

A NOTABLE FAMILY OF ENTIRE INTRINSIC MINIMAL GRAPHS IN THE HEISENBERG GROUP WHICH ARE NOT PERIMETER MINIMIZING

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ABSTRACT. One of the main objectives of this paper is to unravel a new interesting phenomenon of the sub-Riemannian Bernstein problem with respect to its Euclidean ancestor, with the purpose of also indicating a possible line of attack toward its solution. We show that the global intrinsic graphs (1.2) are unstable critical points of the horizontal perimeter. As a consequence of this fact, the study of the stability acquires a central position in the problem itself.

1. Introduction

The development of geometric measure theory in sub-Riemannian spaces has received a strong impulse over the past decade, see [Pa1], [Pa2], [CDG], [KR], [E1], [E2], [E3], [Gro], [GN], [Be], [DS], [DGN1], [AK1], [AK2], [CS1], [A], [FSS1], [Ma1], [FSS2], [Ma2], [CMS], [FSS3], [BRS], [DGN4], [DGN5], [LR], [LM], [FSS4], [Ma3], [CS2] [P1], [P2], [GP], [CG], [CHMY], [CH], [HP1], [HP2], [RR], [BC], [Se1], [Se2], [Mo]. In particular, the papers [GP], [CHMY] and [CH] contain a detailed study of the Bernstein problem in the first Heisenberg group \mathbb{H}^1 , and in more general CR manifolds of real dimension three. Despite the progress made in these latter papers, this problem presently still constitutes a basic open question. One of the main objectives of this paper is to unravel a new interesting phenomenon of the sub-Riemannian Bernstein problem with respect to its Euclidean ancestor, with the purpose of also indicating a possible line of attack toward its solution.

To provide the reader with some perspective, we recall that in the Heisenberg group \mathbb{H}^n a basic discovery of Franchi, Serapioni and Serra Cassano is a structure theorem à la De Giorgi for sets of locally finite horizontal perimeter [FSS1] (see [FSS3] for an extension to Carnot groups of step $r = 2$). To prove the latter they show that the non-isotropic blow-up of such a set at a point of its reduced boundary produces a vertical hyperplane

$$(1.1) \quad P_\gamma = \{(x, y, t) \in \mathbb{H}^n \mid \langle a, x \rangle + \langle b, y \rangle = \gamma\}, \quad a^2 + b^2 \neq 0,$$

(when $\gamma = 0$ these sets are also the maximal subgroups of \mathbb{H}^n). Recalling that the characteristic locus of a hypersurface $\mathcal{S} \subset \mathbb{H}^n$, denoted henceforth by $\Sigma(\mathcal{S})$, is the collection of all points $g \in \mathcal{S}$ at which $T_g\mathcal{S} = H_g\mathbb{H}^n$, where $H\mathbb{H}^n$ denotes the horizontal bundle of \mathbb{H}^n , it is easy to recognize that $\Sigma(P_\gamma) = \emptyset$ for any $\gamma \in \mathbb{R}$. Therefore, in analogy with the Euclidean situation, the cited blow-up result from [FSS1] suggests the natural conjecture that if $\mathcal{S} \subset \mathbb{H}^n$ is a C^2 entire H -minimal graph over some hyperplane, and if $\Sigma(\mathcal{S}) = \emptyset$, then $\mathcal{S} = P_\gamma$ for some γ . Here, following a perhaps unfortunate tradition of the classical situation, H -minimal is intended in the sense that \mathcal{S} is of class C^2 , and the horizontal mean curvature, defined in (2.13) below, vanishes identically as a continuous function on \mathcal{S} . However, for \mathbb{H}^n the situation is very different

Date: September 8, 2006.

Key words and phrases. H -minimal surfaces. Intrinsic graphs. First and second variation formulas.

First author supported in part by NSF CAREER Grant, DMS-0239771.

Second author supported in part by NSF Grant DMS-0300477.

than in Euclidean space. In fact, it was proved in [GP] that there exist non-planar entire H -minimal graphs with empty characteristic locus, thus violating the above plausible conjecture. For instance, the non-planar real analytic surfaces

$$(1.2) \quad \mathcal{S} = \{(x, y, t) \in \mathbb{H}^1 \mid x = y(\alpha t + \beta), \alpha > 0, \beta \in \mathbb{R}\},$$

are H -minimal, and they have empty characteristic locus. Moreover, the surfaces (1.2) are global intrinsic X_1 -graphs in the sense of Franchi, Serapioni and Serra Cassano, see [FSS4] (we stress that intrinsic graphs have empty characteristic locus by definition). This means that there exists a globally defined function $\phi : \mathbb{R}_{u,v}^2 \rightarrow \mathbb{R}$ such that, in the coordinates (x, y, t) , we can parameterize (1.2) as follows

$$(1.3) \quad (x, y, t) = (0, u, v) \circ \phi(u, v)e_1 = (0, u, v) \circ (\phi(u, v), 0, 0) = \left(\phi(u, v), u, v - \frac{u}{2}\phi(u, v) \right),$$

where \circ indicates the non-Abelian group multiplication in \mathbb{H}^1 , see (2.1). In fact, imposing the defining equation $x = y(\alpha t + \beta)$ for \mathcal{S} , with $\alpha > 0$, we obtain the function

$$(1.4) \quad \phi(u, v) = \frac{2u(\alpha v + \beta)}{2 + \alpha u^2},$$

which describes \mathcal{S} as an entire X_1 -graph. We also note that a vertical plane P_γ is a global X_1 -graph if $a \neq 0$ (or an X_2 -graph if $b \neq 0$), with

$$(1.5) \quad \phi(u, v) = \frac{\gamma - bu}{a}, \quad \text{if } a \neq 0.$$

Examples such as (1.2) seem to cast a dim light on the Bernstein problem in \mathbb{H}^1 . There is however a deeper aspect of the problem which has gone unnoticed so far. What could be happening in fact is that, due to the different nature of the relevant perimeter functional, global intrinsic graphs such as (1.2) are only stationary for the horizontal perimeter, but not stable. In this paper we examine this aspect in depth. More precisely, we recall that, thanks to the convexity of the area functional

$$\mathcal{A}(u) = \int_{\Omega} \sqrt{1 + |Du|^2} \, dx, \quad \Omega \subset \mathbb{R}^n,$$

in the classical theory of minimal surfaces any critical point of \mathcal{A} is automatically stable, i.e., it is a local minimizer, see e.g. [CM]. By contrast, we show that the global intrinsic graphs (1.2) are unstable critical points of the horizontal perimeter $P_H(\mathcal{S})$ defined in (2.22) below. As a consequence of this fact, the study of the stability acquires a central position in the problem itself, and our results suggest that, if properly understood from this new perspective, the Bernstein property is still true in \mathbb{H}^1 . Besides their intrinsic interest, we believe that the relevance of our results lies in the method of proof, which is quite general and flexible, and has the potential of being successfully applied to attack the sub-Riemannian Bernstein problem. In Geometric Measure Theory there exist in essence two (different, but equivalent) approaches to stability: the former is based on the so-called method of calibrations, the latter on second variation formulas. Our approach revolves around a general second variation formula established in [DGN3], and on the explicit construction of a continuum of directions along which the intrinsic perimeter strictly decreases. We need to introduce a basic definition.

Definition 1.1. *We say that a C^2 oriented H -minimal surface $\mathcal{S} \subset \mathbb{H}^1$, with $\Sigma(\mathcal{S}) = \emptyset$, is stationary if it has vanishing first variation of the H -perimeter, i.e., if*

$$\mathcal{V}_I^H(\mathcal{S}; \mathcal{X}) \stackrel{\text{def}}{=} \left. \frac{d}{d\lambda} P_H(\mathcal{S}^\lambda) \right|_{\lambda=0} = 0,$$

for any deformation $\mathcal{S} \rightarrow \mathcal{S}^\lambda = \mathcal{S} + \lambda\mathcal{X}$, where $\mathcal{X} \in C_0^2(\mathcal{S}, \mathbb{H}^1)$, with $\mathcal{X} \neq 0$. We say that a stationary \mathcal{S} is stable if the second variation is nonnegative, i.e.,

$$\mathcal{V}_{II}^H(\mathcal{S}; \mathcal{X}) \stackrel{\text{def}}{=} \left. \frac{d^2}{d\lambda^2} P_H(\mathcal{S}^\lambda) \right|_{\lambda=0} \geq 0 ,$$

for any $\mathcal{X} \in C_0^2(\mathcal{S}, \mathbb{H}^1)$, with $\mathcal{X} \neq 0$. If there exists such an \mathcal{X} for which $\mathcal{V}_{II}^H(\mathcal{S}; \mathcal{X}) < 0$, then we say that \mathcal{S} is unstable.

It has been proved in [DGN3] that \mathcal{S} is stationary if and only if \mathcal{S} is H -minimal, see Theorem 3.1 below. In this paper we establish the following result.

Theorem 1.2. *For every $\alpha > 0$ and $\beta \in \mathbb{R}$, the H -minimal global intrinsic X_1 -graphs*

$$(1.6) \quad \mathcal{S} = \{(x, y, t) \in \mathbb{H}^1 \mid x = y(\alpha t + \beta)\} ,$$

are unstable. More precisely, there exist $a \in C_0^\infty(\mathcal{S})$, $a \neq 0$, and $h \in C_0^\infty(\mathcal{S})$, $h \neq 0$, such that with either $\mathcal{X} = aX_1$, or $\mathcal{X} = h\nu_H$, we have $\mathcal{V}_{II}^H(\mathcal{S}; \mathcal{X}) < 0$.

To explain the strategy behind Theorem 1.2 we mention that in Section 2 we collect some preliminary material which constitutes the geometric backbone of the paper. The novel part of the paper is contained in Section 3. An essential ingredient in the proof of Theorem 1.2 is the second variation formula, see Theorem 3.2 below. Combining the latter with some basic sub-Riemannian integration by parts formulas, see Lemmas 3.5 and 3.6, we reduce the study of the stability of (1.2) to checking the validity of some Hardy type inequalities on \mathcal{S} , see Lemmas 3.7 and 3.9. Using the representation of \mathcal{S} as a global graph, we then pull back such Hardy inequalities to ones onto the (y, t) -plane, see Lemma 3.12. Finally, in Lemma 3.13 and Corollary 3.14 we explicitly construct the directions such that, deforming the surface (1.2) along them, the H -perimeter strictly decreases. This establishes the instability of (1.2).

We would like to close this introduction with some conjectures which are suggested by the present work. Consider a C^2 , H -minimal intrinsic X_1 -graph $\mathcal{S} \subset \mathbb{H}^1$. Denoting by \mathcal{B}_ϕ the linearized Burger's operator, whose action on a function $F = F(u, v)$ is given by $\mathcal{B}_\phi(F) = F_u + \phi F_v$, then it was proved in Theorem 1.2 in [ASV] that, provided that $\phi \in C_0^2(\mathbb{R}^2)$, the H -perimeter of \mathcal{S} is given by

$$(1.7) \quad P_H(\mathcal{S}) = \int_{\text{supp}(\phi)} \sqrt{1 + \mathcal{B}_\phi(\phi)^2} \, du \wedge dv .$$

Now, if we think of (1.7) as a functional $P_H(\phi)$ acting on ϕ , one can easily recognize that, given $\zeta \in C_0^\infty(\mathbb{R}^2)$, then the first variation of $P_H(\mathcal{S})$ with respect to the deformation $\mathcal{S} \rightarrow \mathcal{S}^\lambda = \mathcal{S} + \lambda\mathcal{X}$, with $\mathcal{X} = \zeta X_1$, is given by

$$(1.8) \quad \mathcal{V}_I^H(\mathcal{S}; \mathcal{X}) \stackrel{\text{def}}{=} \left. \frac{dP_H(\mathcal{S}^\lambda)}{d\lambda} \right|_{\lambda=0} = \int_{\mathbb{R}^2} \frac{\mathcal{B}_\phi(\phi)}{\sqrt{1 + \mathcal{B}_\phi(\phi)^2}} (\zeta_u + \phi \zeta_v + \zeta \phi_v) \, du \wedge dv .$$

In view of Theorem 1.2, it is natural to make the following conjecture: *Suppose that ϕ , belonging to a suitable Sobolev space, is a local (or even a global) minimizer of (1.7), and therefore in particular also a critical point, then after modification on a set of measure zero, ϕ must be of the type (1.5).*

If we assume a priori that $\phi \in C^2(\mathbb{R}^2)$, then integrating by parts in (1.8) we obtain that

$$(1.9) \quad \mathcal{V}_I^H(\mathcal{S}; \mathcal{X}) = - \int_{\mathbb{R}^2} \zeta \mathcal{B}_\phi \left(\frac{\mathcal{B}_\phi(\phi)}{\sqrt{1 + \mathcal{B}_\phi(\phi)^2}} \right) \, du \wedge dv , \quad \zeta \in C_0^\infty(\mathbb{R}^2) .$$

On the other hand, under the same regularity hypothesis on ϕ one can recognize, see [GS], that

$$(1.10) \quad \mathcal{B}_\phi \left(\frac{\mathcal{B}_\phi(\phi)}{\sqrt{1 + \mathcal{B}_\phi(\phi)^2}} \right) = -\mathcal{H} ,$$

where \mathcal{H} represents the H -mean curvature of \mathcal{S} , defined in (2.13) below. Therefore, a C^2 global intrinsic graph \mathcal{S} is a critical point of (1.7) if and only if \mathcal{S} is H -minimal. As a consequence of these considerations, in the C^2 framework the above conjecture could be reformulated by saying that: *The only C^2 , stable, global intrinsic graphs in \mathbb{H}^1 are the vertical planes.* Finally, we would also like to return to the conjecture in the opening of this introduction and amend it as follows: *In \mathbb{H}^1 the vertical planes (1.1) are the only C^2 , stable, entire H -minimal graphs (over some plane).*

Acknowledgment: The problems treated in this paper were inspired by some stimulating discussions with F. Serra Cassano and R. Serapioni during a visit of the second named author at the University of Trento in April 2005. He would like to thank them for their gracious hospitality. The authors would also like to thank the anonymous referee for his/her careful reading of the manuscript and for some comments which helped to improve the presentation of the paper.

2. Preliminary material

In this section we introduce some relevant notation and definitions from [DGN3] which will be used in the proof of Theorem 1.2. We consider the first Heisenberg group $\mathbb{H}^1 = (\mathbb{R}^3, \circ)$ with group law

$$(2.1) \quad (x, y, t) \circ (x', y', t') = \left(x + x', y + y', t + t' + \frac{1}{2}(xy' - x'y) \right) ,$$

and non-isotropic dilations $\delta_\lambda(x, y, t) = (\lambda x, \lambda y, \lambda^2 t)$, see [S]. Hereafter, we will use the letters $g = (x, y, t)$, $g' = (x', y', t')$, etc., to indicate points in \mathbb{H}^1 . Denoting with $(L_g)_*$ the differential of the left-translation operator $L_g : \mathbb{H}^1 \rightarrow \mathbb{H}^1$ defined by $L_g(g') = g \circ g'$, and letting e_i , $i = 1, 2, 3$, indicate the standard orthonormal basis of \mathbb{R}^3 , one readily verifies that

$$(2.2) \quad \begin{cases} X_1(g) \stackrel{def}{=} (L_g)_*(e_1) = \frac{\partial}{\partial x} - \frac{y}{2} \frac{\partial}{\partial t} , \\ X_2(g) \stackrel{def}{=} (L_g)_*(e_2) = \frac{\partial}{\partial y} + \frac{x}{2} \frac{\partial}{\partial t} , \\ T(g) \stackrel{def}{=} (L_g)_*(e_3) = \frac{\partial}{\partial t} . \end{cases}$$

The three vector fields $\{X_1, X_2, T\}$ generate the Lie algebra \mathfrak{h}^1 of all left-invariant vector fields on \mathbb{H}^1 . They satisfy at every point of \mathbb{H}^1 the non-trivial commutation relation

$$(2.3) \quad [X_1, X_2] = T ,$$

all other commutators being trivial. In view of (2.3), the Heisenberg group constitutes the first (and perhaps most important) prototype of a class of graded nilpotent Lie groups nowadays known as Carnot groups, see [Fo], [S], [Gro], [Pa2], [Be]. We observe explicitly that, if we let $V_1 = \mathbb{R}_{x,y}^2 \times \{0\}_t$, and $V_2 = \{0\}_{x,y} \times \mathbb{R}_t$, then the Heisenberg algebra admits the decomposition $\mathfrak{h}^1 = V_1 \oplus V_2$. We assume hereafter that \mathbb{H}^1 be endowed with a left-invariant Riemannian metric with respect to which $\{X_1, X_2, T\}$ constitute an orthonormal basis. The inner product with respect to this metric will be denoted by $\langle \cdot, \cdot \rangle$. This is the only inner product that will be used in this paper, therefore there will be no confusion with other inner products, such as for

instance the Euclidean one, in \mathbb{R}^3 . The corresponding Levi-Civita connection on \mathbb{H}^1 will be denoted by $\nabla_X Y$. We will denote by $H\mathbb{H}^1$ the subbundle of the tangent bundle $T\mathbb{H}^1$ generated by the distribution $\{X_1, X_2\}$. The horizontal Levi-Civita connection is given as follows. For any $X \in \Gamma(T\mathbb{H}^1)$, $Y \in \Gamma(H\mathbb{H}^1)$ we let

$$(2.4) \quad \nabla_X^H Y = \sum_{i=1}^2 \langle \nabla_X Y, X_i \rangle X_i,$$

and one can easily verify that $\nabla_X^H Y$ is metric preserving and torsion free, in the sense that if we define the horizontal torsion of \mathcal{S} as

$$T^H(X, Y) = \nabla_X^H Y - \nabla_Y^H X - [X, Y]^H,$$

where $[X, Y]^H = \sum_{i=1}^2 \langle [X, Y], X_i \rangle X_i$, then $T^H(X, Y) = 0$. Given a function $f \in C^1(\mathbb{H}^1)$, its Riemannian gradient is given by

$$\nabla f = X_1 f X_1 + X_2 f X_2 + T f T,$$

whereas its horizontal gradient is given by the projection of ∇f onto the subbundle \mathbb{H}^1 , hence

$$\nabla^H f = \langle \nabla f, X_1 \rangle X_1 + \langle \nabla f, X_2 \rangle X_2 = X_1 f X_1 + X_2 f X_2.$$

Given an oriented C^2 surface $S \subset \mathbb{H}^1$, we denote by \mathbf{N} its (non-unit) Riemannian normal with respect to the orthonormal frame $\{X_1, X_2, T\}$, and by $\nu = \mathbf{N}/|\mathbf{N}|$ its Riemannian Gauss map. We consider the quantities

$$(2.5) \quad p = \langle \mathbf{N}, X_1 \rangle, \quad q = \langle \mathbf{N}, X_2 \rangle, \quad \omega = \langle \mathbf{N}, T \rangle, \quad W = \sqrt{p^2 + q^2}.$$

In this paper we adopt the classical non-parametric point of view, see for instance [MM], according to which a C^2 surface $S \subset \mathbb{H}^1$ is a subset of \mathbb{R}^3 which locally coincides with the zero set of a real function. Thus, for every $g_0 \in S$ there exists an open set $\mathcal{O} \subset \mathbb{H}^1$ and a function $\phi \in C^2(\mathcal{O})$ such that: (i) $|\nabla \phi(g)| \neq 0$ for every $g \in \mathcal{O}$; (ii) $S \cap \mathcal{O} = \{g \in \mathcal{O} \mid \phi(g) = 0\}$. We will always assume that S is oriented in such a way that for every $g \in S$ one has

$$\mathbf{N}(g) = \nabla \phi(g) = X_1 \phi(g) X_1 + X_2 \phi(g) X_2 + T \phi(g) T.$$

We note explicitly that, in this situation, the functions p, q, ω defined in (2.5), which are given by $p = X_1 \phi$, $q = X_2 \phi$, $\omega = T \phi$, are not only defined on S , but for every $g_0 \in S$ they belong to $C^1(\mathcal{O})$. This notion of C^2 surface obviously includes the entire intrinsic graphs considered in Theorem 1.2. In fact, in the case of the surfaces S in (1.6), we have (see (3.17)),

$$p = X_1 \phi = 1 + \frac{\alpha}{2} y^2, \quad q = X_2 \phi = -\alpha t - \beta - \frac{\alpha}{2} xy, \quad \omega = T \phi = -\alpha y,$$

and thus in particular the field $\mathbf{N} = pX_1 + qX_2 + \omega T$ belongs to $C^\infty(\mathcal{O}, \mathbb{H}^1)$, with $\mathcal{O} = \mathbb{H}^1$.

We emphasize here that the local defining function ϕ in (i) and (ii) above has a different meaning from the function ϕ in the definition of intrinsic graph in the introduction. Given a surface $S \subset \mathbb{H}^1$, on the set $S \setminus \Sigma(S)$ we define the *horizontal Gauss map* by

$$(2.6) \quad \nu_H = \bar{p} X_1 + \bar{q} X_2,$$

where we have let

$$(2.7) \quad \bar{p} = \frac{p}{W}, \quad \bar{q} = \frac{q}{W}, \quad \text{so that} \quad |\nu_H|^2 = \bar{p}^2 + \bar{q}^2 \equiv 1 \quad \text{on} \quad S \setminus \Sigma(S).$$

Given a point $g_0 \in S \setminus \Sigma(S)$, the horizontal tangent space of S at g_0 is defined by

$$HT_{g_0}(S) = \{v \in H_{g_0}\mathbb{H}^1 \mid \langle v, \nu_H(g_0) \rangle = 0\}.$$

Let us notice that a basis for $HT_{g_0}(S)$ is given by the vector field

$$(2.8) \quad \nu_H^\perp = \bar{q} X_1 - \bar{p} X_2.$$

Given a function $u \in C^1(\mathcal{S})$ one clearly has $\nabla^{H,\mathcal{S}}u(g_0) \in HT_{g_0}(\mathcal{S})$. We next recall some basic definitions from [DGN3].

Definition 2.1. *Let $\mathcal{S} \subset \mathbb{H}^1$ be a C^2 surface, with $\Sigma(\mathcal{S}) = \emptyset$, then we define the horizontal connection on \mathcal{S} as follows. For every $X, Y \in C^1(\mathcal{S}; HT\mathcal{S})$ we let*

$$\nabla_X^{H,\mathcal{S}}Y = \nabla_{\bar{X}}^H\bar{Y} - \langle \nabla_{\bar{X}}^H\bar{Y}, \nu_H \rangle \nu_H ,$$

where $\bar{X}, \bar{Y} \in C^1(\mathbb{H}^1; H\mathbb{H}^1)$ are such that $\bar{X} = X, \bar{Y} = Y$ on \mathcal{S} .

Similarly to the Riemannian case, it is possible to prove that $\nabla_X^{H,\mathcal{S}}Y$ does not depend on the extensions \bar{X}, \bar{Y} . The tangential horizontal gradient of a function $f \in C^1(\mathcal{S})$ is defined as follows

$$(2.9) \quad \nabla^{H,\mathcal{S}}f = \nabla^H\bar{f} - \langle \nabla^H\bar{f}, \nu_H \rangle \nu_H ,$$

where \bar{f} denotes any extension of f to all of \mathbb{H}^1 . The definition of $\nabla^{H,\mathcal{S}}f$ is well-posed since $\nabla^{H,\mathcal{S}}f$ only depends on the values of f on \mathcal{S} . Since $|\nu_H| \equiv 1$ on $\mathcal{S} \setminus \Sigma(\mathcal{S})$, we clearly have $\langle \nabla^{H,\mathcal{S}}f, \nu_H \rangle = 0$, and therefore

$$(2.10) \quad |\nabla^{H,\mathcal{S}}f|^2 = |\nabla^H\bar{f}|^2 - \langle \nabla^H\bar{f}, \nu_H \rangle^2 .$$

All the above definitions are specializations to \mathbb{H}^1 of analogous ones for general Carnot groups, see [DGN3]. The next definition contains the essential geometric concept of horizontal second fundamental form. It is convenient to state it for \mathbb{H}^n , rather than \mathbb{H}^1 .

Definition 2.2. *Let $\mathcal{S} \subset \mathbb{H}^n$ be a C^2 hypersurface with $\Sigma(\mathcal{S}) = \emptyset$, then for every $X, Y \in C^1(\mathcal{S}; HT\mathcal{S})$ we define a tensor field of type $(0, 2)$ on \mathcal{S} , as follows*

$$(2.11) \quad II^{H,\mathcal{S}}(X, Y) = \langle \nabla_{\bar{X}}^H\bar{Y}, \nu_H \rangle \nu_H ,$$

where \bar{X}, \bar{Y} have the same meaning as in Definition 2.1. We call $II^{H,\mathcal{S}}(\cdot, \cdot)$ the horizontal second fundamental form of \mathcal{S} . We also define $\mathcal{A}^{H,\mathcal{S}} : HT\mathcal{S} \rightarrow HT\mathcal{S}$ by letting for every $g \in \mathcal{S}$ and $\mathbf{u}, \mathbf{v} \in HT_g\mathcal{S}$

$$(2.12) \quad \langle \mathcal{A}^{H,\mathcal{S}}\mathbf{u}, \mathbf{v} \rangle = - \langle II^{H,\mathcal{S}}(\mathbf{u}, \mathbf{v}), \nu_H \rangle = - \langle \nabla_{\bar{X}}^H\bar{Y}, \nu_H \rangle ,$$

where $X, Y \in C^1(\mathcal{S}, HT\mathcal{S})$ are such that $X_g = \mathbf{u}, Y_g = \mathbf{v}$, and \bar{X}, \bar{Y} are as above. We call the linear map $\mathcal{A}^{H,\mathcal{S}} : HT_g\mathcal{S} \rightarrow HT_g\mathcal{S}$ the horizontal shape operator. If $\mathbf{e}_1, \dots, \mathbf{e}_{2n-1}$ denotes a local orthonormal frame for $HT\mathcal{S}$, then the matrix of the horizontal shape operator with respect to the basis $\mathbf{e}_1, \dots, \mathbf{e}_{2n-1}$ is given by the $(2n-1) \times (2n-1)$ matrix $-\left[\langle \nabla_{\mathbf{e}_i}^H\mathbf{e}_j, \nu_H \rangle \right]_{i,j=1,\dots,2n-1}$.

If \mathcal{S} has non-empty characteristic locus $\Sigma(\mathcal{S})$, then we consider $\mathcal{S}' = \mathcal{S} \setminus \Sigma(\mathcal{S})$ and define the H -mean curvature of \mathcal{S} at a point $g_0 \in \mathcal{S}'$ as follows

$$(2.13) \quad \mathcal{H} = - \text{trace } \mathcal{A}^{H,\mathcal{S}} = - \sum_{j=1}^{2n-1} \langle \nabla_{\mathbf{e}_i}^H\mathbf{e}_j, \nu_H \rangle .$$

We recall that it was proved in [B], [Ma3] that $\mathcal{H}^{Q-1}(\Sigma(\mathcal{S})) = 0$, where \mathcal{H}^s denotes the s -dimensional Hausdorff measure associated with the horizontal or Carnot-Carathéodory distance of \mathcal{G} , and Q indicates the homogeneous dimension of \mathcal{G} . If $g_0 \in \Sigma(\mathcal{S})$ we let

$$\mathcal{H}(g_0) = \lim_{g \rightarrow g_0, g \in \mathcal{S}' \setminus \Sigma(\mathcal{S})} \mathcal{H}(g) ,$$

provided that such limit exists, finite or infinite. We do not define the H -mean curvature at those points $g_0 \in \Sigma(\mathcal{S})$ at which the limit does not exist. The following result is taken from [DGN3].

Proposition 2.3. *The H -mean curvature of $\mathcal{S} \subset \mathbb{H}^1$ coincides with the function*

$$(2.14) \quad \mathcal{H} = \sum_{i=1}^2 \nabla_i^{H,\mathcal{S}} \langle \nu_H, X_i \rangle = \nabla_1^{H,\mathcal{S}} \bar{p} + \nabla_2^{H,\mathcal{S}} \bar{q} = X_1 \bar{p} + X_2 \bar{q},$$

where \bar{p}, \bar{q} are as in (2.7).

Definition 2.4. *A C^2 surface $\mathcal{S} \subset \mathbb{H}^1$ is called H -minimal if $\mathcal{H} \equiv 0$ as a continuous function on \mathcal{S} .*

In keeping up with the notation of [DGN3] it will be convenient to indicate with $Y\zeta$ and $Z\zeta$ the respective actions of the vector fields ν_H and ν_H^\perp on a function $\zeta \in C_0^1(\mathcal{S} \setminus \Sigma(\mathcal{S}))$, thus

$$(2.15) \quad Y\zeta \stackrel{def}{=} \bar{p} X_1 \zeta + \bar{q} X_2 \zeta, \quad Z\zeta \stackrel{def}{=} \bar{q} X_1 \zeta - \bar{p} X_2 \zeta.$$

The frame $\{Z, Y, T\}$ is orthonormal. It is worth observing that, since the metric tensor $\{g_{ij}\}$ with respect to the inner product $\langle \cdot, \cdot \rangle$ has the property $\det\{g_{ij}\} = 1$, then the (Riemannian) divergence in \mathbb{H}^1 of these vector fields is given by

$$(2.16) \quad \operatorname{div} Y = X_1 \bar{p} + X_2 \bar{q} = \mathcal{H}, \quad \operatorname{div} Z = X_1 \bar{q} - X_2 \bar{p}.$$

Using Cramer's rule one easily obtains from (2.15)

$$(2.17) \quad X_1 \zeta = \bar{q} Z\zeta + \bar{p} Y\zeta, \quad X_2 \zeta = \bar{q} Y\zeta - \bar{p} Z\zeta.$$

One also has

$$(2.18) \quad \nabla_1^{H,\mathcal{S}} \zeta = \bar{q} Z\zeta, \quad \nabla_2^{H,\mathcal{S}} \zeta = -\bar{p} Z\zeta,$$

so that

$$(2.19) \quad |\nabla^{H,\mathcal{S}} \zeta|^2 = (Z\zeta)^2.$$

We notice that

$$(2.20) \quad \bar{q} Z\bar{p} - \bar{p} Z\bar{q} = \mathcal{H}.$$

This can be easily recognized using Proposition 2.3 and (2.17), as follows

$$\mathcal{H} = X_1 \bar{p} + X_2 \bar{q} = \bar{q} Z\bar{p} - \bar{p} Z\bar{q} + \bar{p} Y\bar{p} + \bar{q} Y\bar{q} = \bar{q} Z\bar{p} - \bar{p} Z\bar{q},$$

where we have used the fact that $0 = \frac{1}{2}Y(\bar{p}^2 + \bar{q}^2) = \bar{p}Y\bar{p} + \bar{q}Y\bar{q}$. Finally, we will need the following identity

$$(2.21) \quad \mathcal{H}^2 = (Z\bar{p})^2 + (Z\bar{q})^2.$$

This can be easily proved observing that (2.20) and the identity $\bar{p}^2 + \bar{q}^2 = 1$ give

$$\mathcal{H}^2 = (Z\bar{p})^2 + (Z\bar{q})^2 - (\bar{p}Z\bar{p} + \bar{q}Z\bar{q}) = (Z\bar{p})^2 + (Z\bar{q})^2.$$

In the classical theory of minimal surfaces, the concept of area or perimeter occupies a central position, see [DG1], [DG2], [DCP], [G], [MM]. In sub-Riemannian geometry there exists an appropriate notion of perimeter. Given an open set $\Omega \subset \mathbb{H}^1$ we denote $\mathcal{F}(\Omega) = \{\zeta \in C_0^1(\Omega; H\mathbb{H}^1) \mid \|\zeta\|_{L^\infty(\Omega)} \leq 1\}$. A function $u \in L^1(\Omega)$ is said to belong to $BV_H(\Omega)$ (the space of functions with finite horizontal bounded variation), if

$$\operatorname{Var}_H(u; \Omega) = \sup_{\zeta \in \mathcal{F}(\Omega)} \int_{\Omega} u \operatorname{div}_H \zeta \, dg < \infty.$$

This space becomes a Banach space with the norm $\|u\|_{BV_H(\Omega)} = \|u\|_{L^1(\Omega)} + \text{Var}_H(u; \Omega)$. Given a measurable set $E \subset \mathbb{H}^1$, the H -perimeter of E with respect to the open set $\Omega \subset \mathbb{H}^1$ is defined as follows, see for instance [CDG], and [GN],

$$P_H(E; \Omega) = \text{Var}_H(\chi_E; \Omega) .$$

Given an oriented C^2 surface \mathcal{S} , we will denote by $d\sigma_H$ the H -perimeter measure concentrated on \mathcal{S} . For any Borel subset $E \subset \mathcal{S}$ such that $P_H(E) < \infty$, one has

$$(2.22) \quad \begin{aligned} P_H(E) &= \int_E d\sigma_H = \int_E \sqrt{\langle \boldsymbol{\nu}, X_1 \rangle^2 + \langle \boldsymbol{\nu}, X_2 \rangle^2} d\sigma \\ &= \int_E \frac{\sqrt{\langle \mathbf{N}, X_1 \rangle^2 + \langle \mathbf{N}, X_2 \rangle^2}}{|\mathbf{N}|} d\sigma = \int_E \frac{W}{|\mathbf{N}|} d\sigma , \end{aligned}$$

where in the last equality we have used (2.5). We thus obtain from (2.22)

$$(2.23) \quad d\sigma_H = \frac{W}{|\mathbf{N}|} d\sigma ,$$

where $d\sigma$ denotes the standard surface measure.

3. Proof of Theorem 1.2

This section is devoted to proving Theorem 1.2. In the course of the proof we need to build some auxiliary results which we feel have an independent interest. We begin by recalling the notions of first and second variation of the H -perimeter introduced in Definition 1.1. Classical minimal surfaces are critical points of the perimeter (area functional). It is natural to ask what is the connection between the notion of H -minimal surface and that of H -perimeter. The answer to this question is contained in the following result from [DGN3], see also [DGN2].

Theorem 3.1. *Let $\mathcal{S} \subset \mathbb{H}^1$ be an oriented C^2 surface with $\Sigma(\mathcal{S}) = \emptyset$, then*

$$(3.1) \quad \mathcal{V}_I^H(\mathcal{S}; \mathcal{X}) = \int_{\mathcal{S}} \mathcal{H} \frac{\cos(\mathcal{X} \angle \mathbf{N})}{\cos(\boldsymbol{\nu}_H \angle \mathbf{N})} |\mathcal{X}| d\sigma_H ,$$

where \angle denotes the angle between vectors in the inner product $\langle \cdot, \cdot \rangle$. In particular, \mathcal{S} is stationary if and only if it is H -minimal.

We emphasize that, thanks to the assumption $\Sigma(\mathcal{S}) = \emptyset$, the denominator in the integrand in the right-hand side of (3.1) does not vanish on \mathcal{S} . We mention that versions of Theorem 3.1 have also been obtained independently by other people. An approach based on motion by H -mean curvature can be found in [BC]. When $\mathcal{X} = h\boldsymbol{\nu}_H$, then a proof based on CR-geometry can be found in [CHMY].

A central (and more complex) result for this paper is the following theorem established in [DGN3]. Recalling the function ω defined in (2.5), henceforth we let $\bar{\omega} = \omega/W$.

Theorem 3.2. *Let $\mathcal{S} \subset \mathbb{H}^1$ be a C^2 oriented surface with $\Sigma(\mathcal{S}) = \emptyset$. The second variation of the H -perimeter with respect to the deformation $\mathcal{S} \rightarrow \mathcal{S}^\lambda = \mathcal{S} + \lambda\mathcal{X}$, with $\mathcal{X} = aX_1 + bX_2 + kT \in$*

$C_0^2(\mathcal{S}; \mathbb{H}^1)$, is given by the formula

$$(3.2) \quad \begin{aligned} \mathcal{V}_{II}^H(\mathcal{S}; \mathcal{X}) = & \int_{\mathcal{S}} \left\{ -2(\bar{p}Zb - \bar{q}Za)(Tk - \bar{\omega}Yk) \right. \\ & + (Ta - \bar{\omega}Ya) \left[-2\bar{q}Zk - \bar{q}(a\bar{p} + b\bar{q}) - \bar{p}(a\bar{q} - b\bar{p}) \right] \\ & + (Tb - \bar{\omega}Yb) \left[2\bar{p}Zk + \bar{p}(a\bar{p} + b\bar{q}) - \bar{q}(a\bar{q} - b\bar{p}) \right] \\ & + 2(a\bar{q} - b\bar{p})(\bar{q}Za - \bar{p}Zb)\bar{\omega} \\ & + (Za + \bar{\omega}\bar{p}Zk)^2 + (Zb + \bar{\omega}\bar{q}Zk)^2 \\ & + (a^2 + b^2)\bar{\omega}^2 \\ & + 2\bar{\omega}(aZa + bZb) + 2\bar{\omega}^2(a\bar{p} + b\bar{q})Zk \\ & \left. - (\bar{q}Za - \bar{p}Zb + (a\bar{q} - b\bar{p})\bar{\omega})^2 \right\} d\sigma_H. \end{aligned}$$

Corollary 3.3. *If we choose $a \in C_0^\infty(\mathcal{S})$, $b \equiv k \equiv 0$, and therefore $\mathcal{X} = aX_1$, then the corresponding second variation of the H -perimeter is given by*

$$(3.3) \quad \begin{aligned} \mathcal{V}_{II}^H(\mathcal{S}; \mathcal{X}) = & \int_{\mathcal{S}} \left\{ \bar{p}^2(Za)^2 + \bar{p}^2\bar{\omega}^2 a^2 \right. \\ & \left. + \bar{\omega}Z(a^2) - \bar{p}\bar{q}(T(a^2) - \bar{\omega}Y(a^2)) \right\} d\sigma_H. \end{aligned}$$

Proof. It follows in an elementary fashion from (3.2). One only needs to keep in mind that $\bar{p}^2 + \bar{q}^2 = 1$. □

Corollary 3.4. *Given a C^3 oriented surface $\mathcal{S} \subset \mathbb{H}^1$, with $\Sigma(\mathcal{S}) = \emptyset$, consider the deformation*

$$(3.4) \quad \mathcal{S}^\lambda = \mathcal{S} + \lambda(h\nu_H + kT), \quad h, k \in C_0^2(\mathcal{S}),$$

corresponding to the choice $\mathcal{X} = \bar{p}hX_1 + \bar{q}hX_2$ (notice that $a = \bar{p}h, b = \bar{q}h \in C_0^2(\mathcal{S})$). One has

$$(3.5) \quad \begin{aligned} \mathcal{V}_{II}^H(\mathcal{S}; \mathcal{X}) = & \int_{\mathcal{S}} \left(Zh + \bar{\omega}Zk \right)^2 d\sigma_H \\ & + 2 \int_{\mathcal{S}} h \mathcal{H} (Tk - \bar{\omega}Yk) d\sigma_H \\ & + \int_{\mathcal{S}} \left\{ \bar{\omega}Z(h^2) + 2\mathcal{A}hZk + \mathcal{A}h^2 \right\} d\sigma_H, \end{aligned}$$

where we have set

$$(3.6) \quad \mathcal{A} = (\bar{p}T\bar{q} - \bar{q}T\bar{p}) + \bar{\omega}(\bar{q}Y\bar{p} - \bar{p}Y\bar{q}) + \bar{\omega}^2.$$

Proof. We notice that we presently have

$$\begin{aligned} \bar{p}a + \bar{q}b &= h, & \bar{q}a - \bar{p}b &= 0, \\ aZa + bZb &= hZh, \\ \bar{p}Zb - \bar{q}Za &= h(\bar{p}Z\bar{q} - \bar{q}Z\bar{p}) = -h\mathcal{H}, \end{aligned}$$

where in the last equality we have used the identity (2.20). We also have

$$\begin{aligned}
Ta - \bar{\omega}Ya &= \bar{p}(Th - \bar{\omega}Yh) + h(T\bar{p} - \bar{\omega}Y\bar{p}) , \\
Tb - \bar{\omega}Yb &= \bar{q}(Th - \bar{\omega}Yh) + h(T\bar{q} - \bar{\omega}Y\bar{q}) , \\
(Za + \bar{\omega}\bar{p}Zk)^2 &= \bar{p}^2(Zh + \bar{\omega}Zk)^2 \\
&\quad + (Z\bar{p})^2 h^2 + 2\bar{p}Z\bar{p}h(Zh + \bar{\omega}Zk) , \\
(Zb + \bar{\omega}\bar{q}Zk)^2 &= \bar{q}^2(Zh + \bar{\omega}Zk)^2 \\
&\quad + (Z\bar{q})^2 h^2 + 2\bar{q}Z\bar{q}h(Zh + \bar{\omega}Zk) .
\end{aligned}$$

We next observe that

Substituting these formulas in the right-hand side of (3.2), and using (2.21), we reach the desired conclusion. \square

A different approach to a version of (3.5) based on CR-geometry was found in [CHMY]. To reduce further the expressions in the right-hand side of (3.3), (3.5) we would like to transform the terms containing the derivatives $Z(a^2)$, $T(a^2)$, $Y(a^2)$, $Z(h^2)$. For this, we will use the following basic integration by parts formulas proved in [DGN3].

Lemma 3.5. *Let $\mathcal{S} \subset \mathbb{H}^1$ be a C^2 oriented surface with $\Sigma(\mathcal{S}) = \emptyset$. For any $\zeta \in C_0^1(\mathcal{S})$ one has*

$$\int_{\mathcal{S}} Z\zeta \, d\sigma_H = - \int_{\mathcal{S}} \zeta \bar{\omega} \, d\sigma_H .$$

Lemma 3.6. *With \mathcal{S} as in Lemma 3.5, for any function $\zeta \in C_0^1(\mathcal{S})$ one has*

$$\int_{\mathcal{S}} T\zeta \, d\sigma_H = \int_{\mathcal{S}} Y\zeta \bar{\omega} \, d\sigma_H + \int_{\mathcal{S}} \zeta \bar{\omega} \mathcal{H} \, d\sigma_H .$$

Using Lemma 3.5 we find

$$\begin{aligned}
(3.7) \quad \int_{\mathcal{S}} \bar{\omega}Z(a^2) \, d\sigma_H &= \int_{\mathcal{S}} Z(\bar{\omega}a^2) \, d\sigma_H - \int_{\mathcal{S}} a^2 Z\bar{\omega} \, d\sigma_H \\
&= - \int_{\mathcal{S}} a^2 \bar{\omega}^2 \, d\sigma_H - \int_{\mathcal{S}} a^2 Z\bar{\omega} \, d\sigma_H .
\end{aligned}$$

From Lemma 3.6 we obtain instead

$$\begin{aligned}
(3.8) \quad & - \int_{\mathcal{S}} \bar{p}\bar{q}(T(a^2) - \bar{\omega}Y(a^2)) \, d\sigma_H \\
&= \int_{\mathcal{S}} a^2 \left\{ (\bar{p}T\bar{q} + \bar{q}T\bar{p}) - \bar{\omega}(\bar{p}Y\bar{q} + \bar{q}Y\bar{p}) - \bar{p}\bar{q}\bar{\omega}\mathcal{H} \right\} \, d\sigma_H .
\end{aligned}$$

Substituting (3.7), (3.8) into (3.3), and keeping (2.9) in mind, we finally obtain.

Lemma 3.7. *Let $\mathcal{S} \subset \mathbb{H}^1$ be a C^2 , oriented surface, with $\Sigma(\mathcal{S}) = \emptyset$, then the second variation of the H -perimeter, with respect to deformation $\mathcal{S}^\lambda = \mathcal{S} + \lambda aX_1$, is given by*

$$\begin{aligned}
(3.9) \quad \mathcal{V}_{II}^H(\mathcal{S}; aX_1) &= \int_{\mathcal{S}} \bar{p}^2 |\nabla^{H,\mathcal{S}} a|^2 \, d\sigma_H + \int_{\mathcal{S}} a^2 \left\{ (\bar{p}T\bar{q} + \bar{q}T\bar{p}) - \bar{\omega}(\bar{p}Y\bar{q} + \bar{q}Y\bar{p}) \right. \\
&\quad \left. - \bar{q}^2 \bar{\omega}^2 - Z\bar{\omega} - \bar{p}\bar{q}\bar{\omega}\mathcal{H} \right\} \, d\sigma_H .
\end{aligned}$$

To establish the next lemma we need the following auxiliary result.

Lemma 3.8. *On a surface \mathcal{S} as in Lemma 3.7, one has*

$$- Z\bar{\omega} = \mathcal{A} ,$$

where \mathcal{A} is the quantity defined in (3.6).

Proof. From the definition of $\bar{\omega}$ one has

$$(3.10) \quad - Z\bar{\omega} = \bar{\omega} \frac{ZW}{W} - \frac{Z\omega}{W} .$$

We now claim that

$$(3.11) \quad \frac{Z\omega}{W} = \bar{q} T\bar{p} - \bar{p} T\bar{q} ,$$

and that, furthermore,

$$(3.12) \quad \frac{ZW}{W} = \bar{q} Y\bar{p} - \bar{p} Y\bar{q} + \bar{\omega} .$$

It should be obvious to the reader that, inserting (3.11), (3.12) into (3.10), we obtain the desired conclusion. We are thus left with proving (3.11) and (3.12). For the former, we observe that

$$\langle Z, \mathbf{N} \rangle = 0 .$$

If ϕ denotes a local defining function of \mathcal{S} in the neighborhood of an arbitrary point, we thus have $Z\phi = 0$. Applying T to this identity, we obtain

$$\begin{aligned} 0 = T(Z\phi) &= T(\bar{q}X_1\phi - \bar{p}X_2\phi) = T\bar{q}X_1\phi + \bar{q}TX_1\phi - T\bar{p}X_2\phi - \bar{p}TX_2\phi \\ &= T\bar{q}X_1\phi - T\bar{p}X_2\phi + \bar{q}X_1T\phi - \bar{p}X_2T\phi = pT\bar{q} - qT\bar{p} + Z(T\phi) , \end{aligned}$$

where we have used $[X_i, T] = 0$, $i = 1, 2$. It follows that

$$\frac{Z\omega}{W} = \frac{Z(T\phi)}{W} = \bar{q}T\bar{p} - \bar{p}T\bar{q} ,$$

which proves (3.11). As for (3.12), we have

$$\omega = T\phi = X_1X_2\phi - X_2X_1\phi = X_1(\bar{q}W) - X_2(\bar{p}W) = -(X_2\bar{p} - X_1\bar{q})W + ZW ,$$

from which the desired conclusion follows immediately. □

Using Lemma 3.5 in (3.5) of Corollary 3.4, in combination with Lemma 3.8, we obtain.

Lemma 3.9. *Let $\mathcal{S} \subset \mathbb{H}^1$ be a C^3 , oriented surface, with $\Sigma(\mathcal{S}) = \emptyset$, then the second variation of the H -perimeter with respect to the deformation of $\mathcal{S}^\lambda = \mathcal{S} + \lambda h\nu_H$, is given by*

$$(3.13) \quad \begin{aligned} \mathcal{V}_{II}^H(\mathcal{S}; h\nu_H) &= \int_{\mathcal{S}} \left\{ (Zh)^2 + h^2 (2\mathcal{A} - \bar{\omega}^2) \right\} d\sigma_H \\ &= \int_{\mathcal{S}} \left\{ (Zh)^2 + h^2 [2(\bar{p}T\bar{q} - \bar{q}T\bar{p}) + 2\bar{\omega}(\bar{q}Y\bar{p} - \bar{p}Y\bar{q}) + \bar{\omega}^2] \right\} d\sigma_H , \end{aligned}$$

where \mathcal{A} is defined in (3.6).

After these preparations we turn to the core of the proof of Theorem 1.2. We will focus on the case in which the surface is given by

$$(3.14) \quad \mathcal{S} = \{(x, y, t) \in \mathbb{H}^1 \mid x = y(\alpha t + \beta)\},$$

the other family of surfaces $\mathcal{S} = \{(x, y, t) \in \mathbb{H}^1 \mid y = x(\alpha t + \beta)\}$, with $\alpha < 0$ and $\beta \in \mathbb{R}$, being treated by completely analogous considerations. Our first step in the proof of Theorem 1.2 is to compute the second variation of the H -perimeter for \mathcal{S} . In view of (3.9) in Lemma 3.7, or (3.13) in Lemma 3.9, we need to compute the quantities which appear as the coefficient of a^2 and h^2 in the integral in the right-hand side of the respective formulas. This is the content of the next lemma.

Lemma 3.10. *Let \mathcal{S} be the H -minimal surface given by (3.14), then one has*

$$(3.15) \quad \begin{aligned} & (\bar{p}T\bar{q} + \bar{q}T\bar{p}) - \bar{\omega}(\bar{p}Y\bar{q} + \bar{q}Y\bar{p}) - \bar{q}^2\bar{\omega}^2 \\ & - Z\bar{\omega} - \bar{p}\bar{q}\bar{\omega}\mathcal{H} = - \frac{2\alpha}{W^2(1 + (\alpha t + \beta)^2)}, \end{aligned}$$

$$(3.16) \quad 2(\bar{p}T\bar{q} - \bar{q}T\bar{p}) + 2\bar{\omega}(\bar{q}Y\bar{p} - \bar{p}Y\bar{q}) + \bar{\omega}^2 = - \frac{2\alpha}{W^2}.$$

Proof. We can use the global defining function $\phi(x, y, t) = x - y(\alpha t + \beta)$. As previously stipulated, we assume that \mathcal{S} is oriented in such a way that $\mathbf{N} = \nabla\phi = X_1\phi X_1 + X_2\phi X_2 + T\phi T$. Recalling (2.5), simple calculations based on (2.2) thus give

$$(3.17) \quad p = X_1\phi = 1 + \frac{\alpha}{2}y^2, \quad q = X_2\phi = -\alpha t - \beta - \frac{\alpha}{2}xy, \quad \omega = T\phi = -\alpha y.$$

The second equation in (3.17) becomes on \mathcal{S}

$$(3.18) \quad q = -(\alpha t + \beta) \left(1 + \frac{\alpha}{2}y^2\right).$$

We thus find on \mathcal{S}

$$(3.19) \quad W^2 = |\nabla^H\phi|^2 = \left(1 + \frac{\alpha}{2}y^2\right)^2 (1 + (\alpha t + \beta)^2).$$

Using (2.15) we obtain

$$(3.20) \quad Z(T\phi) = \bar{q}X_1(T\phi) - \bar{p}X_2(T\phi) = \frac{\alpha p}{W} = \frac{\alpha(1 + \frac{\alpha}{2}y^2)}{W} > 0.$$

Next, we have

$$(3.21) \quad X_1p = 0, \quad X_1q = 0, \quad X_2p = \alpha y, \quad X_2q = -\alpha x.$$

This gives

$$(3.22) \quad X_1W = \frac{pX_1p + qX_1q}{W} = 0, \quad X_2W = \frac{pX_2p + qX_2q}{W} = \alpha y (1 + (\alpha t + \beta)^2)^{1/2}.$$

From (2.15) and (3.22) we find

$$(3.23) \quad ZW = \bar{q}X_1W - \bar{p}X_2W = -\frac{p}{W}X_2W = -\alpha y.$$

Combining (3.23) with (3.20) we obtain

$$(3.24) \quad Z\bar{\omega} = \frac{Z(T\phi)}{W} - \frac{T\phi}{W^2}ZW = \frac{\alpha - \frac{\alpha^2}{2}y^2}{W^2}.$$

Using the above formulas it is not difficult to verify that

$$(3.25) \quad Y\bar{p} = Y\bar{q} = 0.$$

We now have from (3.17)

$$TW = \frac{pTp + qTq}{W} = -\frac{\alpha q}{W},$$

and therefore we easily find

$$(3.26) \quad \begin{cases} T\bar{p} = -\frac{\alpha(\alpha t + \beta)}{(1 + \frac{\alpha}{2}y^2)(1 + (\alpha t + \beta)^2)^{3/2}} = -\frac{\alpha(\alpha t + \beta)}{W(1 + (\alpha t + \beta)^2)}, \\ T\bar{q} = -\frac{\alpha}{W(1 + (\alpha t + \beta)^2)}. \end{cases}$$

From (3.17) and (3.26) we conclude that

$$(3.27) \quad \bar{p} T\bar{q} + \bar{q} T\bar{p} = \frac{\alpha(1 + \frac{\alpha}{2}y^2)((\alpha t + \beta)^2 - 1)}{W^2(1 + (\alpha t + \beta)^2)}.$$

From (3.18), (3.24), and (3.27), and elementary computations, we easily reach the conclusion that (3.15) holds. In a similar fashion, we obtain the proof of (3.16) by (3.17), (3.18), (3.19) and (3.26). \square

From Lemmas 3.7, 3.9 and 3.10, we obtain the following corollary.

Corollary 3.11. *Let \mathcal{S} be the H -minimal surface given by (3.14). For any $a \in C_0^\infty(\mathcal{S})$, the second variation along the deformation $\mathcal{S} \rightarrow \mathcal{S}^\lambda = \mathcal{S} + \lambda a X_1$ is given by*

$$(3.28) \quad \mathcal{V}_{II}^H(\mathcal{S}; aX_1) = \int_{\mathcal{S}} \frac{(1 + \frac{\alpha}{2}y^2)^2}{W^2} |\nabla^{H,\mathcal{S}} a|^2 d\sigma_H - 2\alpha \int_{\mathcal{S}} \frac{a^2}{W^2(1 + (\alpha t + \beta)^2)} d\sigma_H.$$

For any $h \in C_0^\infty(\mathcal{S})$, the second variation along the deformation $\mathcal{S} \rightarrow \mathcal{S}^\lambda = \mathcal{S} + \lambda h \nu_H$ is given by

$$(3.29) \quad \mathcal{V}_{II}^H(\mathcal{S}; h\nu_H) = \int_{\mathcal{S}} |\nabla^{H,\mathcal{S}} h|^2 d\sigma_H - 2\alpha \int_{\mathcal{S}} \frac{h^2}{W^2} d\sigma_H.$$

We now consider the global smooth parametrization $\theta : \mathbb{R}^2 \rightarrow \mathbb{R}^3$ of the surface \mathcal{S} given by $\theta(y, t) = (y(\alpha t + \beta), y, t)$. Clearly, $\mathcal{S} = \theta(\mathbb{R}^2)$.

Lemma 3.12. *Let \mathcal{S} be the H -minimal surface given by (3.14). For any $a \in C_0^\infty(\mathcal{S})$, then one has*

$$(3.30) \quad \begin{aligned} \mathcal{V}_{II}^H(\mathcal{S}; aX_1) &= \int_{\mathbb{R}^2} \frac{(1 + \frac{\alpha}{2}y^2) u_y^2}{(1 + (\alpha t + \beta)^2)^{3/2}} dydt \\ &\quad - 2\alpha \int_{\mathbb{R}^2} \frac{u^2}{(1 + \frac{\alpha}{2}y^2)(1 + (\alpha t + \beta)^2)^{3/2}} dydt, \end{aligned}$$

where $u = a \circ \theta \in C_0^\infty(\mathbb{R}^2)$. For any $h \in C_0^\infty(\mathcal{S})$, the one has

$$(3.31) \quad \begin{aligned} \mathcal{V}_{II}^H(\mathcal{S}; h\nu_H) &= \int_{\mathbb{R}^2} \frac{(1 + \frac{\alpha}{2}y^2) u_y^2}{(1 + (\alpha t + \beta)^2)^{1/2}} dydt \\ &\quad - 2\alpha \int_{\mathbb{R}^2} \frac{u^2}{(1 + \frac{\alpha}{2}y^2)(1 + (\alpha t + \beta)^2)^{1/2}} dydt, \end{aligned}$$

where this time we have set $u = h \circ \theta \in C_0^\infty(\mathbb{R}^2)$.

Proof. In order to prove (3.30) we make some reductions. Keeping in mind that from (2.23) we have $d\sigma_H = (|\nabla^H \phi|/|\nabla \phi|)d\sigma = (W/|\nabla \phi|)d\sigma$, from (3.19) we obtain

$$\int_{\mathcal{S}} \frac{a^2}{W^2(1+(\alpha t + \beta)^2)} d\sigma_H = \int_{\mathbb{R}^2} \frac{u^2}{\left(1 + \frac{\alpha}{2}y^2\right)(1+(\alpha t + \beta)^2)^{3/2}} dy dt .$$

In order to express the first integral in the right-hand side of (3.28) as an integral on \mathbb{R}^2 , we compute $|\nabla^{H,\mathcal{S}} a|^2$. We have from (2.9), (2.15), (3.17) and (3.18)

$$(3.32) \quad \begin{aligned} |\nabla^{H,\mathcal{S}} a|^2 &= (Za)^2 = (\bar{q}X_1a - \bar{p}X_2a)^2 \\ &= \frac{((\alpha t + \beta)X_1a + X_2a)^2}{1 + (\alpha t + \beta)^2} . \end{aligned}$$

Now, the chain rule gives $u_y = (\alpha t + \beta)a_x + a_y$, and therefore we see from (2.2) that we have on \mathcal{S}

$$(\alpha t + \beta)X_1a + X_2a = (\alpha t + \beta) a_x + a_y = u_y .$$

From (3.32) we thus conclude that

$$\int_{\mathcal{S}} \bar{p}^2 |\delta_H a|^2 d\sigma_H = \int_{\mathbb{R}^2} \frac{\left(1 + \frac{\alpha}{2}y^2\right) u_y^2}{\left(1 + (\alpha t + \beta)^2\right)^{3/2}} dy dt .$$

This proves (3.30). The proof of (3.31) proceeds analogously, and we omit the details. \square

Lemma 3.13. *Let $\chi_k \in C_0^\infty(\mathbb{R})$ be such that $0 \leq \chi_k(s) \leq 1$, $\chi_k(s) = 0$ for $|s| > 2k$, $\chi_k(s) \equiv 1$ for $|s| < k$, and $|\chi_k'(s)| \leq C/k$ with C independent of k . Define for any $\alpha > 0$*

$$f_k(y) = \frac{\chi_k(y)}{\sqrt{1 + \frac{\alpha}{2}y^2}} .$$

We have $f_k \in C_0^\infty(\mathbb{R})$, and there exists $k_0 \in \mathbb{N}$ such that for all $k > k_0$

$$(3.33) \quad \int_{\mathbb{R}} \frac{f_k(y)^2}{1 + \frac{\alpha}{2}y^2} dy > \frac{1}{2\alpha} \int_{\mathbb{R}} \left(1 + \frac{\alpha}{2}y^2\right) f_k'(y)^2 dy .$$

Proof. We begin by observing that integration of the right-hand side of (3.33) gives

$$(3.34) \quad \begin{aligned} &\int_{\mathbb{R}} \left(1 + \frac{\alpha}{2}y^2\right) f_k'(y)^2 dy \\ &= \int_{\mathbb{R}} (\chi_k')^2 dy + \frac{\alpha^2}{4} \int_{\mathbb{R}} \frac{y^2 \chi_k^2}{\left(1 + \frac{\alpha}{2}y^2\right)^2} dy + \frac{\alpha}{2} \int_{\mathbb{R}} \chi_k^2 \left(\frac{y}{1 + \frac{\alpha}{2}y^2}\right)' dy . \end{aligned}$$

Observing that

$$\left(\frac{y}{1 + \frac{\alpha}{2}y^2}\right)' = \frac{1 - \frac{\alpha}{2}y^2}{\left(1 + \frac{\alpha}{2}y^2\right)^2} ,$$

we conclude from (3.34) and dominated convergence

$$(3.35) \quad \begin{aligned} &\int_{\mathbb{R}} \left(1 + \frac{\alpha}{2}y^2\right) f_k'(y)^2 dy = O\left(\frac{1}{k}\right) + \frac{\alpha}{2} \int_{\mathbb{R}} \frac{\chi_k^2}{\left(1 + \frac{\alpha}{2}y^2\right)^2} dy \\ &\longrightarrow \frac{\alpha}{2} \int_{\mathbb{R}} \frac{1}{\left(1 + \frac{\alpha}{2}y^2\right)^2} dy = \sqrt{\frac{\alpha}{2}} \frac{\pi}{2} , \quad \text{as } k \rightarrow \infty . \end{aligned}$$

On the other hand, again by dominated convergence, we obtain for the integral in the left-hand side of (3.33)

$$(3.36) \quad \int_{\mathbb{R}} \frac{f_k(y)^2}{1 + \frac{\alpha}{2}y^2} dy = \int_{\mathbb{R}} \frac{\chi_k(y)^2}{(1 + \frac{\alpha}{2}y^2)^2} dy \\ \longrightarrow \int_{\mathbb{R}} \frac{1}{(1 + \frac{\alpha}{2}y^2)^2} dy = \frac{\pi}{2} \sqrt{\frac{2}{\alpha}} \quad \text{as } k \rightarrow \infty .$$

In view of (3.35), (3.36), we obtain

$$\int_{\mathbb{R}} \frac{f_k(y)^2}{1 + \frac{\alpha}{2}y^2} dy - \frac{1}{2\alpha} \int_{\mathbb{R}} \left(1 + \frac{\alpha}{2}y^2\right) f'_k(y)^2 dy \longrightarrow \frac{3\pi}{8} \sqrt{\frac{2}{\alpha}} .$$

From this the conclusion readily follows. □

As an immediate consequence of Lemma 3.13 and Fubini's theorem we obtain.

Corollary 3.14. *For $k \in \mathbb{N}$, $k \geq k_0$, define $u_k(y, t) = f_k(y)\chi_k(t)$, where χ_k and f_k and k_0 are as in Lemma 3.13. One has*

$$(3.37) \quad \int_{\mathbb{R}^2} \frac{u_k(y, t)^2}{(1 + \frac{\alpha}{2}y^2)(1 + (\alpha t + \beta)^2)^{3/2}} dy dt > \frac{1}{2\alpha} \int_{\mathbb{R}^2} \frac{1 + \frac{\alpha}{2}y^2}{(1 + (\alpha t + \beta)^2)^{3/2}} \left(\frac{\partial u_k(y, t)}{\partial y}\right)^2 dy dt .$$

We are finally ready to give the

Proof of Theorem 1.2. Let u_k , $k \geq k_0$, be as in Corollary 3.14. Define $a_k : \mathbb{H}^1 \rightarrow \mathbb{R}$, $a_k \in C_0^\infty(\mathbb{H}^1)$, as follows

$$a_k(x, y, t) = \frac{\chi_k(y)\chi_k(t)\chi_k(x - y(\alpha t + \beta))}{\sqrt{1 + \frac{\alpha}{2}y^2}} .$$

We observe that $a_k(\theta(y, t)) = u_k(y, t)\chi_k(0) = u_k(y, t)$. At this point, appealing to (3.30) in Lemma 3.12 and to Corollary 3.14, we conclude that for every fixed $k \geq k_0$, we have for the deformation $\mathcal{S} \rightarrow \mathcal{S}^\lambda = \mathcal{S} + \lambda a_k X_1$

$$\mathcal{V}_{II}^H(\mathcal{S}; a_k X_1) = \frac{d^2}{d\lambda^2} P_H(\mathcal{S}^\lambda) \Big|_{\lambda=0} < 0 .$$

This proves that \mathcal{S} cannot be a local minimizer of the H -perimeter for compactly supported deformations along X_1 . In a similar way, using (3.30) in Lemma 3.12 and Corollary 3.14, we see that \mathcal{S} cannot be a local minimizer for deformations along the horizontal normal ν_H . In particular, since every global minimizer is also a local one, \mathcal{S} cannot be a global minimizer either. □

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