturbation δa is supported in a set which is *non-trapping* w.r.t. that flow.
- If δa is supported in a trapping set ; well posedness and uniqueness (even up to a finite dimensional set) may be lost.
 In the radial case, at least, they are lost.
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When a is known, it is well known that f can be reconstructed uniquely, even by means of explicit formulas: BUKHGEIM, KAZANTSEV & ARBUZOV; NOVIKOV; NATTERER.

For this reason, some of the numerical attempts to do a reconstruction are focused on recovery, or getting a good approximation of *a* first, instead of treating (a, f) as a pair. Then they get a better approximation for *a*, etc. Sometimes this is called *attenuation correction*. In clinical applications, additional X-rays are taken to reconstruct a(x) first. Eliminating or reducing those additional X-rays remains an important problem.

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Clinical scans with the wrong and the right attenuation



Figure : SPECT cardiac scans reconstructed assuming a = 0 (top), and with an approximation of the actual a (bottom)

a SPECT/CT scanner



Figure : The Siemens Symbia SPECT/CT scanner

a CT (only) scanner



Figure : The Siemens Somatom Sensation Spirit CT scanner

No much progress in the mathematical understanding of the identification problem so far. A related but not identical problem for finding both a constant attenuation and the source in the exponential X-ray transform has been solved by SOLMON and HERTLE. The main result in SOLMON is, roughly speaking, that specific pairs of constant *a* and radial *f* cannot be distinguished but all other pairs can.

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NATTERER also viewed the problem as a range characterization problem: if the ranges of X_{a_1} and X_{a_2} happen to be the same, for example, then there cannot be uniqueness. If they intersect at the origin only, there is. Range conditions, in example, in a work by NOVIKOV, have been viewed as a possible tool for solving the problem, both numerically, for example by BRONNIKOV; and analytically, as in the recent work by JOLLLIVET & BAL. Numerical reconstructions have been tried, too, with variable success, by

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Linearization

Linearize. Notation: $\delta X_{a,f}$ acting on $(\delta a, \delta f)$. Another notation:

$$I_w f(x,\theta) = \int w(x+t\theta,\theta)f(x+t\theta)\,\mathrm{d}t,$$

which is the weighted X-ray transform of f with weight $w(x, \theta)$. Note that X_a is of the same type but with a more special weight: $X_a = I_{e^{-Ba}}$.

Linearization

$$\delta X_{a,f}(\delta a, \delta f) = I_w \delta a + X_a \delta f,$$

where

$$w(x,\theta) = -\int_{-\infty}^{0} e^{-Ba(x+t\theta,\theta)} f(x+t\theta) dt.$$

Another way to write w:

$$w=-e^{-Ba}u,$$

with *u* solving the transport equation $(\theta \cdot \partial_x + a) = f$, u = 0 for $\theta \cdot x \ll 1$:

$$u(x,\theta) = \int_{-\infty}^{0} e^{-\int_{t}^{0} a(x+s\theta) \,\mathrm{d}s} f(x+t\theta) \,\mathrm{d}t.$$

A more general problem: inverting a sum of two weighted X-ray transforms

In other words,

$$\delta X_{a,f}(\delta a, \delta f) = I_{w_1} \delta a + I_{w_2} \delta f,$$

where

$$w_1 = -e^{-Ba}u, \quad w_2 = e^{-Ba}.$$

This brings us to the more general problem:

Inverting a sum of two weighted X-ray transforms

Given two weights $w_{1,2}(x,\theta)$, and

 $I(g_1, g_2) = I_{w_1}g_1 + I_{w_2}g_2,$

find g_1 and g_2 .

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Inverting a sum of two weighted X-ray transforms

Given two weights $w_{1,2}(x,\theta)$, and

$$\mathcal{I}(g_1,g_2) = I_{w_1}g_1 + I_{w_2}g_2,$$

find g_1 and g_2 .

The first impression is that this might be too much to ask for, but the second impression is that we seem to have two equations (integrals in the θ and the $-\theta$ directions) for two unknowns. So this might work, if a certain determinant is not zero. On the other hand, that determinant vanishes for any x (and some θ), as we will see in a moment.

To study Ker \mathcal{I} , take the Fourier transform of

$$(I_{w_1}g_1+I_{w_2}g_2)(z,\pm\theta)=0, \quad z\perp\theta,$$

w.r.t. z, to get

$$\left(w_1(x,\pm D^{\perp}/|D|) + \text{l.o.t.}\right)g_1 + \left(w_2(x,\pm D^{\perp}/|D|) + \text{l.o.t.}\right)g_2 = 0.$$

Here, "l.o.t." = "lower order terms". This is actually a 2×2 system of ΨDO equations.

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The determinant of the principal symbol is given by the following

Hamiltonian

$$p_0(x,\xi) = W(x,\xi^{\perp}/|\xi|),$$

where

$$W(x,\theta) = w_1(x,\theta)w_2(x,-\theta) - w_1(x,-\theta)w_2(x,\theta).$$

This function is of fundamental importance. Since W is an odd function of θ , it has zeros for any x! Therefore, p_0 cannot be elliptic in any domain. The Hamiltonian flow of p_0 then plays a fundamental role by the Hörmander's propagation of singularities theorem. We call the projections of the Hamiltonian curves on the *x*-space **rays**.

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A radial example

Choose

$$w_1 = \frac{1}{2}\theta \cdot x, \quad w_2 = 1.$$

Then

$$W(x,\theta) = \theta \cdot x, \quad |\xi|p_0 = x_1\xi_2 - x_2\xi_1.$$

Therefore,

$$|D|p_0(x,D) = x_1D_2 - x_2D_1 = -\mathrm{i}\partial/\partial\phi,$$

where ϕ is the polar angle. Bicharacteritics:

 $x = R(\cos t, \sin t), \quad \xi = \lambda(\sin t, \cos t), \quad R \ge 0, \lambda \ne 0.$

The rays are the circles $|x| = R \ge 0$, including the origin.

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We will apply the theory of Ψ DOs of real principal type. If P is such an operator, singularities (points of the wave front set WF(g)) of the solution Pg = 0 occupy whole bicharacteristics. Under the a priori assumption $\operatorname{supp} g \subset K$, we can actually recover WF(g) if all bicharacteristics over K leave K eventually.

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We call *K* non-trapping (for *p*₀) if there is no complete bicharacteristic over *K*.

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We call K non-trapping (for p_0) if there is no complete bicharacteristic over K.

The non-trapping condition plays a fundamental role in the theory of solvability of Pu = h in K. In our case, we get

If a priori, supp $g \in K$, and K is non-trapping, then $\mathcal{I}g = 0 \Longrightarrow g \in C^{\infty}$.

We can make it more precise:

If supp $g \in K$, and K is non-trapping, then $\mathcal{I}g \in H^s \Longrightarrow g \in H^{s-3/2}$.

The operator ${\mathcal I}$ is of order -1/2, so there is a loss of one derivative there.

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Trapping and non-trapping sets in the "radial example"





The rays are the circles |x| = R.

Non-trapping, left; trapping, right.

In the first case, the problem of inverting \mathcal{I} is well posed (uniqueness and stability); in the second one — it is not (infinite dimensional kernel).

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Back to the linearized map $\delta X_{a,f}$

What does this mean for $\delta X_{a,f}$? Look for (x, θ) so that

 $u(x,\theta)=u(-x,\theta),$

where, as before, u is the solution of the transport eqn. (an attenuated integral of f). If a = 0, then this just means that x is the midpoint of the chord below. Then (x, θ) is a zero of W, and $(x, \pm \theta^{\perp})$ are characteristic. The green curve represents a ray.



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Assume δa supported in a non-trapping compact set K. Then

- ► Knowing $\delta X_{a,f}(\delta a, \delta f)$, we can recover the singularities with a loss of one derivative; there is an estimate.
- Ker $\delta X_{a,f}$ is finite dimensional and is in $C_0^{\infty}(K)$.
- If $\delta X_{a,f}$ is injective on K, then there is a stability estimate:

 $\|\delta a\|_{H^{s}(K)} + \|\delta a\|_{H^{s}(K)} \leq C \|\delta X_{a,f}(\delta a, \delta f)\|_{H^{s+3/2}(Z)}.$

For this estimate is preserved with a uniform C under a slightly stronger condition: $K \subset \Omega$ with Ω pseudo-convex: (any compact subset is non-trapping, and \forall compact $K_1 \subset \Omega$, there exists a compact set $K_2 \subset \Omega$ so that every ray in Ω having endpoints over K_1 , lies entirely in K_2).

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Explicit conditions for injectivity of $\delta X_{a,f}$ (and therefore, for stability):

▶ Local Condition: With $W_0 = u(x, \theta) - u(x, -\theta)$, if

$$W_0(x_0, \theta) = 0 \implies \partial_{\theta^{\perp}} W_0(x_0, \theta) \neq 0,$$

then $\delta X_{a,f}$ is injective on some neighborhood of x_0 . This property guarantees that the rays have non-zero speed at x_0 , i.e., x_0 is non-trapping.

If Ba and u are analytic in K, then δX_{a,f} is injective on K.

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▶ If *Ba* and *u* are analytic in *K*, then $\delta X_{a,f}$ is injective on *K*.

Theorem 2 (local uniqueness and stability)

Fix a₀, f₀. Let

- ► $a_j a_0$, $f_j f_0$, j = 1, 2 be supported in a compact set $K \subset \Omega$ — non-trapping (a bit more is needed: pseudo-convex).
- $\delta X_{a,f}$ be injective on K.
- ▶ *a_j*, *f_j* are close to *a*₀, *f*₀ in the sense

$$\|B(a_j-a_0)\|_{C^k(\bar{\Omega}\times S^1)}+\|u_j-u_0\|_{C^k(\bar{\Omega}\times S^1)}\leq \varepsilon, \quad j=1,2, \quad k\gg 1.$$

Then, if $\varepsilon \ll 1$, $X_{a_1}f_1 = X_{a_2}f_2$ implies $a_1 = a_2$ and $f_1 = f_2$. Moreover, there is Hölder stability.

The estimate above is just a smallness condition for $a_j - a_0$ and $f_j - f_0$ but a weaker one.

Theorem 2 (local uniqueness and stability)

Fix a_0 , f_0 . Let

- ► $a_j a_0$, $f_j f_0$, j = 1, 2 be supported in a compact set $K \subset \Omega$ — non-trapping (a bit more is needed: pseudo-convex).
- $\delta X_{a,f}$ be injective on K.
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A radial example

We study the linearization δX w.r.t. (a, f) near

$$a=0, \quad f=\mathbf{1}_{B(0,1)},$$

for perturbations of those a and f supported in B(0,1) only. Recall

$$\delta X_{a,f}(\delta a, \delta f) = I_w \delta a + I_1 \delta f$$

with

$$w(x,\theta) = -\sqrt{1-(\theta^{\perp}\cdot x)^2} - \theta\cdot x$$

Then

 $W_0 = -2\theta \cdot x.$

The Hamiltonian H, up to a constant factor, is as in the previous example. Therefore, $-2|\xi|H$ is the symbol of

$$x_1D_2 - x_2D_1 = -\mathrm{i}\partial/\partial\phi,$$

 ϕ is the polar angle in the x space. The rays are the concentric circles |x| = R, $R \ge 0$, including the degenerate case x = 0. As before, $K \subset B(0,1)$ is non-trapping, if and only if K does not contain an entire circle of that kind.

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Figure : Left: A non-trapping (and pseudo-convex) set. Small enough perturbations supported in a non-trapping set are recoverable. Actually, $\{\operatorname{supp} \delta a \subset \operatorname{non-trapping}\}$ only is enough.

In the whole ball (which is trapping) the kernel of $\delta X(\delta a, \delta f)$ consists of radial δa and δf supported there connected by (using the Radon transform notation $Rf(p, \omega)$)

$$\sqrt{1-p^2}R\delta a - R\delta f = 0,$$



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Next two theorems serve as an example that for non-trapping domains, the uniqueness may fail (by more than a finite dimensional space).

Theorem 3 (non-uniqueness for the linearization)

Let $f \in C_0^{\infty}$ be radial. Then $\delta X_{0,f}$ has an infinite dimensional kernel.

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One can write an explicit integral formula to compute δf , given δa .

Theorem 4 (non-uniqueness for the non-linear problem)

Let $a \in C_0^\infty$ and $f \in C_0^\infty$ be radial. Then there exists a radial $f_0 \in C_0^\infty$ so that

 $X_a f = X_0 f_0.$

Again, one can write an explicit integral formula to compute f_0 , given *a* and *f*. A simple but a non-constructive proof is to observe that the l.h.s. is an even C_0^{∞} function of r = |x|, and therefore, in the range of X_0 .

The method (developed with Luo and Jianliang) works!



Figure : "Good guess" $a_0 = 0$, $f_0 = 1$. Top row: original *a* and *f*; bottom row: the reconstructed ones with good guesses

Non-uniqueness for radial (a, f)



Figure : Non-uniqueness for radial (a, f). Top row: a, bottom row: f. Fist column: exact a and f; second and third columns: computed a, f with different initial guesses. The reconstructions are totally wrong.

Ill-posedness for perturbed radial (a, f)



Figure : Ill-posedness for perturbed radial (a, f). Fist column: exact a and f; second and third columns: computed a, f with different choices of initial guesses. The reconstructions are very poor.

III posedness with circular rays



Figure : Ill posedness with circular rays. Top row: a, bottom row: f. Fist column: exact a and f; second and third columns: computed a, f with different initial guesses. The artifacts are circular, dictated by the Hamiltonian flow.

Stabilized example with circular rays



Figure : A stabilized example with circular rays. A constraint condition for supp *a* to be in y < 0 was used. The initial guesses were the same as in Figure 8. Top row: *a*, bottom row: *f*. Fist column: exact *a* and *f*; second and third columns: computed *a*, *f* with different initial guesses.

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In the examples above, the determinant $W(x,\xi)$ has a non-degenerate characteristic variety W = 0. The Hamiltonian flow (projected on the x-space) consists of concentric circles.

It is possible to have an open set in the phase space, where W = 0. Then those zeros are stationary and we can expect high instability. In the examples above, the determinant $W(x,\xi)$ has a non-degenerate characteristic variety W = 0. The Hamiltonian flow (projected on the x-space) consists of concentric circles.

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Large open set of zeros of W



Figure : Large open set of unstable points. Top row: *a*, bottom row: *f*. Fist column: exact *a* and *f*; 2nd and 3rd column: computed with guess zero for both *a* and *f* with the LBFGS/AG solver. With the LBFGS solver, the recovered *a* takes values in the range (-0.31, 0.70) instead of (0, 1), while with the AG solver, that range is (-0.25, 0.34). Note the black spot in the reconstructed *f*.

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A smaller open set of zeros of W



Figure : The phantom was moved closer to a corner compared to the previous example. This decreases the set of the unstable points (the zeros of W. Top row: *a*, bottom row: *f*. Fist column: exact *a* and *f*; second column: computed with guess zero for both *a* and *f* with the LBFGS solver; computed with guess zero for both *a* and *f* with the AG solver.



Figure : We removed two "walls". This changes the properties of W dramatically. The set of zeros is much smaller.



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Another less unstable case



Figure : Back to four "walls" but the have linearly changing density. The zeros of W are (almost) a set of lower dimension: two opposite θ 's at each x.

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