

## LECTURE 15

### 15. GLOBAL SOLUTIONS

In this lecture we study the so-called global solutions, i.e. solutions defined in the whole space, with an additional assumption that they grow quadratically at infinity. More precisely, we consider elements of the class  $P_\infty(x_0, M)$  which satisfy

- $\|D^2u\|_{L^\infty(\mathbb{R}^n)} \leq M$ ,
- $x_0 \in \Gamma(u)$ .

The global solutions may exist by their own, but most importantly they may appear as blowups of one or a sequence of functions with variable centers, i.e. limits of rescalings

$$u_{x_j, r_j}^j(x) = \frac{u^j(x_j + r_j x) - u^j(x_j)}{r_j^2}.$$

We will first study the global solution for the classical obstacle problem, then generalize the results for Problems A, B and at the end of this lecture we will study the case of Problem C.

#### 15.1. Classical Obstacle Problem.

**Theorem 15.1.** *Let  $u \in P_\infty(M)$  be a global solution of Problem A and assume that  $u \geq 0$  in  $\mathbb{R}^n$ . Then  $u$  is a convex function in  $\mathbb{R}^n$ , i.e.*

$$\partial_{ee}u(x) \geq 0, \quad \text{for any direction } e \text{ and } x \in \mathbb{R}^n$$

*In particular, the set  $\{u = 0\}$  is convex.*

*Proof.* Fix any direction  $e$ . Without loss of generality suppose that  $e = e_n = (0, \dots, 0, 1)$ . Assume, on the contrary, that

$$-m := \inf_{\Omega} \partial_{nn}u < 0,$$

and let  $x_j \in \Omega$  be a minimizing sequence for the value  $-m$ , i.e.

$$\lim_{j \rightarrow \infty} \partial_{nn}u(x_j) = -m.$$

Let  $d_j = \text{dist}(x_j, \Gamma)$  and consider the rescalings

$$u_j(x) = u_{x_j, d_j}(x) = \frac{1}{d_j^2} u(x_j + d_j x).$$

Observe that  $B_1 \subset \Omega(u_j)$  and the free boundary  $\Gamma(u_j)$  contains at least one point on  $\partial B_1$ . Since also  $\|D^2u_j\|$  are uniformly bounded we have the uniform estimates

$$|u_j(x)| \leq \frac{M}{2} (R+1)^2$$

for all  $R > 0$  and therefore we can extract a subsequence converging in  $C_{\text{loc}}^{1,\alpha}(\mathbb{R}^n)$  to a global solution  $u_0$  of Problem A. The assumption  $u \geq 0$ , implies that  $u_0 \geq 0$  and therefore,  $\Omega(u_0) = \{u_0 > 0\}$ . Moreover, similarly to  $u_j$ , observe that since  $B_1 \subset \Omega(u_0)$ , and  $\partial B_1$  contains at least one free boundary point.

Next observe, since all functions  $u_j$  satisfy  $\Delta u_j = 1$  in  $B_1$ , the convergence to  $u_0$  can be assumed to be at least in  $C_{\text{loc}}^2(B_1)$ . Hence, the limit function  $u_0$  satisfies

$$\Delta u_0 = 1, \quad \partial_{nn} u_0 \geq -m \quad \text{in } B_1, \quad \partial_{nn} u_0(0) = -m.$$

Since  $\partial_{nn} u_0$  is harmonic in  $B_1$ , the minimum principle implies that  $\partial_{nn} u_0 \equiv -m$  in  $B_1$ . In fact we have even more,  $\partial_{nn} u_0 = -m$  in the connected component of  $\Omega(u_0)$  which contains  $B_1$ . Hence we obtain the representations

$$(15.1) \quad \partial_n u_0(x) = g_1(x') - mx_n, \quad x' = (x_1, \dots, x_{n-1})$$

and

$$(15.2) \quad u_0(x) = g_2(x') + g_1(x')x_n - \frac{m}{2}x_n^2,$$

in  $B_1$ . Now let us choose a point  $(x', 0) \in B_1$  and start moving in the direction  $e_n$ . Observe that as long as we stay in  $\Omega(u_0)$ , we still have  $\partial_{nn} u = -m$  and therefore still have the representations (15.1)–(15.2). However, sooner or later we will reach  $\partial\Omega(u_0)$ , otherwise if  $x_n$  becomes very large (15.2) will imply  $u_0 < 0$ , contrary to our assumption. Since  $u_0 = |\nabla u_0| = 0$  on  $\partial\Omega(u_0)$ , from (15.1) we obtain that the first value  $\xi(x')$  of  $x_n$  for which we arrive at  $\partial\Omega(u_0)$  is given by

$$\xi(x') = \frac{g_1(x')}{m}.$$

Hence from (15.2) we deduce that

$$g_2(x') = -\frac{g_1(x')^2}{2m}.$$

Now, the representation (15.2) takes the form

$$u_0(x) = -\frac{m}{2}(x_n - \xi(x'))^2,$$

which is not possible since  $u_0 \geq 0$ . This concludes the proof.  $\square$

**15.2. Problems A, B.** Next, our goal is to generalize Theorem 15.1 for global solutions of Problems A, B. We will consider two cases: when the complement of  $\Omega$  is bounded and when it is unbounded.

**15.2.1. The compact complement case.** Assume now we have  $u \in P_\infty(x_0, M)$  for which  $\Omega^c$  is compact.

**Lemma 15.2.** *Let  $u \in P_\infty(x_0, M)$  be a global solution of Problem A, B, such that  $\Omega^c$  is compact and  $\text{Int } \Omega^c \neq \emptyset$ . Then  $x_0$  is a low energy point.*

*Proof.* Suppose, towards a contradiction, that  $x_0$  is a high energy point. Consider then a so-called “shrink-down” of  $u$  with a fixed center at  $x_0$ , i.e. sequence of rescalings

$$u_k(x) = u_{x_0, R_k}(x) = \frac{u(x_0 + R_k x) - u(x_0)}{R_k^2}$$

for  $R_k \rightarrow \infty$  which converges to a global solution  $u_\infty$ . Similarly to blowups with fixed centers (Theorem 10.3), it is not hard to show that  $u_\infty$  is a homogeneous global solution, as a simple corollary of Weiss’s monotonicity formula (see Lecture 10). The same monotonicity formula implies

$$(15.3) \quad \begin{aligned} \alpha_n = \omega(x_0) &\leq W(r, u, x_0) \leq \lim_{R_k \rightarrow \infty} W(R_k, u, x_0) \\ &= \lim_{R_k \rightarrow \infty} W(1, u_k) = W(1, u_\infty). \end{aligned}$$

On the other hand, we know that for homogeneous global solutions  $W$  can take only two values:  $\alpha_n/2$  and  $\alpha_n$ , hence  $W(1, u_\infty) = \alpha_n$ . global solution and  $W(1, u_\infty) \leq \alpha_n$ . This, combined with Hence, from (15.3) implies that  $W(r, u, x_0) = \alpha_n$  for any  $r > 0$ . Thus, by Theorem 10.2  $u$  must be homogeneous with respect to the point  $x_0$  and the classification of homogeneous solutions implies that  $u$  must be a polynomial solution. This contradicts the assumption  $\text{Int } \Omega^c \neq \emptyset$ .  $\square$

**Lemma 15.3.** *Let  $u$  be as in Lemma 15.2 . Then  $\Omega^c$  will consist of finite union of connected components  $\Omega_i^c$ ,  $i = 1, \dots, N$  with  $C^1$  boundaries and nonempty interiors such that  $u$  is constant in  $\Omega_i^c$ .*

*Proof.* Note that every point on  $\partial\Omega$  is of low energy. Applying now Theorems 12.4 and 13.1 we obtain the desired structure for  $\Omega^c$ .  $\square$

**Lemma 15.4.** *Let  $u$  be as in Lemma 15.2 and suppose that*

$$(15.4) \quad \sup_{\Omega^c} u = 0$$

*Then, for a suitable choice of the origin in  $\Omega^c$ , the function*

$$r \mapsto \frac{u(rx)}{r^2}$$

*is nondecreasing, for any fixed  $x$ .*

*Proof.* We will give the proof for  $n \geq 3$ . Denote by  $V$  the Newtonian potential of  $\Omega^c$ , i.e.

$$V(x) = \int_{\Omega^c} \frac{c_n}{|x-y|^{n-2}} dy.$$

Then  $V$  is bounded and superharmonic in  $\mathbb{R}^n$  and harmonic in  $\Omega$ . By the maximum principle, there is at least one point  $\zeta_0 \in \Omega^c$  such that

$$V(\zeta_0) \geq V(x) \quad \text{for all } x \in \mathbb{R}^n.$$

Set the origin at  $\zeta_0$ .

Since

$$\Delta(u - V) = 1$$

in the sense of distributions and all second order partial derivatives of  $u - V$  are bounded harmonic functions, the Hessian of  $u - V$  is a constant matrix, by Liouville's theorem. Hence  $u - V$  is a polynomial of degree two. Set

$$P(x) = u(x) - V(x) - u(0) + V(0).$$

Note that  $|\nabla V(0)| = |\nabla u(0)| = 0$ . Hence  $P(0) = |\nabla P(0)| = 0$ , this implies that  $P$  is homogeneous. Now consider the function

$$h(x) = x \cdot \nabla u(x) - 2u(x).$$

$h$  is continuous in  $\mathbb{R}^n$  and for all  $x \neq 0$  fixed,

$$\frac{d}{dr} \left( \frac{u(rx)}{r^2} \right) = \frac{1}{r^3} h(rx).$$

We will show that  $h$  is non-negative in  $\mathbb{R}^n$ . In fact

$$h(x) = -2u(x) \geq 0, \quad \forall x \in \Omega^c.$$

On the other hand, by the homogeneity of  $P$ ,

$$h(x) = x \cdot \nabla V(x) - 2V(x) + 2V(0) - 2u(0)$$

then

$$\lim_{|x| \rightarrow \infty} h(x) = 2V(0) - 2u(0) \geq 0.$$

Since  $h$  harmonic in  $\Omega$ , by the minimum principle,  $h$  is positive in  $\Omega$ .  $\square$

Now we can prove the main result of this section.

**Theorem 15.5.** *Let  $u \in P_\infty(M)$  be such that  $\Omega^c$  is compact and  $\text{Int } \Omega^c \neq \emptyset$ . Suppose also that (15.4) holds. Then  $u \geq 0$  in  $\mathbb{R}^n$  and  $u$  is a convex function. In particular  $\{u = 0\}$  is a convex set.*

*Proof.* Choose the origin as in Lemma 15.4.

Consider first the case of Problem A. We claim that there exists small  $\rho > 0$  such that  $u \geq 0$  in  $B_\rho$ . Indeed, if  $0 \in \text{Int } \Omega^c$ , this is immediate. If  $0 \in \Gamma$ , then it is a low energy point by Lemma 15.2 and therefore the statement follows from Lemma 12.3. Now, invoking Lemma 15.4, we conclude

$$0 \leq u(\rho x) \leq \frac{\rho^2}{R^2} u(Rx), \quad x \in B_1, \quad R > \rho,$$

i.e.  $u \geq 0$  everywhere in  $\mathbb{R}^n$ . Then we invoke Theorem 15.1.

In the case of Problem B, we observe that the set  $\{u \leq 0\}$  is star-like and therefore connected. Let now use the structure of  $\Omega^c$ . If  $\Omega_i^c$ ,  $i = 1, \dots, N$  are the components as in Lemma 15.3 and  $u = c_i$  there, then  $c_i \leq 0$  by the assumption (15.4). On the other hand since,  $u$  is subharmonic, we must have either  $u = 0$  in the interior of  $\{u \leq 0\}$  or  $u < 0$ . The latter is impossible, since it will imply that  $c_i < 0$  for all  $i = 1, \dots, N$  (recall that  $\Omega_i^c$  have nonempty interiors), which contradicts (15.4). Therefore we must have

$$\{u \leq 0\} = \{u = 0\}.$$

and we arrive at the situation of Problem A.  $\square$

15.2.2. *Global solutions with unbounded  $\Omega^c$ .*

**Theorem 15.6.** *Let  $u \in P_\infty(M)$  such that  $\Omega^c$  is unbounded and has nonempty interior. Then, there is a  $a \in \mathbb{R}$  such that  $u \geq a$  and  $\Omega^c = \{u = a\}$ .*

*In particular, by Theorem 15.1  $\Omega^c$  is convex.*

*Proof.* Suppose that some shrink-down  $u_\infty$  of  $u$  at 0 is a half space solution. Then, arguing as in the proof of Lemma 15.2 we will have that  $u - u(0)$  is a half space solution. Hence the theorem follows in this case.

Now, if no shrink-down is a half-space solution, we may assume  $u_\infty$  is a polynomial. The assumption  $\text{Int } \Omega^c \neq \emptyset$  prevents  $u$  from being a polynomial.

Since  $\Omega^c$  is unbounded, there exists a sequence  $x_j \in \partial\Omega$  tending to  $\infty$ . In this case we may scale by  $R_j = |x_j|$  so as to obtain, in the limit, a global solution with a free boundary point  $e$  on the unit sphere. By homogeneity then the ray  $\{re : r > 0\}$ , must lie in the free boundary. Since  $u_\infty$  is a homogeneous quadratic polynomial, this is possible only if  $\partial_e u_\infty \equiv 0$ . Consider now the Alt-Caffarelli-Friedman monotonicity functional

$$\phi_e(r, u) := \Phi(r, (\partial_e u)^+, (\partial_e u)^-).$$

Since  $(\partial_e u)^\pm$  are subharmonic by Lemma 6.2, from ACF monotonicity formula we have that

$$0 \leq \phi_e(r, u) \leq \phi_e(\infty, u) = \phi_e(1, u_\infty) = 0.$$

Hence, either  $(\partial_e u)^+$  or  $(\partial_e u)^-$  must vanish identically and we may assume without loss of generality that  $\partial_e u \geq 0$  (otherwise we replace  $e$  by  $-e$ ).

Next, without loss of generality assume  $e = e_n$  and (after changing the origin) that

$$B_r(0) \subset \Omega^c(u).$$

Then  $u \equiv u(0)$  in  $B_r(0)$ . Moreover, by monotonicity in the direction  $e_n$  we have that  $u \leq u(0)$  in the half-infinite cylinder  $B'_r(0) \times (-\infty, 0)$ , where  $B'_r(0)$  stands for a ball in  $\mathbb{R}^{n-1}$ . Since  $u$  is subharmonic, the maximum principle implies now that

$$u(x', x_n) = u(0) \quad \text{for } x' \in B'_r(0), x_n \leq 0.$$

Define now a  $(n-1)$ -dimensional solution

$$\hat{u}(x') = \lim_{x_n \rightarrow -\infty} u(x', x_n)$$

First, we notice that the limit exists by the monotonicity in the direction  $e_n$ . Next, the limit is finite, since  $B'_r(0) \times (-\infty, 0] \subset \Omega^c$  which gives the estimate

$$|u(x) - u(0)| \leq \frac{M}{2} |x'|^2.$$

Thus,  $\hat{u}$  is a  $(n-1)$ -dimensional solution with a quadratic growth at infinity. Also note that

$$B'_r(0) \subset \Omega^c(\hat{u}).$$

First, suppose that  $\hat{u}$  is either a half space solution, or falls into the hypotheses of Theorem 15.5. Then  $\hat{u}$  is convex and non-negative. Since  $u(x', x_n) - u(0) \geq \hat{u}(x') \geq 0$  we conclude the proof by applying Theorem 15.1 to  $u(x) - u(0)$ .

Next, if the lower dimensional solution  $\hat{u}$  is neither of the above it must fall into the third category analyzed above. Hence we repeat our argument and translate  $\hat{u}$  again in a new direction and reduce the dimension further. Finally, by induction, we need to classify the one dimensional solutions. However, the only one-dimensional solutions are  $x_1^2/2$ ,  $(x_1^+)^2/2$ , or two separated solutions of the latter, which are all nonnegative.  $\square$

### 15.3. Problem C.

**Theorem 15.7.** *Let  $u \in P_\infty(M)$  be a solution of Problem C such that the origin is a branching point, i.e.  $0 \in \partial\{u > 0\} \cap \partial\{u < 0\}$  and  $|\nabla u(0)| = 0$ . Then  $u$  is a two-plane solution*

$$u(x) = \frac{\lambda_+}{2} (x \cdot e)_+^2 - \frac{\lambda_-}{2} (x \cdot e)_-^2$$

for a certain direction  $e$ .

*Proof.* The proof follows from the classification of homogeneous global solutions in Theorem 11.1 and the following shrink-down argument.

Consider a limit

$$u_\infty(x) = \lim_{R_j \rightarrow \infty} \frac{u(R_j x)}{R_j^2}$$

over a certain sequence  $R_j \rightarrow \infty$ . Then Weiss's monotonicity formula implies that  $u_\infty$  is a homogeneous global solution. Since we still have that 0 is a branching

free-boundary point for  $u_\infty$ , Theorem 11.1 implies that  $u_\infty$  is a two-plane solution for a certain direction  $e$ .

In particular, for any direction  $\nu$ ,  $\partial_\nu u_\infty$  does not change sign in  $\mathbb{R}^n$  and therefore

$$\phi_\nu(r, u_\infty) = \Phi(r, (\partial_\nu u_\infty)^+, (\partial_\nu u_\infty)^-) = 0.$$

On the other hand, by the ACF monotonicity formula we have

$$0 \leq \phi_\nu(r, u) \leq \phi_\nu(\infty, u) = \phi_\nu(1, u_\infty) = 0,$$

implying  $\partial_\nu u$  does not change sign. Since this holds for all directions  $\nu$  we conclude  $u$  is one dimensional and hence can be computed as earlier.  $\square$