# PRECISE LOCAL ESTIMATES FOR DIFFERENTIAL EQUATIONS DRIVEN BY FRACTIONAL BROWNIAN MOTION: HYPOELLIPTIC CASE 

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#### Abstract

This article is concerned with stochastic differential equations driven by a $d$ dimensional fractional Brownian motion with Hurst parameter $H>1 / 4$ and understood in the rough paths sense. Whenever the coefficients of the equation satisfy a uniform hypoellipticity condition, we establish a sharp local estimate on the associated control distance function and a sharp local lower estimate on the density of the solution. Our methodology relies heavily on the rough paths structure of the equation.


1. Introduction. We will split our Introduction into two parts. In Section 1.1 we recall some background on the stochastic analysis of stochastic differential equations driven by a fractional Brownian motion. In Section 1.2 we describe our main results. Section 1.3 is then devoted to a brief explanation about the methodology we have used in order to obtain our main results.
1.1. Background and motivation. One way to envision Malliavin calculus is to see it as a geometric and analytic framework on an infinite dimensional space (namely, the Wiener space) equipped with a Gaussian measure. This is already apparent in Malliavin's seminal contribution [23] giving a probabilistic proof of Hörmander's theorem. The same point of view has then been pushed forward in the celebrated series of papers by Kusuoka and Stroock [18-20] which set up the basis for densities and probabilities expansions for diffusion processes within this framework.

On the other hand, the original perspective of Lyons' rough path theory (cf. [21, 22]) is quite different. Summarizing very briefly, it asserts that a reasonable differential calculus with respect to a noisy process $X$ can be achieved as long as one can define enough iterated integrals of $X$. One of the first processes to which the theory has been successfully applied is a fractional Brownian motion $B$ with Hurst parameter $H$ on which we shall focus in the present paper. The process $B$ can be seen as a natural generalization of Brownian motion allowing any kind of Hölder regularity. We are interested in the following differential equation driven by $B$ :

$$
\left\{\begin{array}{l}
d X_{t}=\sum_{\alpha=1}^{d} V_{\alpha}\left(X_{t}\right) d B_{t}^{\alpha}, \quad 0 \leq t \leq 1,  \tag{1.1}\\
X_{0}=x \in \mathbb{R}^{N}
\end{array}\right.
$$

Here, the $V_{\alpha}$ 's are $C_{b}^{\infty}$ vector fields, and the Hurst parameter is assumed to satisfy $H>1 / 4$. In this setting, putting together the results contained in [9] and [22], the stochastic differential equation (1.1) can be understood in the framework of rough path theory.

[^0]With the solution of (1.1) in hand, a natural problem one can think of is the following: can we extend the aforementioned analytic studies on Wiener's space to the process $B$ ? In particular, can we complete Kusuoka-Stroock's program in the fractional Brownian motion setting? This question has received a lot of attention in the recent years, and previous efforts along this line include Hörmander-type theorems for the process $X$, defined by (1.1) (cf. [2, $7,8]$ ), and some upper Gaussian bounds on the density $p(t, x, y)$ of $X_{t}(c f .[3,16])$ as well as $\operatorname{Varadhan-type~estimates~for~} \log (p(t, x, y))$ in small time [4]. One should stress at this point that the road from the Brownian to the fractional Brownian case is far from being trivial. This is essentially due to the lack of independence of the fBm increments and Markov property as well as to the technically demanding characterization of the Cameron-Martin space whenever $B$ is not a Brownian motion. We shall go back to those obstacles throughout the article.

Our contribution can be seen as a step in the direction mentioned above. More specifically, we shall obtain a sharp local estimate on the associated control distance function and some sharp local bounds for the density of $X_{t}$ under hypoelliptic conditions on the vector fields $V_{\alpha}$. This will be achieved thanks to a combination of geometric and analytic tools which can also be understood as a mix of stochastic analysis and rough path theory. We describe our main results more precisely in the next subsection.
1.2. Statement of main results. Let us recall that equation (1.1) is our main object of concern. We are typically interested in the degenerate case where the vector fields $V=$ $\left\{V_{1}, \ldots, V_{d}\right\}$ satisfy the uniform hypoellipticity assumption to be defined shortly. This is a standard degenerate setting where one can expect that the solution of the SDE (1.1) admits a smooth density with respect to the Lebesgue measure. As mentioned in Section 1.1, we wish to obtain quantitative information for the density in this context.

We first formulate the uniform hypoellipticity condition which will be assumed throughout the rest of the paper. For $l \geq 1$, define $\mathcal{A}(l)$ to be the set of words over letters $\{1, \ldots, d\}$ with length, at most, $l$ (including the empty word) and $\mathcal{A}_{1}(l) \triangleq \mathcal{A}(l) \backslash\{\varnothing\}$. Denote $\mathcal{A}_{1}$ as the set of all nonempty words. Given a word $\alpha \in \mathcal{A}_{1}$, we define the vector field $V_{[\alpha]}$ inductively by $V_{[i]} \triangleq V_{i}$ and $V_{[\alpha]} \triangleq\left[V_{i}, V_{[\beta]}\right]$ for $\alpha=(i, \beta)$ with $i$ being a letter and $\beta \in \mathcal{A}_{1}$.

Uniform Hypoellipticity Assumption. The vector fields $\left(V_{1}, \ldots, V_{d}\right)$ are $C_{b}^{\infty}$, and there exists an integer $l_{0} \geq 1$ such that

$$
\begin{equation*}
\inf _{x \in \mathbb{R}^{N}} \inf _{\eta \in S^{N-1}}\left\{\sum_{\alpha \in \mathcal{A}_{1}\left(l_{0}\right)}\left\langle V_{[\alpha]}(x), \eta\right\rangle_{\mathbb{R}^{N}}^{2}\right\}>0 . \tag{1.2}
\end{equation*}
$$

The smallest such $l_{0}$ is called the hypoellipticity constant for the vector fields.
Under condition (1.2) it was proved by Cass-Friz [7] and Cass-Hairer-Litterer-Tindel [8] that the solution to the SDE (1.1) admits a smooth density $y \mapsto p(t, x, y)$ with respect to the Lebesgue measure on $\mathbb{R}^{N}$ for all $(t, x) \in(0,1] \times \mathbb{R}^{N}$. Our contribution aims at getting quantitative small time estimates for $p(t, x, y)$.

In order to describe our bounds on the density $p(t, x, y)$, let us recall that the small time behavior of $p(t, x, y)$ is closely related to the so-called control distance function associated with the vector fields. This fact was already revealed in the Varadhan-type asymptotics result proved by Baudoin-Ouyang-Zhang [4],

$$
\begin{equation*}
\lim _{t \rightarrow 0} t^{2 H} \log p(t, x, y)=-\frac{1}{2} d(x, y)^{2} \tag{1.3}
\end{equation*}
$$

The control distance function $d(x, y)$ in (1.3), which plays a prominent role in our paper, is defined as follows.

Definition 1.1. Let $\overline{\mathcal{H}}$ be the Cameron-Martin space of the fractional Brownian motion. For any $h \in \overline{\mathcal{H}}$, denote by $\Phi_{t}(x ; h)$ the solution to the ODE

$$
\begin{equation*}
d x_{t}=\sum_{\alpha=1}^{d} V_{\alpha}\left(x_{t}\right) d h_{t}^{\alpha}, \quad 0 \leq t \leq 1, \text { with } x_{0}=x \tag{1.4}
\end{equation*}
$$

In addition, for $x, y \in \mathbb{R}^{N}$ and $\Phi_{t}(x ; h)$, defined as in (1.4), set

$$
\begin{equation*}
\Pi_{x, y} \triangleq\left\{h \in \overline{\mathcal{H}}: \Phi_{1}(x ; h)=y\right\} \tag{1.5}
\end{equation*}
$$

to be the space of Cameron-Martin paths which join $x$ to $y$ in the sense of differential equations. The control distance function $d(x, y)=d_{H}(x, y)$ is defined by

$$
d(x, y) \triangleq \inf \left\{\|h\|_{\overline{\mathcal{H}}}: h \in \Pi_{x, y}\right\}, \quad x, y \in \mathbb{R}^{N}
$$

According to (1.3), one can clearly expect that the Cameron-Martin structure and the control distance function will play an important role in understanding the small time behavior of $p(t, x, y)$. However, unlike the diffusion case and due to the complexity of the CameronMartin structure of a fractional Brownian motion, the function $d(x, y)$ is far from being a metric, and its shape is not clear. Our first main result is thus concerned with the local behavior of $d(x, y)$. It establishes a comparison between $d$ and the Euclidian distance.

THEOREM 1.2. Suppose that the vector fields $\left\{V_{1}, \ldots, V_{d}\right\}$ satisfy the uniform hypoellipticity condition (1.2) with hypoellipticity constant $l_{0}$. Let d be the control distance function given in Definition 1.1. There exist constants $C_{1}, C_{2}, \delta>0$, where $C_{1}, C_{2}$ depend only on $H$, $l_{0}$ and the vector fields and where $\delta$ depends only on $l_{0}$ and the vector fields such that

$$
\begin{equation*}
C_{1}|x-y| \leq d(x, y) \leq C_{2}|x-y|^{\frac{1}{I_{0}}} \tag{1.6}
\end{equation*}
$$

for all $x, y \in \mathbb{R}^{N}$ with $|x-y|<\delta$.
With the technique we use to prove Theorem 1.2 together with some further effort, we are, in fact, able to establish a stronger result, namely, the local equivalence of $d$ to the sub-Riemannian distance induced by the vector fields $\left\{V_{1}, \ldots, V_{d}\right\}$. More specifically, let us write the distance given in Definition 1.1 as $d_{H}(x, y)$ in order to emphasize the dependence on the Hurst parameter $H$. Our second main result asserts that all the distances $d_{H}$ are locally equivalent.

THEOREM 1.3. Under the same assumptions as in Theorem 1.2, for any $H_{1}, H_{2} \in$ $(1 / 4,1)$, there exist constants $C=C\left(H_{1}, H_{2}, V\right)>0$ and $\delta>0$ such that

$$
\begin{equation*}
\frac{1}{C} d_{H_{1}}(x, y) \leq d_{H_{2}}(x, y) \leq C d_{H_{1}}(x, y) \tag{1.7}
\end{equation*}
$$

for all $x, y \in \mathbb{R}^{N}$ with $|x-y|<\delta$. In particular, all distances $d_{H}$ are locally equivalent to $d_{\mathrm{BM}} \equiv d_{1 / 2}$, where $d_{\mathrm{BM}}$ stands for the control distance function of the system (1.4) driven by a Brownian motion, that is, the sub-Riemannian distance induced by the vector fields $\left\{V_{1}, \ldots, V_{d}\right\}$.

The local equivalence of distances, stated in Theorem 1.3, plays a crucial role in our later analysis of the density $p(t, x, y)$, in particular, in the proof of Theorem 1.6 below (e.g., see the proof of Lemma 4.15 below). Moreover, we believe Theorem 1.3 and the tools we developed in order to prove this theorem may be of independent interest and have other applications as well.

REMARK 1.4. In the special case when (1.1) reads as $d X_{t}=X_{t} \otimes d B_{t}$, that is, when $X_{t}$ is the truncated signature of $B$ up to order $l>0$, it is proved in [1] that all $d_{H}(x, y)$ are globally equivalent. The proof crucially depends on the fact that the signature of $B$ is homogeneous with respect to the dilation operator on $G^{(l)}\left(\mathbb{R}^{d}\right)$, the free nilpotent Lip group over $\mathbb{R}^{d}$ of order $l$. In the current general nonlinear case, the local equivalence is much more technically challenging. In addition, we think that the global equivalence of the distances $d_{H}$ does not hold.

Our third main result asserts that the density $p(t, x, y)$ of $X_{t}$ is strictly positive everywhere whenever $t>0$. It generalizes, for the first time, the result of [3], Theorem 1.4, to a general hypoelliptic case by affirming that Hypothesis 1.2 in that theorem is always verified under our assumption (1.2). Recall that a distribution over a differentiable manifold is a smooth choice of subspace of the tangent space at every point with constant dimension.

THEOREM 1.5. Let $\left\{V_{1}, \ldots, V_{d}\right\}$ be a family of $C_{b}^{\infty}$-vector fields on $\mathbb{R}^{N}$ which span a distribution over $\mathbb{R}^{N}$ and satisfy the uniform hypoellipticity assumption (1.2). Let $X_{t}^{x}$ be the solution to the stochastic differential equation (1.1), where $B_{t}$ is a d-dimensional fractional Brownian motion with Hurst parameter $H>1 / 4$. Then, for each $t \in(0,1]$, the density of $X_{t}$ is everywhere strictly positive.

As we will see in Section 4.1, the proof of the above result is based on finite dimensional geometric arguments, such as the classical Sard theorem as well as a general positivity criteria for densities on the Wiener space. After our preprint was released, further results on strict positivity were obtained very recently in [17] using the notion of $\mathcal{K}$-regularity. However, the latter reference focuses exclusively on the strict positivity of density for Gaussian rough differential equations. In contrast, the proof of Theorem 1.5 comes as a natural and short byproduct of our global approach.

Let us now turn to a description of our last main result. It establishes a sharp local lower estimate for the density function $p(t, x, y)$ of the solution to the $\operatorname{SDE~(1.1)~in~small~time.~}$

THEOREM 1.6. Under the uniform hypoellipticity assumption (1.2), let $p(t, x, y)$ be the density of the random variable $X_{t}$ defined by equation (1.1). There exist some constants $C, \tau>0$, depending only on $H, l_{0}$ and the vector fields $V_{\alpha}$, such that

$$
\begin{equation*}
p(t, x, y) \geq \frac{C}{\left|B_{d}\left(x, t^{H}\right)\right|}, \tag{1.8}
\end{equation*}
$$

for all $(t, x, y) \in(0,1] \times \mathbb{R}^{N} \times \mathbb{R}^{N}$, satisfying the following local condition involving the distance d introduced in Definition 1.1,

$$
d(x, y) \leq t^{H} \quad \text { and } \quad t<\tau
$$

In relation (1.8), $B_{d}\left(x, t^{H}\right) \triangleq\left\{z \in \mathbb{R}^{N}: d(x, z)<t^{H}\right\}$ denotes the ball with respect to the distance $d$, and $|\cdot|$ stands for the Lebesgue measure.

The sharpness of Theorem 1.6 can be seen from the fractional Brownian motion case, that is, when $N=d$ and $V=\mathrm{Id}$. We also point out that, in terms of uniform local lower estimates, the cone $d(x, y) \leq t^{H}$ is the regime where the estimate appears to be interesting, as it controls the rate of explosion (singularity) of the density as $t \rightarrow 0^{+}$. Finally, let us stress again that the developments of the above main results are intrinsically connected. Indeed, as we will see, the technique we use to prove Theorem 1.2 will be an essential ingredient for establishing Theorem 1.3. In addition, Theorem 1.3 and Theorem 1.5 provide two essential ingredients toward the proof of Theorem 1.6.
1.3. Strategy and outlook. Let us say a few words about the methodology we have used in order to obtain our main results. Although we will describe our overall strategy with more details in Section 4, let us mention here that it is based on the reduction of the problem to a finite dimensional one, plus some geometric-type arguments. Let us also highlight the fact that a much simpler strategy can be used for the elliptic case, as explained in our companion paper [15].

More specifically, the key point in our proofs is that the solution $X_{t}$ to (1.1) can be approximated by a simple enough function of the so-called truncated signature of order $l$ for the fractional Brownian motion $B$. This object is formally defined, for a given $l \geq 1$, as the following $\bigoplus_{k=0}^{l}\left(\mathbb{R}^{d}\right)^{\otimes k}$-valued process:

$$
\Gamma_{t}=1+\sum_{k=1}^{l} \int_{0<t_{1}<\cdots<t_{k}<t} d B_{t_{1}} \otimes \cdots \otimes d B_{t_{k}}
$$

and it enjoys some convenient algebraic and analytic properties. The truncated signature is the main building block of the rough path theory (see, e.g., [22]) and was also used in [20] in a Malliavin calculus context. Part of our challenge in the current contribution is to combine the properties of the process $\Gamma$, together with the Cameron-Martin space structure related to the fractional Brownian motion $B$, in order to achieve sharp estimates for the control distance function as well as the density of solution.

As mentioned above, the truncated signature gives rise to a $l$ th order local approximation of $X_{t}$ in a neighborhood of its initial condition $x$. Namely, if $V_{\left(i_{1}, \ldots, i_{k}\right)}$ denotes the composition of $V_{i_{1}}$ up to $V_{i_{k}}$ considered as differential operators and if we set

$$
\begin{equation*}
F_{l}\left(\Gamma_{t}, x\right) \triangleq \sum_{k=1}^{l} \sum_{i_{1}, \ldots, i_{k}=1}^{d} V_{\left(i_{1}, \ldots, i_{k}\right)}(x) \int_{0<t_{1}<\cdots<t_{k}<t} d B_{t_{1}}^{i_{1}} \cdots d B_{t_{k}}^{i_{k}} \tag{1.9}
\end{equation*}
$$

then classical rough paths considerations assert that $F_{l}\left(\Gamma_{t}, x\right)$ is an approximation of order $t^{H l}$ of $X_{t}$ for small $t$. In the sequel we will heavily rely on some degeneracy properties of $F_{l}$ derived from the uniform hypoelliptic assumption (1.2) in order to get the following information:
(i) One can construct a path $h$ in the Cameron-Martin space of $B$ which joins $x$ and any point $y$ in a small enough neighborhood of $x$. This task is carried out thanks to a complex iteration procedure, whose building block is the nondegeneracy of the function $F_{l}$. It is detailed in Section 3.2. In this context, observe that the computation of the Cameron-Martin norm of $h$ also requires a substantial effort. This will be the key step in order to prove Theorem 1.2 and Theorem 1.3 concerning the distance $d$ given in Definition 1.1.
(ii) The proof of the lower bound, given in Theorem 1.6, also hinges heavily on the approximation $F_{l}$ given by (1.9). Indeed, the preliminary results about the density of $\Gamma_{t}$, combined with the nondegeneracy of $F_{l}$, yield good properties for the density of $F_{l}\left(\Gamma_{t}, x\right)$. One is then left with the task of showing that $F_{l}\left(\Gamma_{t}, x\right)$ approximates $X_{t}$ properly at the density level.

In conclusion, although the steps performed in the remainder of the article might look technically and computationally involved, they rely on a natural combination of analytic and geometric bricks as well as a reduction to a finite dimensional problem. Let us also highlight the fact that our next challenge is to iterate the local estimates presented here in order to get Gaussian-type lower bounds for the density $p(t, x, y)$ of $X_{t}$. This causes some further complications, due to the complex (non-Markovian) dependence structure for the increments of the fractional Brownian motion $B$. We defer this important problem to a future project. Eventually, we mention that this note is an abridged and self-contained version of our original draft on the same topic. For further computational details, the reader is referred to [14].

REMARK 1.7. As the reader might have observed, our equation (1.1) does not have a drift component. If one wishes to add a drift $V_{0}\left(X_{t}\right) d t$ to the equation, it first means that the approximation function $F_{l}$ in (1.9) should include mixed integrals involving both $d t$ and $d B_{t}$ differentials. This could possibly be achieved thanks to tree-type notation, similar to what is done, for example, in [26]. Once the approximation map $F_{l}$ is constructed, some subtle effects coming from the drift term are observed in the hypoelliptic case and should be taken into account in our density estimates. We refer to [5] for an analysis of the small noise asymptotics for the probability of sets $A$ which are nonhorizontally accessible, meaning that the drift $V_{0}$ is needed to reach those sets. Our density bounds should go along the same lines in the hypoelliptic case with drift. As in Kusuoka and Stroock [20], we have refrained to incorporate those elaborate developments in the current paper for the sake of conciseness and simplicity. However, they are certainly worth considering for a future contribution.

Organization of the present paper. In Section 2 we present some basic notions from the analysis of fractional Brownian motion and rough path theory. In Section 3 we develop the proofs of Theorem 1.2 and Theorem 1.3 concerning the control distance function. In Section 4 we develop the proof of Theorem 1.6 concerning the density of solution. Theorem 1.5 is proved in the first key step toward proving Theorem 1.6 in Section 4.1.

Notation. Throughout the rest of this paper, we will use "Letter ${ }_{\text {subscript" }}$ to denote constants whose value depend only on the objects specified in the "subscript" and may differ from line to line. Unless otherwise stated, a constant will implicitly depend on $H, V, l_{0}$. We will always omit the dependence on dimension.
2. Preliminary results. This section is devoted to some preliminary notions on the stochastic analysis of fractional Brownian motion. We also recall some basic facts about rough path solutions to noisy differential equations.
2.1. Stochastic analysis of fractional Brownian motion. Let us start by recalling the definition of fractional Brownian motion.

Definition 2.1. A d-dimensional fractional Brownian motion with Hurst parameter $H \in(0,1)$ is an $\mathbb{R}^{d}$-valued continuous centered Gaussian process $B_{t}=\left(B_{t}^{1}, \ldots, B_{t}^{d}\right)$ whose covariance structure is given by

$$
\begin{equation*}
\mathbb{E}\left[B_{s}^{i} B_{t}^{j}\right]=\frac{1}{2}\left(s^{2 H}+t^{2 H}-|s-t|^{2 H}\right) \delta_{i j} \triangleq R(s, t) \delta_{i j} \tag{2.1}
\end{equation*}
$$

This process is defined and analyzed in numerous articles (cf. [11, 27, 28], for instance) to which we refer for further details. We always assume that the fractional Brownian motion $B$ is defined on a complete probability space $(\Omega, \mathcal{F}, \mathbb{P})$ where $\mathcal{F}$ is generated by $\left\{B_{t}: t \in[0, T]\right\}$. In this section we focus on basic stochastic analysis notions which will be used in the sequel.

In order to introduce the Hilbert spaces which will feature in the sequel, we consider a onedimensional fractional Brownian motion $\left\{B_{t}: 0 \leq t \leq 1\right\}$ with Hurst parameter $H \in(0,1)$ defined on $(\Omega, \mathcal{F}, \mathbb{P})$. The discussion here can be easily adapted to the multidimensional setting with arbitrary time horizon $[0, T]$. Let $\mathcal{C}_{1}$ be the associated first order Wiener chaos, that is, $\mathcal{C}_{1} \triangleq \overline{\operatorname{Span}\left\{B_{t}: 0 \leq t \leq 1\right\}}$ in $L^{2}(\Omega, \mathbb{P})$. We will frequently identify a Hilbert space with its dual in the canonical way without further mentioning.

DEFINITION 2.2. Let $W$ be the space of continuous paths $w:[0,1] \rightarrow \mathbb{R}^{1}$ with $w_{0}=0$. Define $\overline{\mathcal{H}}$ to be the space of elements $h \in W$ that can be written as

$$
\begin{equation*}
h_{t}=\mathbb{E}\left[B_{t} Z\right], \quad 0 \leq t \leq 1, \tag{2.2}
\end{equation*}
$$

where $Z \in \mathcal{C}_{1}$. We equip $\overline{\mathcal{H}}$ with an inner product structure given by

$$
\left\langle h, h^{\prime}\right\rangle_{\overline{\mathcal{H}}} \triangleq \mathbb{E}\left[Z \cdot Z^{\prime}\right], \quad h, h^{\prime} \in \overline{\mathcal{H}}
$$

whenever $h, h^{\prime}$ are defined by (2.2) for two random variables $Z, Z^{\prime} \in \mathcal{C}_{1}$. The Hilbert space $\left(\overline{\mathcal{H}},\langle\cdot, \cdot\rangle_{\overline{\mathcal{H}}}\right)$ is called the Cameron-Martin subspace of the fractional Brownian motion.

One of the advantages of working with fractional Brownian motion is that a convenient analytic description of $\overline{\mathcal{H}}$ in terms of fractional calculus is available. We refer to the aforementioned references [11, 27, 28] for the definition of fractional derivatives and some further characterizations of $\overline{\mathcal{H}}$. We mention one useful fact here: there is an isomorphism $K$ between $L^{2}([0,1])$ and $I_{0+}^{H+1 / 2}\left(L^{2}([0,1])\right)$, defined by

$$
K \varphi \triangleq \begin{cases}C_{H} \cdot I_{0^{+}}^{1}\left(t^{H-\frac{1}{2}} \cdot I_{0^{+}}^{H-\frac{1}{2}}\left(s^{\frac{1}{2}-H} \varphi(s)\right)(t)\right), & H>\frac{1}{2}  \tag{2.3}\\ C_{H} \cdot I_{0^{+}}^{2 H}\left(t^{\frac{1}{2}-H} \cdot I_{0^{+}}^{\frac{1}{2}-H}\left(s^{H-\frac{1}{2}} \varphi(s)\right)(t)\right), & H \leq \frac{1}{2}\end{cases}
$$

where $C_{H}$ is a universal constant depending only on $H$. Then, the space $\overline{\mathcal{H}}$ can be identified with $I_{0^{+}}^{H+1 / 2}\left(L^{2}([0,1])\right)$, and the Cameron-Martin norm is given by

$$
\begin{equation*}
\|h\|_{\overline{\mathcal{H}}}=\left\|K^{-1} h\right\|_{L^{2}([0,1])} \tag{2.4}
\end{equation*}
$$

Next, we mention a variational embedding theorem for the Cameron-Martin subspace $\overline{\mathcal{H}}$ which will be used in a crucial way. The case when $H>1 / 2$ is a simple exercise, while the case when $H \leq 1 / 2$ was treated in [12]. From a pathwise viewpoint this allows us to integrate a fractional Brownian path against a Cameron-Martin path or vice versa (cf. [29]) and to make sense of ordinary differential equations driven by a Cameron-Martin path (cf. [21]).

Proposition 2.3. If $H>\frac{1}{2}$, then $\overline{\mathcal{H}} \subseteq C_{0}^{H}\left([0,1] ; \mathbb{R}^{d}\right)$, the space of H-Hölder continuous paths. If $H \leq \frac{1}{2}$, then for any $q>(H+1 / 2)^{-1}$, we have $\overline{\mathcal{H}} \subseteq C_{0}^{q-\mathrm{var}}\left([0,1] ; \mathbb{R}^{d}\right)$, the space of continuous paths with finite $q$-variation. In addition, the above inclusions are continuous embeddings.

For the sake of conciseness, we only recall a few notation on Malliavin calculus which feature prominently in our future considerations. The reader is referred to [27] for a thorough introduction to Malliavin calculus techniques. We use $\mathbf{D}_{h}$ to denote the Malliavin derivative along a direction $h \in \overline{\mathcal{H}}$. Respectively, higher order derivatives are denoted by $\mathbf{D}_{h_{1}, \ldots, h_{k}}^{k}$. The corresponding Sobolev spaces are written as $\mathbb{D}^{k, p}(\overline{\mathcal{H}})$, and we also set $\mathbb{D}^{\infty}(\overline{\mathcal{H}})=\bigcap_{p \geq 1} \bigcap_{k \geq 1} \mathbb{D}^{k, p}(\overline{\mathcal{H}})$.

DEFINITION 2.4. Let $F=\left(F^{1}, \ldots, F^{n}\right)$ be a random vector whose components are in $\mathbb{D}^{\infty}(\overline{\mathcal{H}})$. Define the Malliavin covariance matrix of $F$ by

$$
\begin{equation*}
\gamma_{F}=\left(\left\langle\mathbf{D} F^{i}, \mathbf{D} F^{j}\right\rangle_{\overline{\mathcal{H}}}\right)_{1 \leq i, j \leq n} \tag{2.5}
\end{equation*}
$$

Then, $F$ is said to be nondegenerate if $\gamma_{F}$ is invertible a.s. and

$$
\left(\operatorname{det} \gamma_{F}\right)^{-1} \in \bigcap_{p \geq 1} L^{p}(\Omega)
$$

It is a fundamental result in Malliavin calculus that the law of a nondegenerate random vector $F$ admits a smooth density with respect to the Lebesgue measure on $\mathbb{R}^{n}$ (cf. [27], Theorem 2.1.4).
2.2. Free nilpotent groups. Next, we introduce some basic algebraic structure that plays an essential role in the rough path analysis of equation (1.1). The reader is referred to [13] for a more systematic presentation. We start with the following notation.

Notation 2.5. The truncated tensor algebra of order $l$ over $\mathbb{R}^{d}$ is denoted by $T^{(l)}$. Under addition and tensor product, $\left(T^{(l)},+, \otimes\right)$ is an associative algebra. The set of homogeneous Lie polynomials of degree $k$ is denoted as $\mathcal{L}_{k}$, and the free nilpotent Lie algebra of order $l$ is denoted as $\mathfrak{g}^{(l)}$. The exponential map on $T^{(l)}$ is defined by

$$
\begin{equation*}
\exp (a) \triangleq \sum_{k=0}^{\infty} \frac{1}{k!} a^{\otimes k} \in T^{(l)} \tag{2.6}
\end{equation*}
$$

where the sum is indeed locally finite and hence well defined. $T^{(l)}$ is equipped with a natural inner product induced by the Euclidean structure on $\mathbb{R}^{d}$ and the Hilbert-Schmidt tensor norm on each of the tensor products. To be more precise, the inner product on $\left(\mathbb{R}^{d}\right)^{\otimes k}$ is induced by

$$
\left\langle v_{1} \otimes \cdots \otimes v_{k}, w_{1} \otimes \cdots \otimes w_{k}\right\rangle_{\left(\mathbb{R}^{d}\right)^{\otimes k}} \triangleq\left\langle v_{1}, w_{1}\right\rangle_{\mathbb{R}^{d}} \cdots\left\langle v_{k}, w_{k}\right\rangle_{\mathbb{R}^{d}} .
$$

The components $\left(\mathbb{R}^{d}\right)^{\otimes k}$ and $\left(\mathbb{R}^{d}\right)^{\otimes m}(k \neq m)$ are assumed to be orthogonal.
The following algebraic structure is critical in rough path theory.
DEFINITION 2.6. The free nilpotent Lie group $G^{(l)}$ of order $l$ is defined by

$$
G^{(l)} \triangleq \exp \left(\mathfrak{g}^{(l)}\right) \subseteq T^{(l)}
$$

The exponential function is a diffeomorphism under which $\mathfrak{g}^{(l)}$ in Notation 2.5 is the Lie algebra of $G^{(l)}$.

It will be useful in the sequel to have some basis available for the algebras introduced above. Recall that $\mathcal{A}(l)$ (respectively, $\left.\mathcal{A}_{1}(l)\right)$ denotes the set of words (respectively, nonempty words) over $\{1, \ldots, d\}$ of length, at most, $l$. We set $\mathrm{e}_{(\varnothing)} \triangleq 1$, and for each word $\alpha=\left(i_{1}, \ldots, i_{r}\right) \in \mathcal{A}_{1}(l)$, we set

$$
\begin{equation*}
\mathrm{e}_{(\alpha)} \triangleq \mathrm{e}_{i_{1}} \otimes \cdots \otimes \mathrm{e}_{i_{r}} \quad \text { and } \quad \mathrm{e}_{[\alpha]} \triangleq\left[\mathrm{e}_{i_{1}}, \ldots\left[\mathrm{e}_{i_{r-2}},\left[\mathrm{e}_{i_{r-1}}, \mathrm{e}_{i_{r}}\right]\right]\right] \tag{2.7}
\end{equation*}
$$

where $\left\{\mathrm{e}_{1}, \ldots, \mathrm{e}_{d}\right\}$ denotes the canonical basis of $\mathbb{R}^{d}$. Then, $\left\{\mathrm{e}_{(\alpha)}: \alpha \in \mathcal{A}(l)\right\}$ forms an orthonormal basis of $T^{(l)}$ under the Hilbert-Schmidt tensor norm. In addition, we also have $\mathfrak{g}^{(l)}=\operatorname{Span}\left\{\mathrm{e}_{[\alpha]}: \alpha \in \mathcal{A}_{1}(l)\right\}$.

As a closed subspace, $\mathfrak{g}^{(l)}$ inherits a canonical Hilbert structure from $T^{(l)}$ which makes it into a flat Riemannian manifold. The associated volume measure $d u$ (the Lebesgue measure) on $\mathfrak{g}^{(l)}$ is left invariant with respect to the product induced from the group structure on $G^{(l)}$ through the exponential diffeomorphism. In addition, for each $\lambda>0$, there is a dilation operation $\delta_{\lambda}: T^{(l)} \rightarrow T^{(l)}$, induced by $\delta_{\lambda}(a) \triangleq \lambda^{k} a$ if $a \in\left(\mathbb{R}^{d}\right)^{\otimes k}$, which satisfies the relation $\delta_{\lambda} \circ \exp =\exp \circ \delta_{\lambda}$ when restricted on $\mathfrak{g}^{(l)}$. Thanks to the fact that $\delta_{\lambda}(a)=\lambda^{k} a$ for any $a \in\left(\mathbb{R}^{d}\right)^{\otimes k}$ and recalling that $\mathcal{L}_{k}$ is introduced in Notation 2.5 , one can easily show that

$$
\begin{equation*}
d u \circ \delta_{\lambda}^{-1}=\lambda^{-v} d u \quad \text { where } v \triangleq \sum_{k=1}^{l} k \operatorname{dim}\left(\mathcal{L}_{k}\right) \tag{2.8}
\end{equation*}
$$

We always fix the Euclidean norm on $\mathbb{R}^{d}$ in the remainder of the paper. As far as the free nilpotent group $G^{(l)}$ is concerned, there are several useful metric structures. Among them we will use an extrinsic metric $\rho_{\mathrm{HS}}$ which can be defined easily due to the fact that $G^{(l)}$ is a subspace of $T^{(l)}$. Namely, for $g_{1}, g_{2} \in G^{(l)}$, we set

$$
\begin{equation*}
\rho_{\mathrm{HS}}\left(g_{1}, g_{2}\right) \triangleq\left\|g_{2}-g_{1}\right\|_{\mathrm{HS}}, \quad g_{1}, g_{2} \in G^{(l)} \tag{2.9}
\end{equation*}
$$

where the right hand side is induced from the Hilbert-Schmidt norm on $T^{(l)}$.
2.3. Path signatures and the fractional Brownian rough path. The stochastic differential equation (1.1), driven by a fractional Brownian motion $B$, is standardly solved in the rough paths sense. In this section we recall some basic facts about this notion of solution. We will also give some elements of rough paths expansions which are at the heart of our methodology in the present paper.

The link between free nilpotent groups and noisy equations, like (1.1), is made through the notion of signature. Recall that a continuous map $\mathbf{x}:\left\{(s, t) \in[0,1]^{2}: s \leq t\right\} \rightarrow T^{(l)}$ is called a multiplicative functional if, for $s<u<t$, one has $\mathbf{x}_{s, t}=\mathbf{x}_{s, u} \otimes \mathbf{x}_{u, t}$. A particular occurrence of this kind of map is given when one considers a path $w$ with finite variation and sets for $s \leq t$,

$$
\begin{equation*}
\mathbf{w}_{s, t}^{n}=\int_{s<u_{1}<\cdots<u_{n}<t} d w_{u_{1}} \otimes \cdots \otimes d w_{u_{n}} \tag{2.10}
\end{equation*}
$$

Then, the so-called truncated signature path of order $l$ associated with $w$ is the functional $S_{l}(w)_{,,,}:\left\{(s, t) \in[0,1]^{2}: s \leq t\right\} \rightarrow T^{(l)}$ defined by

$$
\begin{equation*}
S_{l}(w)_{s, t}:=1+\sum_{n=1}^{l} \mathbf{w}_{s, t}^{n} . \tag{2.11}
\end{equation*}
$$

It can be shown that the functional $S_{l}(w)_{\text {,., }}$ is multiplicative and takes values in the free nilpotent group $G^{(l)}$. The truncated signature of order $l$ for $w$ is the tensor element $S_{l}(w)_{0,1} \in$ $G^{(l)}$. It is simply denoted as $S_{l}(w)$.

A rough path can be seen as a generalization of the truncated signature path (2.11) to the nonsmooth situation. Specifically, the definition of Hölder rough paths can be summarized as follows.

DEFINITION 2.7. Let $\gamma \in(0,1)$. The space of weakly geometric $\gamma$-Hölder rough paths is the set of multiplicative functionals $\mathbf{x}:\left\{(s, t) \in[0,1]^{2}: s \leq t\right\} \rightarrow G^{[1 / \gamma]}$ such that

$$
\begin{equation*}
\|\mathbf{x}\|_{\gamma ; \mathrm{HS}} \triangleq \sum_{i=1}^{[1 / \gamma]} \sup _{0 \leq s<t \leq 1} \frac{\left\|\mathbf{x}_{s, t}^{i}\right\|_{\mathrm{HS}}}{|t-s|^{i \gamma}}<\infty \tag{2.12}
\end{equation*}
$$

An important subclass of weakly geometric $\gamma$-Hölder rough paths is the set of geometric $\gamma$-Hölder rough paths. These are multiplicative paths $\mathbf{x}$ with values in $G^{\lfloor 1 / \gamma\rfloor}$ such that $\|\mathbf{x}\|_{\gamma ; \text { HS }}$ is finite and such that there exists a sequence $\left\{x_{\varepsilon}: \varepsilon>0\right\}$ with $x_{\varepsilon} \in C^{\infty}\left([0,1] ; \mathbb{R}^{d}\right)$ satisfying

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0}\left\|\mathbf{x}-S_{[1 / \gamma]}\left(x_{\varepsilon}\right)\right\|_{\gamma ; \mathrm{HS}}=0 \tag{2.13}
\end{equation*}
$$

The notion of signature allows us to define a more intrinsic notion of distance on the free group $G^{(l)}$. This metric, known as the Carnot-Caratheodory metric, is defined by

$$
\rho_{\mathrm{CC}}\left(g_{1}, g_{2}\right) \triangleq\left\|g_{1}^{-1} \otimes g_{2}\right\|_{\mathrm{CC}}, \quad g_{1}, g_{2} \in G^{(l)}
$$

where the CC-norm $\|\cdot\|_{\text {CC }}$ is defined by

$$
\begin{equation*}
\|g\|_{\mathrm{CC}} \triangleq \inf \left\{\|w\|_{1-\mathrm{var}}: w \in C^{1-\mathrm{var}}\left([0,1] ; \mathbb{R}^{d}\right) \text { and } S_{l}(w)=g\right\} \tag{2.14}
\end{equation*}
$$

It can be shown that the infimum in (2.14) is attainable.
REMARK 2.8. It is well known that, for any $g \in G^{(l)}$, one can find a piecewise linear path $w$ such that $S_{l}(w)=g$ (cf. [13], for instance). In addition, a simple reparametrization allows one to take $w$ to be smooth with derivative compactly supported in $(0,1)$.

The HS and CC metrics are equivalent as seen from the following so-called ball-box estimate (cf. [13], Proposition 7.49).

PROPOSITION 2.9. Let $\rho_{\mathrm{HS}}$ and $\rho_{\mathrm{CC}}$ be the distances on $G^{(l)}$, respectively, defined by (2.9) and (2.14). For each $l \geq 1$, there exists a constant $C=C_{l}>0$ such that

$$
\begin{equation*}
\rho_{\mathrm{CC}}\left(g_{1}, g_{2}\right) \leq C \max \left\{\rho_{\mathrm{HS}}\left(g_{1}, g_{2}\right), \rho_{\mathrm{HS}}\left(g_{1}, g_{2}\right)^{\frac{1}{l}} \cdot \max \left\{1,\left\|g_{1}\right\|_{\mathrm{CC}}^{1-\frac{1}{l}}\right\}\right\} \tag{2.15}
\end{equation*}
$$

and

$$
\rho_{\mathrm{HS}}\left(g_{1}, g_{2}\right) \leq C \max \left\{\rho_{\mathrm{CC}}\left(g_{1}, g_{2}\right)^{l}, \rho_{\mathrm{CC}}\left(g_{1}, g_{2}\right) \cdot \max \left\{1,\left\|g_{1}\right\|_{\mathrm{CC}}^{l-1}\right\}\right\}
$$

for all $g_{1}, g_{2} \in G^{(l)}$. In particular,

$$
\|g\|_{\mathrm{CC}} \leq 1 \quad \Longrightarrow \quad\|g-\mathbf{1}\|_{\mathrm{HS}} \leq C\|g\|_{\mathrm{CC}}
$$

and

$$
\|g-\mathbf{1}\|_{\mathrm{HS}} \leq 1 \quad \Longrightarrow \quad\|g\|_{\mathrm{CC}} \leq C\|g-\mathbf{1}\|_{\mathrm{HS}}^{\frac{1}{\tau}}
$$

One of the main applications of rough path theory is to extend most stochastic calculus tools to a large class of Gaussian processes. The following result, borrowed from [9, 12], establishes this link for fractional Brownian motion.

Proposition 2.10. Let B be a fractional Brownian motion with Hurst parameter $H>$ $1 / 4$. Then, $B$ admits a lift $\mathbf{B}$ as a geometric rough path of order $[1 / \gamma]$ for any $\gamma<H$.

Let us now turn to the definition of rough differential equations. There are several equivalent ways to introduce this notion, among which we will choose to work with Taylor-type expansions since they are more consistent with our later developments. To this aim, let us first consider a bounded variation path $w$ and the following ordinary differential equation driven by $w$ :

$$
\begin{equation*}
d x_{t}=\sum_{\alpha=1}^{d} V_{\alpha}\left(x_{t}\right) d w_{t}^{\alpha} \tag{2.16}
\end{equation*}
$$

where the $V_{\alpha}$ 's are $C_{b}^{\infty}$ vector fields. For any given word $\alpha=\left(i_{1}, \ldots, i_{r}\right)$ over the letters $\{1, \ldots, d\}$, define the vector field $V_{(\alpha)} \triangleq\left(V_{i_{1}} \cdots\left(V_{i_{r-2}}\left(V_{i_{r-1}} V_{i_{r}}\right)\right)\right)$, where we have identified a vector field with a differential operator, so that $V_{i} V_{j}$ means differentiating $V_{j}$ along direction $V_{i}$. Classically, a formal Taylor expansion of the solution $x_{t}$ to (2.16) is given by

$$
\begin{equation*}
x_{s, t} \sim \sum_{k=1}^{\infty} \sum_{i_{1}, \ldots, i_{k}=1}^{d} V_{\left(i_{1}, \ldots, i_{k}\right)}\left(x_{s}\right) \int_{s<u_{1}<\cdots<u_{k}<t} d w_{u_{1}}^{i_{1}} \cdots d w_{u_{k}}^{i_{k}} \tag{2.17}
\end{equation*}
$$

where we have set $x_{s, t}=x_{t}-x_{s}$. This expansion can be rephrased in more geometrical terms. Specifically, we define the following Taylor approximation function on $\mathfrak{g}^{(l)}$. Recall that the sets of words $\mathcal{A}(l), \mathcal{A}_{1}(l)$ are introduced at the beginning of Section 1.2.

DEFINITION 2.11. For each $l \geq 1$, we define the Taylor approximation function $F_{l}$ : $\mathfrak{g}^{(l)} \times \mathbb{R}^{N} \rightarrow \mathbb{R}^{N}$ of order $l$ associated with the ODE (2.16) by

$$
F_{l}(u, x) \triangleq \sum_{\alpha \in \mathcal{A}_{1}(l)} V_{(\alpha)}(x) \cdot(\exp u)^{\alpha}, \quad(u, x) \in \mathfrak{g}^{(l)} \times \mathbb{R}^{N}
$$

where the exponential function is defined on $T^{(l)}$ by (2.6) and $(\exp (u))^{\alpha}$ is the coefficient of $\exp (u)$ with respect to the tensor basis element $\mathrm{e}_{(\alpha)}$ introduced in (2.7). We also say that $u \in \mathfrak{g}^{(l)}$ joins $x$ to $y$ in the sense of Taylor approximation if $y=x+F_{l}(u, x)$.

With Definition 2.11 in hand, we can recast the formal expansion (2.17) (truncated at an arbitrary degree $l$ ) in the following way:

$$
\begin{equation*}
x_{s, t} \sim F_{l}\left(\log \left(S_{l}(w)_{s, t}\right), x_{s}\right) \tag{2.18}
\end{equation*}
$$

where the function $\log$ is the inverse of the exponential map for $G^{(l)}$, and $S_{l}(w)_{s, t}$ is the truncated signature path of $w$ defined by (2.11). In order to define rough differential equations, a natural idea is to extend this approximation scheme to rough paths. We get a definition which is stated below in the fractional Brownian motion case.

Definition 2.12. Let $B$ be a fractional Brownian motion with Hurst parameter $H>$ $1 / 4$, and let $\gamma \in(1 / 4, H)$ be fixed. Consider the rough path lift $\mathbf{B}$ of $B$, as in Proposition 2.10. Let $\left\{V_{\alpha}: 1 \leq \alpha \leq d\right\}$ be a family of $C_{b}^{\infty}$ vector fields on $\mathbb{R}^{N}$. We say that $X$ is a solution to the rough differential equation (1.1) if, for all $(s, t) \in[0,1]^{2}$ such that $s<t$, we have

$$
\begin{equation*}
X_{s, t}=F_{[1 / \gamma]-1}\left(\log \left(S_{[1 / \gamma]-1}(\mathbf{B})_{s, t}\right), X_{s}\right)+R_{S, t} \tag{2.19}
\end{equation*}
$$

where $R_{s, t}$ is an $\mathbb{R}^{N}$-valued remainder such that there exists $\varepsilon>0$ satisfying

$$
\sup _{0 \leq s<t \leq 1} \frac{\left|R_{s, t}\right|}{|t-s|^{1+\varepsilon}}<\infty
$$

Roughly speaking, Definition 2.12 says that the expansion of the solution $X$ to a rough differential equation should coincide with (2.17) up to a remainder with Hölder regularity greater than 1. This approach goes back to Davie [10], and it can be shown to coincide with more classical notions of solutions. The following existence and uniqueness result is fundamental in rough path theory.

Proposition 2.13. Under the same conditions as in Definition 2.12, there exists a unique solution to equation (1.1) understood in the sense of (2.19).
3. Local estimate for the control distance function. In this section we develop the proofs of Theorem 1.2 and Theorem 1.3. In contrast to the elliptic case that is treated in [15], a major difficulty in the hypoelliptic case is that one cannot explicitly construct a CameronMartin path joining two points in the sense of Definition 1.1 in any easy way. The analysis of Cameron-Martin norms also becomes more involved. We detail the steps in what follows, starting with some preliminary lemmas.
3.1. Preliminary results. As we mentioned above, it is difficult to explicitly construct a Cameron-Martin path joining $x$ to $y$ in the sense of differential equation in the hypoelliptic case. However, it is possible to find some $u \in \mathfrak{g}^{(l)}$ joining $x$ to $y$ in the sense of Taylor approximation, that is, $y=x+F_{l}(u, x)$ where the function $F_{l}(u, x)$ is introduced in Definition 2.11. This is the content of the following lemma proved in [20], Lemma 3.15. Recall that $l_{0}$ is the hypoellipticity constant in the assumption (1.2).

LEMmA 3.1. For each $l \geq l_{0}$, there exist constants $r, A>0$, depending only on $l$ and the vector fields, and $a C_{b}^{\infty}$-function

$$
\Psi_{l}:\left\{u \in \mathfrak{g}^{(l)}:\|u\|_{\mathrm{HS}}<r\right\} \times \mathbb{R}^{N} \times\left\{\eta \in \mathbb{R}^{N}:|\eta|<r\right\} \rightarrow \mathfrak{g}^{(l)}
$$

such that, for all $u, x, \eta$ in the domain of $\Psi_{l}$, we have:
(i) $\Psi_{l}(u, x, 0)=u$;
(ii) $\left\|\Psi_{l}(u, x, \eta)-u\right\|_{\mathrm{HS}} \leq A|\eta|$;
(iii) $F_{l}\left(\Psi_{l}(u, x, \eta), x\right)=F_{l}(u, x)+\eta$.

REMARK 3.2. The property (ii) is not explicitly stated in [20], Lemma 3.15. It is only clear when one develops the construction of $\Psi_{l}$ carefully.

The intuition behind the function $\Psi_{l}$ can be described as follows. Let $y \triangleq x+F_{l}(u, x)$ so that $u$ joins $x$ to $y$ in the sense of Taylor approximation. Then, $v \triangleq \Psi_{l}(u, x, \eta)$ joins $x$ to $y+\eta$, that is, $x+F_{l}(v, x)=y+\eta$. In particular, $\Psi_{l}(0, x, y-x)$ gives an element in $\mathfrak{g}^{(l)}$, which joins $x$ to $y$ in the sense of Taylor approximation, provided $|y-x|<r$.

We begin with some preliminary steps toward the proof of Theorem 1.2, namely, the lower bound on the control distance function $d(\cdot, \cdot)$ and the upper bound for the case $H<1 / 2$.

Lemma 3.3. Assume that the vector fields in equation (1.1) satisfy the uniform hypoellipticity assumption (1.2) with constant $l_{0}$. Let $d=d_{H}$ be the control distance introduced in Definition 1.1. Then, the following bounds hold true:
(i) For all $H \in(1 / 4,1)$ and $x, y$ such that $|x-y| \leq 1$, we have

$$
d(x, y) \geq C_{1}|x-y|
$$

(ii) Whenever $H \in(1 / 4,1 / 2)$, we have

$$
d(x, y) \leq C_{2}|x-y|^{\frac{1}{L_{0}}}
$$

Proof. Claim (i) follows from the exact same argument as in [15], Theorem 3.3, for the elliptic case. Claim (ii) stems from the fact that, when $H<1 / 2$, we have

$$
\begin{equation*}
d(x, y) \leq C_{H} d_{\mathrm{BM}}(x, y) \tag{3.1}
\end{equation*}
$$

where $d_{B M}$ stands for the distance for the Brownian motion case. Note that (3.1) can be easily justified by the fact that

$$
d(x, y) \leq\|h\|_{\overline{\mathcal{H}}} \leq C_{H}\|h\|_{W^{1,2}}
$$

for any $h$ joining $x$ and $y$ in the differential equation sense (cf. [15], Lemma 2.8). Then, with (3.1) in hand our claim (ii) follows from the Brownian hypoelliptic analysis [20].

From now on, we focus on the case $H>1 / 2$. It is not surprising that this is the harder case since the Cameron-Martin subspace $\overline{\mathcal{H}}$ gets smaller as $H$ increases. First, we need to make use of the following scaling property of the Cameron-Martin norm. The proof is immediate by using the relation (2.4) and a simple change of variables. We denote $\overline{\mathcal{H}}([0, T])$ (respectively, $\left.d_{T}(x, y)\right)$ as the Cameron-Martin subspace (respectively, the control distance function) associated with fractional Brownian motion over [ $0, T]$.

Lemma 3.4. Let $0<T_{1}<T_{2}$. Given $h \in \overline{\mathcal{H}}\left(\left[0, T_{1}\right]\right)$, define $\tilde{h}_{t} \triangleq h_{T_{1} t / T_{2}}$ for $0 \leq t \leq T_{2}$. Then, $\tilde{h} \in \overline{\mathcal{H}}\left[0, T_{2}\right]$, and

$$
\begin{equation*}
\|\tilde{h}\|_{\overline{\mathcal{H}}\left(\left[0, T_{2}\right]\right)}=\left(\frac{T_{1}}{T_{2}}\right)^{H}\|h\|_{\overline{\mathcal{H}}\left(\left[0, T_{1}\right]\right)} \tag{3.2}
\end{equation*}
$$

In particular, we have

$$
\begin{equation*}
d_{1}(x, y)=T^{H} d_{T}(x, y) \quad \forall T>0, x, y \in \mathbb{R}^{N} . \tag{3.3}
\end{equation*}
$$

We also need the following lemma about the free nilpotent group $G^{(l)}$ which allows us to choose a "regular" path $\gamma$ with $S_{l}(\gamma)=u$ for all $u \in G^{(l)}$.

Lemma 3.5. Let $l \geq 1$. For each $M>0$, there exists a constant $C=C_{l, M}>0$ such that, for every $u \in G^{(l)}$ with $\|u\|_{\mathrm{CC}} \leq M$, one can find a smooth path $\gamma:[0,1] \rightarrow \mathbb{R}^{d}$ which satisfies:
(i) $S_{l}(\gamma)=u$;
(ii) $\dot{\gamma}$ is supported on $[1 / 3,2 / 3]$;
(iii) $\|\ddot{\gamma}\|_{\infty ;[0,1]} \leq C$.

Proof. We first prove the claim for a generic element $u \in \exp \left(\mathcal{L}_{k}\right)$, seen as an element of $G^{(k)}$. Let $\left\{a_{1}, \ldots, a_{d_{k}}\right\}$ be a basis of $\mathcal{L}_{k}$ where $d_{k} \triangleq \operatorname{dim} \mathcal{L}_{k}$. Given $u \in \exp \left(\mathcal{L}_{k}\right)$, we can write $u=\exp (a)$ with

$$
\begin{equation*}
a=\lambda_{1} a_{1}+\cdots+\lambda_{d_{k}} a_{d_{k}} \in \mathcal{L}_{k} \tag{3.4}
\end{equation*}
$$

for some $\lambda_{1}, \ldots, \lambda_{d_{k}} \in \mathbb{R}$. Since we assume that $\|u\|_{\mathrm{CC}} \leq M$, according to the ball-box estimate (cf. Proposition 2.9) and the fact that $a \in \mathcal{L}_{k}$, we have

$$
\begin{equation*}
\|a\|_{\mathrm{HS}}=\|u-\mathbf{1}\|_{\mathrm{HS}} \leq C_{1, l, M} \tag{3.5}
\end{equation*}
$$

Moreover, $\mathcal{L}_{k}$ is a finite dimensional vector space on which all norms are equivalent. Thus, relation (3.5) yields

$$
\begin{equation*}
\max _{1 \leq i \leq d_{k}}\left|\lambda_{i}\right| \leq C_{2, l, M} \tag{3.6}
\end{equation*}
$$

Now, recall from Remark 2.8 that, for each $a_{i}$ in (3.4), one can choose a smooth path $\alpha_{i}:[0,1] \rightarrow \mathbb{R}^{d}$ such that $S_{k}\left(\alpha_{i}\right)=\exp \left(a_{i}\right)$ and $\dot{\alpha}_{i}$ is supported on $[1 / 3,2 / 3]$. Set

$$
R_{k} \triangleq \max \left\{\left\|\ddot{\alpha}_{i}\right\|_{\infty ;[0,1]}: 1 \leq i \leq d_{k}\right\} .
$$

Note that $R_{k}$ is a constant depending only on $k$. We construct a smooth path $\gamma:\left[0, d_{k}\right] \rightarrow \mathbb{R}^{d}$ by

$$
\begin{equation*}
\gamma \triangleq\left(\left|\lambda_{1}\right|^{\frac{1}{k}} \alpha_{1}^{\operatorname{sgn}\left(\lambda_{1}\right)}\right) \sqcup \cdots \sqcup\left(\left|\lambda_{d_{k}}\right|^{\frac{1}{k}} \alpha_{d_{k}}^{\operatorname{sgn}\left(\lambda_{d_{k}}\right)}\right) \tag{3.7}
\end{equation*}
$$

where $\alpha_{i}^{-1}$ denotes the reverse of $\alpha_{i}$, and $\sqcup$ denotes path concatenation. Then, $\dot{\gamma}$ is obviously compactly supported, and we also claim that $S_{k}(\gamma)=u$. Indeed, it follows from (3.7) that

$$
\begin{align*}
S_{k}(\gamma) & =S_{k}\left(\left|\lambda_{1}\right|^{\frac{1}{k}} \alpha_{1}^{\operatorname{sgn}\left(\lambda_{1}\right)}\right) \otimes \cdots \otimes S_{k}\left(\left|\lambda_{d_{k}}\right|^{\frac{1}{k}} \alpha_{d_{k}}^{\operatorname{sgn}\left(\lambda_{d_{k}}\right)}\right) \\
& =\delta_{\left|\lambda_{1}\right|^{\frac{1}{k}}}\left(S_{k}\left(\alpha_{1}^{\operatorname{sgn}\left(\lambda_{1}\right)}\right)\right) \otimes \cdots \otimes \delta_{\left|\lambda_{d_{k}}\right|^{\frac{1}{k}}}\left(S_{k}\left(\alpha_{d_{k}}^{\operatorname{sgn}\left(\lambda_{d_{k}}\right)}\right)\right)  \tag{3.8}\\
& =\delta_{\left|\lambda_{1}\right|^{\frac{1}{k}}}\left(\exp \left(\operatorname{sgn}\left(\lambda_{1}\right) a_{1}\right)\right) \otimes \cdots \otimes \delta_{\left|\lambda_{d_{k}}\right|^{\frac{1}{k}}}\left(\exp \left(\operatorname{sgn}\left(\lambda_{d_{k}}\right) a_{d_{k}}\right)\right)
\end{align*}
$$

where we have used the properties of the dilation, recalled in Section 2.2, and the relation between signatures and $G^{(l)}$, given in (2.10)-(2.11). In addition, since each element $\exp \left(\lambda_{i} a_{i}\right)$ above sits in $\exp \left(\mathcal{L}_{k}\right)$, the tensor product in $G^{(k)}$ is reduced to

$$
\begin{equation*}
S_{k}(\gamma)=\exp \left(\lambda_{1} a_{1}\right) \otimes \cdots \otimes \exp \left(\lambda_{d_{k}} a_{d_{k}}\right)=\exp (a)=u \tag{3.9}
\end{equation*}
$$

We have thus found a path $\gamma$ with compactly supported derivative such that $S_{k}(\gamma)=u$. In addition, from the definition of $R_{k}$ and (3.6) we have

$$
\|\ddot{\gamma}\|_{\infty ;\left[0, d_{k}\right]} \leq R_{k} \cdot\left(\max _{1 \leq i \leq d_{k}}\left|\lambda_{i}\right|\right)^{\frac{1}{k}} \leq C_{3, l, M}
$$

By suitable rescaling and adding trivial pieces on both ends, if necessary, we may assume that $\gamma$ is defined on $[0,1]$ and $\dot{\gamma}$ is supported on [1/3, 2/3]. In this way we have

$$
\|\ddot{\gamma}\|_{\infty ;[0,1]} \leq C_{k} \cdot C_{3, k, M} \triangleq C_{4, k, M}
$$

where $C_{k}$ is the constant coming from the rescaling. Therefore, our assertion (i)-(iii) holds for $u$ which are elements of $\exp \left(\mathcal{L}_{k}\right)$.

With the help from the previous special case, we now prove the lemma by induction on $l$. The case when $l=1$ is obvious, as we can simply choose $\gamma$ to be a straight line segment. Suppose now that the claim is true on $G^{(l-1)}$. We let $M>0$ and $u \in G^{(l)}$ with $\|u\|_{\mathrm{CC}} \leq M$. Define $v \triangleq \pi^{(l-1)}(u)$ where $\pi^{(l-1)}: G^{(l)} \rightarrow G^{(l-1)}$ is the canonical projection. We obviously have

$$
\|v\|_{\mathrm{CC}} \leq\|u\|_{\mathrm{CC}} \leq M
$$

where the CC-norm of $v$ is taken on the group $G^{(l-1)}$. According to the induction hypothesis, there exists a constant $C_{l-1, M}$ such that we can find a smooth path $\alpha:[0,1] \rightarrow \mathbb{R}^{d}$ which satisfies (i)-(iii) in the assertion of Lemma 3.5, for $v=S_{l-1}(\alpha)$ and constant $C_{l-1, M}$. Define

$$
\begin{equation*}
w \triangleq\left(S_{l}(\alpha)\right)^{-1} \otimes u \tag{3.10}
\end{equation*}
$$

where the tensor product is defined on $G^{(l)}$. Then, note that, owing to the fact that $\|u\|_{\mathrm{CC}} \leq$ $M$, we have

$$
\|w\|_{\mathrm{CC}} \leq\left\|S_{l}(\alpha)\right\|_{\mathrm{CC}}+\|u\|_{\mathrm{CC}} \leq\|\alpha\|_{1-\mathrm{var} ;[0,1]}+\|u\|_{\mathrm{CC}} \leq \frac{1}{2}\|\ddot{\alpha}\|_{\infty ;[0,1]}+M .
$$

Therefore, thanks to the induction procedure applied to $v=S_{l-1}(\alpha)$, we get

$$
\|w\|_{\mathrm{CC}} \leq \frac{1}{2} C_{l-1, M}+M \triangleq C_{5, l, M}
$$

We claim that $w \in \exp \left(\mathcal{L}_{l}\right)$. This can be proved in the following way:
(i) Write $u=\exp \left(l_{0}+l_{h}\right)$, where $l_{0} \in \mathfrak{g}^{(l-1)}$ and $l_{h} \in \mathcal{L}_{l}$. Recall $v \triangleq \pi^{(l-1)}(u)$. We argue that $v=\exp \left(l_{0}\right) \in G^{(l-1)}$ as follows: since $l_{h} \in \mathcal{L}_{l}$, any product of the form $l_{h}^{p} \otimes l_{0}^{q}=0$ whenever $p, q>0$. Taking into account the definition (2.6) of the exponential function, we get that

$$
\begin{equation*}
u=\exp \left(l_{0}+l_{h}\right) \quad \Longrightarrow \quad v=\exp \left(l_{0}\right) \in G^{(l-1)} \tag{3.11}
\end{equation*}
$$

(ii) Recall that our induction hypothesis asserts that $v=S_{l-1}(\alpha)$; thus, according to (3.11), we have $S_{l-1}(\alpha)=\exp \left(l_{0}\right)$. Thanks to the same kind of argument as in (i), we get $S_{l}(\alpha)=\exp \left(l_{0}+l_{h}^{\prime}\right) \in G^{(l)}$ for some $l_{h}^{\prime} \in \mathcal{L}_{l}$.
(iii) In order to conclude that $w \in \exp \left(\mathcal{L}_{l}\right)$, we go back to relation (3.10) which can now be read as

$$
w=\left(\exp \left(l_{0}+l_{h}^{\prime}\right)\right)^{-1} \otimes \exp \left(l_{0}+l_{h}\right)
$$

According to Campbell-Baker-Hausdorff formula and taking into account the fact that

$$
\left[l_{0}, l_{0}\right]=\left[l_{0}, l_{h}\right]=\left[l_{0}, l_{h}^{\prime}\right]=\left[l_{h}, l_{h}^{\prime}\right]=0 \in \mathfrak{g}^{(l)}
$$

we conclude that $w=\exp \left(l_{h}-l_{h}^{\prime}\right)$, and thus $w \in \exp \left(\mathcal{L}_{l}\right)$.
We are now ready to summarize our information and to conclude our induction procedure. Namely, for $u \in G^{(l)}$, we can recast relation (3.10) as

$$
\begin{equation*}
u=S_{l}(\alpha) \otimes w \tag{3.12}
\end{equation*}
$$

and we have just proved that $w \in \exp \left(\mathcal{L}_{l}\right)$. Hence, relation (3.9) asserts that $w$ can be written as $w=S_{l}(\beta)$, where $\beta:[0,1] \rightarrow \mathbb{R}^{d}$ satisfying relation (i)-(iii) in Lemma 3.5 with $C=$ $C_{6, l, M}$. Now, set $\gamma \triangleq \alpha \sqcup \beta$, and rescale it so that it is defined on $[0,1]$ and its derivative path is supported on $[1 / 3,2 / 3]$. Then, recalling our decomposition (3.12), we have

$$
S_{l}(\gamma)=S_{l}(\alpha) \otimes S_{l}(\beta)=S_{l}(\alpha) \otimes w=u
$$

and, moreover, the following upper bound holds true:

$$
\|\ddot{\gamma}\|_{\infty ;[0,1]} \leq 36 \max \left\{\|\ddot{\alpha}\|_{\infty ;[0,1]},\|\ddot{\beta}\|_{\infty ;[0,1]}\right\} \leq C_{7, l, M} .
$$

Therefore, our induction procedure is established which finishes the proof.
We conclude this subsection by stating a convention on the group $G^{(l)}$ which will ease notation in our future computations.

Convention 3.6. Since $\mathfrak{g}^{(l)}$ is a finite dimensional vector space on which differential calculus is easier to manage, we will frequently identify $G^{(l)}$ with $\mathfrak{g}^{(l)}$ through the exponential diffeomorphism without further mention. In this way, for instance, $S_{l}(w)=u$ means $S_{l}(w)=$ $\exp (u)$ if $u \in \mathfrak{g}^{(l)}$. The same convention will apply to other similar relations when the meaning is clear from context. For norms on $\mathfrak{g}^{(l)}$, we denote $\|u\|_{\mathrm{CC}} \triangleq\|\exp (u)\|_{\mathrm{CC}}$. As for the HSnorm, note that

$$
C_{1, l}\|u\|_{\mathrm{HS}} \leq\|\exp (u)-\mathbf{1}\|_{\mathrm{HS}} \leq C_{2, l}\|u\|_{\mathrm{HS}}
$$

for all $u \in \mathfrak{g}^{(l)}$ satisfying $\|\exp (u)-\mathbf{1}\|_{\mathrm{HS}} \wedge\|u\|_{\mathrm{HS}} \leq 1$. Therefore, up to a constant depending only on $l$, the notation $\|u\|_{\text {HS }}$ can either mean the HS-norm of $u \operatorname{or} \exp (u) \mathbf{- 1}$. This will not matter because we are only concerned with local estimates. The same convention applies to the distance functions $\rho_{\mathrm{CC}}$ and $\rho_{\mathrm{HS}}$.
3.2. Proof of Theorem 1.2. In this section we give the details to complete the proof of Theorem 1.2, namely, the local comparison between the distance $d(\cdot, \cdot)$ and the Euclidean distance. Thanks to Lemma 3.3, we only focus on the upper bound for $H>1 / 2$.

Recall that $\Psi_{l}(u, x, \eta)$ is the function given by Lemma 3.1. This function allows us to construct elements in $\mathfrak{g}^{(l)}$ joining two points in the sense of Taylor approximation locally. In what follows, we take $l=l_{0}$ (where $l_{0}$ stands for the hypoellipticity constant), and we will omit the subscript $l$ for simplicity (e.g., $F=F_{l}$ and $\Psi=\Psi_{l}$ ). We will also identify $G^{(l)}$ with $\mathfrak{g}^{(l)}$ in the way mentioned in Convention 3.6. We now divide our proof into several steps.
Step 1: Construction of an approximating sequence. Let $\delta<r$ be a constant to be chosen later on, where $r$ is the constant appearing in the domain of $\Psi$ in Lemma 3.1. Consider $x, y \in \mathbb{R}^{N}$ with $|x-y|<\delta$.

We are going to construct three sequences

$$
\left\{x_{m}\right\} \subseteq \mathbb{R}^{N}, \quad\left\{u_{m}\right\} \subseteq \mathfrak{g}^{\left(l_{0}\right)}, \quad\left\{h_{m}\right\} \subseteq C^{\infty}\left([0,1] ; \mathbb{R}^{d}\right)
$$

inductively. We start with $x_{1} \triangleq x$ and define the rest of them by the following general procedure in the order

$$
u_{1} \rightarrow h_{1} \rightarrow x_{2} \rightarrow u_{2} \rightarrow h_{2} \rightarrow x_{3} \rightarrow \cdots
$$

To this aim, suppose we have already defined $x_{m}$. Set

$$
\begin{equation*}
u_{m} \triangleq \Psi\left(0, x_{m}, y-x_{m}\right) \quad \text { and } \quad \bar{u}_{m} \triangleq \delta_{\left\|u_{m}\right\|_{\mathrm{CC}}^{-1}} u_{m} . \tag{3.13}
\end{equation*}
$$

By Lemma 3.1 the first condition in (3.13) states that $u_{m}$ is an element of $\mathfrak{g}^{\left(l_{0}\right)}$ such that

$$
\begin{equation*}
x_{m}+F\left(u_{m}, x_{m}\right)=y, \tag{3.14}
\end{equation*}
$$

while the second condition in (3.13) ensures that $\left\|\bar{u}_{m}\right\|_{\text {CC }}=1$. Once $u_{m}$ is defined, we construct $h_{m}$ in the following way: let $\bar{h}_{m}:[0,1] \rightarrow \mathbb{R}^{d}$ be the smooth path given by Lemma 3.5 such that $S_{l_{0}}\left(\bar{h}_{m}\right)=\bar{u}_{m}, \dot{\bar{h}}_{m}$ is supported on $[1 / 3,2 / 3]$, and $\left\|\ddot{\bar{h}}_{m}\right\|_{\infty ;[0,1]} \leq C_{l_{0}}$. Define

$$
\begin{equation*}
h_{m} \triangleq\left\|u_{m}\right\|_{\mathrm{CC}} \bar{h}_{m} \tag{3.15}
\end{equation*}
$$

so that the truncated signature of $h_{m}$ is exactly $u_{m}$ (here, recall the Convention 3.6). More specifically, we have

$$
S_{l_{0}}\left(h_{m}\right)=S_{l_{0}}\left(\left\|u_{m}\right\|_{\mathrm{CC}} \cdot \bar{h}_{m}\right)=\delta_{\left\|u_{m}\right\|_{\mathrm{CC}}}\left(S_{l_{0}}\left(\bar{h}_{m}\right)\right)=\delta_{\left\|u_{m}\right\|_{\mathrm{CC}}}\left(\bar{u}_{m}\right)=u_{m}
$$

Taking into account the definition (2.14) of the CC-norm, it is immediate that

$$
\begin{equation*}
\left\|u_{m}\right\|_{\mathrm{CC}} \leq\left\|h_{m}\right\|_{1-\mathrm{var} ;[0,1]} \leq\left\|u_{m}\right\|_{\mathrm{CC}}\left\|\bar{h}_{m}\right\|_{1-\mathrm{var} ;[0,1]} \leq C_{l_{0}}\left\|u_{m}\right\|_{\mathrm{CC}} \tag{3.16}
\end{equation*}
$$

where the last inequality stems from the fact that $\bar{h}_{m}$ has a bounded second derivative. Eventually, we define

$$
\begin{equation*}
x_{m+1} \triangleq \Phi_{1}\left(x_{m} ; h_{m}\right) \tag{3.17}
\end{equation*}
$$

where we recall that $\Phi_{t}(x ; h)$ is the solution flow of the $\operatorname{ODE}(2.16)$ driven by $h$ over $[0,1]$.
Step 2: Checking the condition $\left|y-x_{m}\right|<r$. Recall that, in Lemma 3.1, we have to impose $\|u\|_{\text {HS }}<r$ and $|\eta|<r$ in order to apply $\Psi$. In the context of (3.13), it means that we should make sure that

$$
\begin{equation*}
\left|y-x_{m}\right|<r \quad \text { for all } m . \tag{3.18}
\end{equation*}
$$

We will now choose $\delta_{1}$ small enough such that if $|y-x|<\delta_{1}$, then (3.18) is satisfied. This will guarantee that $u_{m}$ is well defined by Lemma 3.1, and we will also be able to write down several useful estimates for $x_{m}$ and $u_{m}$. Our first condition on $\delta_{1}$ is that $\delta_{1} \leq r$, so that if $|x-y|<\delta_{1}$, we can define $u_{1}$ by a direct application of Lemma 3.1. We will now prove by induction that if $\delta_{1}$ is chosen small enough, then condition (3.18) is satisfied. To this aim, assume that $\left|x_{m}-y\right|<\delta_{1}$. Then, one can apply Lemma 3.1 in order to define $u_{m}, h_{m}$ and $x_{m+1}$. We also get the following estimate:

$$
\begin{equation*}
\left\|u_{m}\right\|_{\mathrm{HS}} \leq A\left|x_{m}-y\right|<A \delta_{1} \tag{3.19}
\end{equation*}
$$

where $A$ is the constant appearing in Lemma 3.1. In addition, let us require $\delta_{1} \leq 1 / A$ so that $\left\|u_{m}\right\|_{\mathrm{HS}} \leq 1$. Recalling relations (3.14) and (3.17), we get

$$
\left|x_{m+1}-y\right|=\left|\Phi_{1}\left(x_{m}, h_{m}\right)-x_{m}-F\left(S_{l_{0}}\left(h_{m}\right), x_{m}\right)\right| .
$$

Thus, applying, successively, the Taylor-type estimate of [13], Proposition 10.3, and relation (3.16), we end up with

$$
\left|x_{m+1}-y\right| \leq C_{V, l_{0}}\left\|h_{m}\right\|_{1-\mathrm{var} ;[0,1]}^{1+l_{0}} \leq C_{V, l_{0}}\left\|u_{m}\right\|_{\mathrm{CC}}^{1+l_{0}}
$$

The quantity $\left\|u_{m}\right\|_{\text {CC }}$ above can be bounded thanks to the ball-box estimate of Proposition 2.9 for which we observe that the dominating term in (2.15) is $\rho_{\mathrm{HS}}\left(g_{1}, g_{2}\right)^{1 / l_{0}}$ since our element $u_{m}$ is bounded by one in HS-norm. We get

$$
\left|x_{m+1}-y\right| \leq C_{V, l_{0}}\left\|u_{m}\right\|_{\mathrm{CC}}^{1+l_{0}} \leq C_{V, l_{0}}\left\|u_{m}\right\|_{\mathrm{HS}}^{1+\frac{1}{l_{0}}} \leq C_{V, l_{0}} A^{1+\frac{1}{l_{0}}}\left|x_{m}-y\right|^{1+\frac{1}{l_{0}}}
$$

Summarizing our considerations so far, we have obtained the estimate

$$
\begin{equation*}
\left|x_{m+1}-y\right| \leq C_{1, V, l_{0}}\left\|u_{m}\right\|_{\mathrm{CC}}^{1+l_{0}} \leq C_{2, V, l_{0}}\left|x_{m}-y\right|^{1+\frac{1}{l_{0}}} \tag{3.20}
\end{equation*}
$$

On top of the inequalities $\delta_{1}<r$ and $\delta_{1} \leq 1 / A$ imposed previously, we will also assume that $C_{2, V, l_{0}} \delta_{1}^{1 / l_{0}} \leq 1 / 2$ which easily yields the relation

$$
\begin{equation*}
\left|x_{m+1}-y\right| \leq \frac{1}{2}\left|x_{m}-y\right|<\frac{1}{2} \delta_{1}<\delta_{1} \tag{3.21}
\end{equation*}
$$

For our future computations we will thus set

$$
\delta_{1} \triangleq r \wedge A^{-1} \wedge\left(2 C_{2, V, l_{0}}\right)^{-l_{0}}
$$

According to our bound (3.21), we can guarantee that if $|x-y|<\delta_{1}$, then $\left|x_{m}-y\right|<\delta_{1}<r$ for all $m$. In addition, an easy induction procedure performed on inequality (3.21) leads to the following relation, valid for all $m \geq 1$ :

$$
\begin{equation*}
\left|x_{m}-y\right| \leq 2^{-(m-1)}|x-y| \tag{3.22}
\end{equation*}
$$

Together with the second inequality of (3.20), we obtain that

$$
\begin{equation*}
\left\|u_{m}\right\|_{\mathrm{CC}} \leq C_{3, V, l_{0}} 2^{-\frac{m}{l_{0}}}|x-y|^{\frac{1}{l_{0}}} \quad \forall m \geq 1 \tag{3.23}
\end{equation*}
$$

We will now choose a constant $\delta_{2} \leq \delta_{1}$ such that the sequence $\left\{\left\|u_{m}\right\|_{\mathrm{CC}} ; m \geq 1\right\}$ is decreasing with $m$ when $|x-y|<\delta_{2}$. This property will be useful for our future considerations. Toward this aim, observe that applying, successively, (2.15), (3.19) and (3.20), we get

$$
\begin{equation*}
\left\|u_{m+1}\right\|_{\mathrm{CC}} \leq C_{l_{0}}\left\|u_{m+1}\right\|_{\mathrm{HS}}^{\frac{1}{l_{0}}} \leq C_{4, V, l_{0}}\left\|u_{m}\right\|_{\mathrm{CC}}^{1+\frac{1}{L_{0}}} \tag{3.24}
\end{equation*}
$$

Hence, invoking the second inequality in (3.20), we have

$$
\begin{equation*}
\left\|u_{m+1}\right\|_{\mathrm{CC}} \leq C_{5, V, l_{0}}|x-y|^{\frac{1}{I_{0}^{2}}}\left\|u_{m}\right\|_{\mathrm{CC}} \tag{3.25}
\end{equation*}
$$

Therefore, let us consider a new constant $\delta_{2}>0$ such that

$$
C_{5, V, l_{0}} \delta_{2}^{\frac{1}{l_{0}^{2}}}<1
$$

If we choose $|x-y|<\delta$ with $\delta \triangleq \delta_{1} \wedge \delta_{2}$, equation (3.25) can be recast as

$$
\begin{equation*}
\left\|u_{m+1}\right\|_{\mathrm{CC}} \leq\left\|u_{m}\right\|_{\mathrm{CC}} \tag{3.26}
\end{equation*}
$$

Note that $\delta=\delta_{1} \wedge \delta_{2}$ depends only on $l_{0}$ and the vector fields but not on the Hurst parameter $H$. We have thus shown that the application of Lemma 3.1 is valid in our context.

Step 3: Construction of a path joining $x$ and $y$ in the sense of differential equation. Our next aim is to obtain a path $\tilde{h}$ joining $x$ and $y$ along the flow of equation (2.16). The first step in this direction is to rescale $h_{m}$ in a suitable way. Namely, set $a_{1} \triangleq 0$, and, for $m \geq 1$, define recursively the following sequence:

$$
a_{m+1} \triangleq \sum_{k=1}^{m}\left\|u_{m}\right\|_{\mathrm{CC}}, \quad I_{m} \triangleq\left[a_{m}, a_{m+1}\right], \quad I \triangleq \bigcup_{m=1}^{\infty} I_{m}
$$

It is clear that $\left|I_{m}\right|=\left\|u_{m}\right\|_{\text {CC }}$ and $I$ is a compact interval since the sequence $\left\{\left\|u_{m}\right\|_{\mathrm{CC}}: m \geq\right.$ $1\}$ is summable according to (3.23). We also define a family of functions $\left\{\tilde{h}_{m}, m \geq 1\right\}$ by

$$
\begin{equation*}
\tilde{h}_{m}(t) \triangleq h_{m}\left(\frac{t-a_{m}}{a_{m+1}-a_{m}}\right), \quad t \in I_{m} \tag{3.27}
\end{equation*}
$$

and the concatenation of the first $\tilde{h}_{m}$ 's is

$$
\begin{equation*}
\tilde{h}^{(m)} \triangleq \tilde{h}_{1} \sqcup \cdots \sqcup \tilde{h}_{m}:\left[0, a_{m+1}\right] \rightarrow \mathbb{R}^{d} \tag{3.28}
\end{equation*}
$$

We will now bound the derivative of $\tilde{h}_{m}$. Specifically, we first use equation (3.27) to get

$$
\sup _{m \geq 1}\left\|\dot{\tilde{h}}^{(m)}\right\|_{\infty ;\left[0, a_{m+1}\right]}=\sup _{m \geq 1}\left\|\dot{\tilde{h}}_{m}\right\|_{\infty ; I_{m}}=\sup _{m \geq 1} \frac{1}{\left|I_{m}\right|} \cdot\left\|\dot{h}_{m}\right\|_{\infty ;[0,1]}
$$

Then, resort to relation (3.15) which yields

$$
\sup _{m \geq 1}\left\|\dot{\tilde{h}}^{(m)}\right\|_{\infty ;\left[0, a_{m+1}\right]}=\sup _{m \geq 1}\left\{\frac{\left\|u_{m}\right\|_{\mathrm{CC}}}{\left|I_{m}\right|} \cdot\left\|\dot{\bar{h}}_{m}\right\|_{\infty ;[0,1]}\right\}
$$

Since $\left\|u_{m}\right\|_{\mathrm{CC}}=\left|I_{m}\right|$, we end up with

$$
\begin{equation*}
\sup _{m \geq 1}\left\|\dot{\tilde{h}}^{(m)}\right\|_{\infty ;\left[0, a_{m+1}\right]}=\sup _{m \geq 1}\left\{\left\|\dot{\bar{h}}_{m}\right\|_{\infty ;[0,1]}\right\} \leq C_{l_{0}} \tag{3.29}
\end{equation*}
$$

where the last inequality stems from the fact that $\left\|\ddot{\bar{h}}_{m}\right\|_{\infty ;[0,1]} \leq C_{l_{0}}$.
We can now proceed to the construction of the announced path joining $x$ and $y$. Namely, set

$$
\begin{equation*}
\tilde{h} \triangleq \bigsqcup_{m=1}^{\infty} \tilde{h}_{m}: I \rightarrow \mathbb{R}^{d} \tag{3.30}
\end{equation*}
$$

Then, according to (3.29), we have that $\tilde{h}$ is a smooth function from $I$ to $\mathbb{R}^{d}$. We claim that $\Phi_{1}(x ; \tilde{h})=y$, where $\Phi$ has to be understood in the sense of equation (1.4). Indeed, set

$$
z_{t}=\Phi_{t}(x ; \tilde{h}), \quad t \in I
$$

From the construction of $x_{m}$ in (3.17) and the fact that $\left.\tilde{h}\right|_{\left[0, a_{m+1}\right]}=\tilde{h}^{(m)}$, asserted in (3.30), we have

$$
\begin{equation*}
x_{m+1}=x+\sum_{\alpha=1}^{d} \int_{0}^{a_{m+1}} V_{\alpha}\left(z_{t}\right) d \tilde{h}_{t}^{\alpha} \tag{3.31}
\end{equation*}
$$

Since $x_{m+1} \rightarrow y$ as $m \rightarrow \infty$, which can be easily seen from (3.22), one can take limits in (3.31) to conclude that

$$
y=x+\sum_{\alpha=1}^{d} \int_{0}^{|I|} V_{\alpha}\left(z_{t}\right) d \tilde{h}_{t}^{\alpha} .
$$

Therefore, $\tilde{h}$ is a smooth path joining $x$ and $y$ in the sense of differential equations.
Step 4: Strategy for the upper bound. According to relation (3.3) in Lemma 3.4 on the scaling property, we have

$$
\begin{align*}
d(x, y) & =|I|^{H} d_{|I|}(x, y) \leq|I|^{H}\|\tilde{h}\|_{\overline{\mathcal{H}}([0,|I|])} \\
& =\lim _{m \rightarrow \infty}\left(\left(\sum_{k=1}^{m}\left|I_{k}\right|\right)^{H}\left\|\tilde{h}^{(m)}\right\|_{\overline{\mathcal{H}}\left(\left[0, a_{m+1}\right]\right)}\right), \tag{3.32}
\end{align*}
$$

where the last relation stems from the definition (3.30) of $\tilde{h}$.
In order to estimate the right-hand side of (3.32), we use the relation (2.4) for the CameronMartin norm to get

$$
\left\|\tilde{h}^{(m)}\right\|_{\overline{\mathcal{H}}\left(\left[0, a_{m+1}\right]\right)}^{2}=\left\|K^{-1} \tilde{h}^{(m)}\right\|_{L^{2}\left(\left[0, a_{m+1}\right] ; d t\right)}^{2}
$$

Using the definition of $K$, given by (2.3), we are led to

$$
\left\|\tilde{h}^{(m)}\right\|_{\tilde{\mathcal{H}}\left(\left[0, a_{m+1}\right]\right)}^{2}=C_{H} \int_{0}^{a_{m+1}}\left|t^{H-\frac{1}{2}} D_{0+}^{H-\frac{1}{2}}\left(s^{\frac{1}{2}-H} \dot{\tilde{h}}^{(m)}(s)\right)(t)\right|^{2} d t
$$

By the formula for the fractional derivative (cf. [15], equation (2.4)), we obtain

$$
\begin{aligned}
\left\|\tilde{h}^{(m)}\right\|_{\tilde{\mathcal{H}}\left(\left[0, a_{m+1}\right]\right)}^{2}= & C_{H} \cdot \int_{0}^{a_{m+1}} \left\lvert\, t^{H-\frac{1}{2}}\left(t^{1-2 H} \dot{\tilde{h}}^{(m)}(t)\right.\right. \\
& \left.+\left(H-\frac{1}{2}\right) \int_{0}^{t} \frac{t^{\frac{1}{2}-H} \dot{\tilde{h}}^{(m)}(t)-s^{\frac{1}{2}-H} \dot{\tilde{h}}^{(m)}(s)}{(t-s)^{H+\frac{1}{2}}} d s\right)\left.\right|^{2} d t
\end{aligned}
$$

We now split the interval $\left[0, a_{m+1}\right]$ as $\left[0, a_{m+1}\right]=\bigcup_{k=0}^{m} I_{k}$ and use the elementary inequality $(a+b+c)^{2} \leq 3\left(a^{2}+b^{2}+c^{2}\right)$ in order to get

$$
\begin{equation*}
\left\|\tilde{h}^{(m)}\right\|_{\overline{\mathcal{H}}\left(\left[0, a_{m+1}\right]\right)}^{2} \leq Q_{1}+Q_{2}+Q_{3} \tag{3.33}
\end{equation*}
$$

with

$$
\begin{align*}
& Q_{1} \triangleq \sum_{k=1}^{m} \int_{I_{k}}\left|t^{H-\frac{1}{2}}\left(t^{1-2 H} \dot{\tilde{h}}_{k}(t)\right)\right|^{2} d t \triangleq \sum_{k=1}^{m} Q_{1, k},  \tag{3.34}\\
& Q_{2} \triangleq \sum_{k=1}^{m} \int_{I_{k}}\left|t^{H-\frac{1}{2}} \sum_{l=1}^{k-1} \int_{I_{l}} \frac{t^{\frac{1}{2}-H} \dot{\tilde{h}}_{k}(t)-s^{\frac{1}{2}-H} \dot{\tilde{h}}_{l}(s)}{(t-s)^{H+\frac{1}{2}}} d s\right|^{2} d t \triangleq \sum_{k=1}^{m} Q_{2, k},  \tag{3.35}\\
& Q_{3} \triangleq \sum_{k=1}^{m} \int_{I_{k}}\left|t^{H-\frac{1}{2}} \int_{a_{k}}^{t} \frac{t^{\frac{1}{2}-H} \dot{\tilde{h}}_{k}(t)-s^{\frac{1}{2}-H} \dot{\tilde{h}}_{k}(s)}{(t-s)^{H+\frac{1}{2}}} d s\right|^{2} d t \triangleq \sum_{k=1}^{m} Q_{3, k} \tag{3.36}
\end{align*}
$$

We are now reduced to bound the above three terms. For the sake of conciseness we will mainly focus on $Q_{2}$ which is the most demanding in terms of singularities. We leave to the patient reader the nonrewarding task of checking details for $Q_{1}$ and $Q_{3}$ or refer to [14] for the complete details.
Step 5: Bound for $Q_{2}$. In order to estimate $Q_{2}$, we handle each $Q_{2, k}$ in (3.35) separately, and we resort to the elementary change of variables

$$
u \triangleq \frac{t-a_{k}}{a_{k+1}-a_{k}} \quad \text { and } \quad v \triangleq \frac{s-a_{l}}{a_{l+1}-a_{l}}
$$

We also express the terms $\dot{\tilde{h}}_{k}$ in (3.35) in terms of $\dot{\bar{h}}_{k}$. Thanks to some easy algebraic manipulations, we get

$$
\begin{equation*}
Q_{2, k}=\int_{0}^{1}\left|\sum_{l=1}^{k-1} \int_{0}^{1} \frac{\frac{\dot{h}_{k}(u)}{\left|I_{k}\right|}-\left(\frac{a_{k}+u\left|I_{k}\right|}{a_{l}+v\left|I_{l}\right|}\right)^{H-\frac{1}{2}} \cdot \frac{\dot{\vec{h}}_{l}(v)}{\left|I_{l}\right|}}{\left(a_{k}+u\left|I_{k}\right|-a_{l}-v\left|I_{l}\right|\right)^{H+\frac{1}{2}}}\right| I_{l}|d v|^{2}\left|I_{k}\right| d u . \tag{3.37}
\end{equation*}
$$

In the expression above, notice that, for $l \leq k-1$, we have

$$
a_{k}+u\left|I_{k}\right|-a_{l}-v\left|I_{l}\right|=q_{k, l}(u, v),
$$

where

$$
\begin{equation*}
q_{k, l}(u, v)=(1-v)\left|I_{l}\right|+\left|I_{l+1}\right|+\cdots+\left|I_{k-1}\right|+u\left|I_{k}\right| . \tag{3.38}
\end{equation*}
$$

Therefore, invoking the trivial bounds $a_{k}+u\left|I_{k}\right| \leq \sum_{j_{1}=1}^{k}\left|I_{j_{1}}\right|$ and $a_{l}+v\left|I_{l}\right| \geq \sum_{j_{2}=1}^{l-1}\left|I_{j_{2}}\right|$ and bounding trivially the differences by sums, we obtain

$$
\begin{equation*}
Q_{2, k} \leq C_{H} \int_{0}^{1}\left|\sum_{l=1}^{k-1} \int_{0}^{1} \frac{\left|\dot{\bar{h}}_{k}(u)\right|+\left(\frac{\sum_{j_{1}=1}^{k}\left|I_{j_{1}}\right|}{\sum_{j_{2}=1}^{l-1}\left|j_{j_{2}}\right|}\right)^{H-\frac{1}{2}} \cdot\left|\dot{\bar{h}}_{l}(v)\right|}{\left|q_{k, l}(u, v)\right|^{H+\frac{1}{2}}}\right| I_{l}|d v|^{2}\left|I_{k}\right| d u \tag{3.39}
\end{equation*}
$$

In order to obtain a sharp estimate in (3.39), we want to take advantage of the fact that $\dot{\bar{h}}_{l}$ is supported on $[1 / 3,2 / 3]$. This allows to avoid the singularities in $u, v$ close to 0 and 1 . We thus introduce the intervals

$$
J_{1} \triangleq[0,1 / 3], \quad J_{2} \triangleq[1 / 3,2 / 3], \quad J_{3} \triangleq[2 / 3,1]
$$

and decompose the expression (3.39) as follows:

$$
Q_{2, k} \leq C_{H} \sum_{p, q=1}^{3} L_{k, p, q}
$$

where the quantity $L_{k, p, q}$ is defined by

$$
\begin{equation*}
L_{k, p, q} \triangleq \int_{J_{p}}\left|\sum_{l=1}^{k-1} \int_{J_{q}} \frac{\left|\dot{\bar{h}}_{k}(u)\right|+\left(\frac{\left|I_{1}\right|+\cdots+\left|I_{k}\right|}{\left|I_{1}\right|+\cdots+\mid I_{l-1}}\right)^{H-\frac{1}{2}} \cdot\left|\dot{\bar{h}}_{l}(v)\right|}{\left|q_{k, l}(u, v)\right|^{H+\frac{1}{2}}}\right| I_{l}|d v|^{2}\left|I_{k}\right| d u, \tag{3.40}
\end{equation*}
$$

for all $p, q=1,2,3$. Notice again that, since all the $\dot{\bar{h}}_{k}$ are supported on $[1 / 3,2 / 3]$, the only nonvanishing $L_{k, p, q}$ 's are those for which $p=2$ or $q=2$. Let us show how to handle the terms $L_{k, p, q}$, given by (3.40), according to $q=1,2$ and $q=3$.

Whenever $q=1$ or $q=2$, regardless of the value of $p$, it is easily seen from (3.38) that we can bound $q_{k, l}(u, v)$ from below uniformly by $C \sum_{j=l}^{k-1}\left|I_{j}\right|$. Thanks again to the fact that $\dot{\bar{h}}_{k}$ is uniformly bounded for all $k$, we obtain

$$
\begin{aligned}
& \frac{\left|\dot{\bar{h}}_{k}(u)\right|+\left(\frac{\left|I_{1}\right|+\cdots+\left|I_{k}\right|}{\left|I_{1}\right|+\cdots+\left|I_{l-1}\right|}\right)^{H-\frac{1}{2}} \cdot\left|\dot{\bar{h}}_{l}(v)\right|}{q_{k, l}(u, v)^{H+\frac{1}{2}}} \\
& \quad \leq \frac{C_{H, l_{0}}}{\left(\left|I_{l}\right|+\cdots+\left|I_{k-1}\right|\right)^{H+\frac{1}{2}}} \cdot\left(\frac{\left|I_{1}\right|+\cdots+\left|I_{k}\right|}{\left|I_{1}\right|+\cdots+\left|I_{l-1}\right|}\right)^{H-\frac{1}{2}} .
\end{aligned}
$$

Summing the above quantity over $l$ and integrating over [0, 1], we end up with

$$
L_{k, p, q} \leq C_{H, l_{0}}\left|I_{k}\right| \cdot\left(\sum_{l=1}^{k-1} \frac{\left|I_{l}\right|}{\left(\left|I_{l}\right|+\cdots+\left|I_{k-1}\right|\right)^{H+\frac{1}{2}}} \cdot\left(\frac{\left|I_{1}\right|+\cdots+\left|I_{k}\right|}{\left|I_{1}\right|+\cdots+\left|I_{l-1}\right|}\right)^{H-\frac{1}{2}}\right)^{2}
$$

By lower bounding the quantity $\left|I_{1}\right|+\cdots+\left|I_{l-1}\right|$ above uniformly by $\left|I_{1}\right|$, we get

$$
\begin{equation*}
L_{k, p, q} \leq C_{H, l_{0}}\left|I_{k}\right| \cdot\left(\frac{\sum_{j=1}^{k}\left|I_{j}\right|}{\left|I_{1}\right|}\right)^{2 H-1} \cdot\left(\sum_{l=1}^{k-1}\left|I_{l}\right|^{\frac{1}{2}-H}\right)^{2} . \tag{3.41}
\end{equation*}
$$

Recall that we have shown in (3.26) that $m \mapsto\left\|u_{m}\right\|_{\text {CC }}$ is a decreasing sequence. Since $\left\|u_{m}\right\|_{\mathrm{CC}}=\left|I_{m}\right|$, we can bound uniformly $\sum_{l=1}^{k-1}\left|I_{l}\right|^{1 / 2-H}$ by $k\left|I_{1}\right|^{1 / 2-H}$ and $\left|I_{1}\right|^{-1} \sum_{j=1}^{k}\left|I_{j}\right|$ by $k$. Plugging this information into (3.41), we obtain,

$$
\begin{equation*}
L_{k, p, q} \leq C_{H, l_{0}} k^{2 H+1}\left|I_{k}\right|^{2(1-H)} \tag{3.42}
\end{equation*}
$$

which is our bound for $L_{k, p, q}$ when $q \in\{1,2\}$.
Let us now bound $L_{k, p, q}$ for $q=3$ and $p=2$. In this case, going back to the definition (3.40) of $L_{k, p, q}$, we have that $\dot{\bar{h}}_{l}(v)=0$ for $v \in J_{q}$. Thus, we get

$$
\begin{align*}
L_{k, 2,3} & =\int_{J_{2}}\left|\sum_{l=1}^{k-1} \int_{J_{3}} \frac{\left|\dot{\bar{h}}_{k}(u)\right|\left|I_{l}\right| d v}{\left((1-v)\left|I_{l}\right|+\left|I_{l+1}\right|+\cdots+\left|I_{k-1}\right|+u\left|I_{k}\right|\right)^{H+\frac{1}{2}}}\right|^{2}\left|I_{k}\right| d u  \tag{3.43}\\
& \leq C_{H, l_{0}}\left|\sum_{l=1}^{k-1} \int_{J_{3}} \frac{\left|I_{l}\right| d v}{\left((1-v)\left|I_{l}\right|+\left|I_{l+1}\right|+\cdots+\left|I_{k}\right|\right)^{H+\frac{1}{2}}}\right|^{2} \cdot\left|I_{k}\right|,
\end{align*}
$$

where we have used the boundedness of $\dot{\bar{h}}_{k}$ for the second inequality. We can now evaluate the above $v$-integral explicitly, which yields

$$
\begin{aligned}
& \int_{J_{3}} \frac{\left|I_{l}\right| d v}{\left((1-v)\left|I_{l}\right|+\left|I_{l+1}\right|+\cdots+\left|I_{k}\right|\right)^{H+\frac{1}{2}}} \\
& \quad=\frac{1}{\left(H-\frac{1}{2}\right)}\left(\frac{1}{\left(\left|I_{l+1}\right|+\cdots+\left|I_{k}\right|\right)^{H-\frac{1}{2}}}-\frac{1}{\left(\frac{1}{3}\left|I_{l}\right|+\left|I_{l+1}\right|+\cdots+\left|I_{k}\right|\right)^{H-\frac{1}{2}}}\right) \leq \frac{C_{H}}{\left|I_{k}\right|^{H-\frac{1}{2}}},
\end{aligned}
$$

where the second inequality is obtained by lower bounding trivially $\left|I_{l+1}+\cdots+\left|I_{k}\right|\right.$ by $| I_{k} \mid$. Summing this inequality over $l$ and plugging this information into (3.43), we get

$$
\begin{equation*}
L_{k, 2,3} \leq C_{H, l_{0}}\left|I_{k}\right|\left(\frac{k}{\left|I_{k}\right|^{H-\frac{1}{2}}}\right)^{2} \leq C_{H, l_{0}} k^{2 H+1}\left|I_{k}\right|^{2(1-H)} \tag{3.44}
\end{equation*}
$$

Summarizing our considerations in this step, we have handled the cases $q=1,2$ and $(q, p)=$ $(3,2)$ in (3.42) and (3.44), respectively. Therefore, we obtain

$$
\begin{equation*}
Q_{2} \leq C_{H, l_{0}} \sum_{k=1}^{m} k^{2 H+1}\left|I_{k}\right|^{2(1-H)} \tag{3.45}
\end{equation*}
$$

Step 6: Conclusion. Let us go back to the decomposition (3.33), plug in our bound (3.45) on $Q_{2}$ and recall that similar bounds are available for $Q_{1}$ and $Q_{3}$. We get

$$
\left\|\tilde{h}^{(m)}\right\|_{\tilde{\mathcal{H}}\left(\left[0, a_{m+1}\right]\right)}^{2} \leq C_{H}\left(Q_{1}+Q_{2}+Q_{3}\right) \leq C_{H, l_{0}} \sum_{k=1}^{m} k^{2 H+1}\left|I_{k}\right|^{2(1-H)}
$$

In addition, we have $\left|I_{k}\right|=\left\|u_{k}\right\|_{\mathrm{CC}}$, and relation (3.23) asserts that $k \mapsto\left\|u_{k}\right\|_{\mathrm{CC}}$ decays exponentially. Thus, we get

$$
\begin{align*}
& \left(\sum_{k=1}^{m}\left|I_{k}\right|\right)^{2 H}\left\|\tilde{h}^{(m)}\right\|_{\tilde{\mathcal{H}}\left(\left[0, a_{m+1}\right]\right)}^{2} \\
& \quad \leq C_{H, l_{0}}\left(\sum_{k=1}^{m}\left|I_{k}\right|\right)^{2 H}\left(\sum_{k=1}^{m} k^{2 H+1}\left|I_{k}\right|^{2(1-H)}\right)  \tag{3.46}\\
& \quad \leq C_{H, V, l_{0}}\left(\sum_{k=1}^{m} 2^{-\frac{k}{l_{0}}}\right)^{2 H} \cdot\left(\sum_{k=1}^{m} k^{2 H+1} 2^{-\frac{2(1-H)}{l_{0}} k}\right) \cdot|x-y|^{\frac{2}{l_{0}}} \\
& \quad \leq C_{H, V, l_{0}}|x-y|^{\frac{2}{\left.\right|_{0}}}
\end{align*}
$$

where we have trivially bounded the partial geometric series for the last step. Moreover, the left-hand side of (3.46) converges to a quantity which is lower bounded by $d^{2}(x, y)$ as $m \rightarrow \infty$, thanks to (3.32). Therefore, letting $m \rightarrow \infty$ in (3.46), we have obtained

$$
\begin{equation*}
d(x, y)^{2} \leq C_{H, V, l_{0}}|x-y|^{\frac{2}{l_{0}}} \tag{3.47}
\end{equation*}
$$

which concludes our proof of Theorem 1.2.
3.3. Proof of Theorem 1.3. The technique developed in the previous proof also provides part of the essential analysis for proving Theorem 1.3, namely, the local Lipschitz equivalence of all control distance functions with different Hurst parameters. Our main idea for proving Theorem 1.3 is to show that the control distance function $d(x, y)$ is locally Lipschitz equivalent to the "distance" function defined by

$$
\begin{equation*}
g(x, y) \triangleq \inf \left\{\|u\|_{\mathrm{CC}}: u \in \mathfrak{g}^{\left(l_{0}\right)} \text { and } x+F_{l_{0}}(u, x)=y\right\} \tag{3.48}
\end{equation*}
$$

Note that $g(x, y)$ is canonical in the sense that it does not depend on the Hurst parameter $H$.
The following lemma leads to one direction of the comparison.
Lemma 3.7. There exist constants $C, \kappa>0$ such that, for any $u \in \mathfrak{g}^{\left(l_{0}\right)}$ with $\|u\|_{\mathrm{HS}}<\kappa$, we have

$$
\begin{equation*}
d\left(x, x+F_{l_{0}}(u, x)\right) \leq C\|u\|_{\mathrm{CC}} \tag{3.49}
\end{equation*}
$$

Proof. We only consider the case when $H>1 / 2$, as the other case follows from Lemma 2.8 in [15] and the result for the diffusion case proved in [20]. We use the same notation as in the proof of Theorem 1.2. In particular, we set up an inductive procedure as in that proof, starting by setting $u_{1} \triangleq u, x_{1} \triangleq x$ and $y \triangleq x+F_{l_{0}}(u, x)$. Choose $\kappa_{1}>0$ so that

$$
\begin{equation*}
\|u\|_{\mathrm{CC}}<\kappa_{1} \quad \Longrightarrow \quad|y-x|<\delta \tag{3.50}
\end{equation*}
$$

where $\delta$ is the constant arising in the proof of Theorem 1.2. By constructing successively elements $u_{m} \in \mathfrak{g}^{(l)}$ and intervals $I_{m}$, we obtain exactly, as in (3.24), that

$$
\begin{equation*}
\left|I_{m}\right|=\left\|u_{m}\right\|_{\mathrm{CC}} \leq C_{V, l_{0}}\left\|u_{m-1}\right\|^{1+\frac{1}{l_{0}}}=C_{V, l_{0}}\left|I_{m-1}\right|^{1+\frac{1}{l_{0}}} \tag{3.51}
\end{equation*}
$$

In addition, along the same lines as (3.46) and (3.47) we have

$$
\begin{equation*}
d\left(x, x+F_{l_{0}}(u, x)\right)^{2} \leq C_{H, l_{0}} \lim _{m \rightarrow \infty}\left(\sum_{k=1}^{m}\left|I_{k}\right|\right)^{2 H}\left(\sum_{k=1}^{m} k^{2 H+1}\left|I_{k}\right|^{2(1-H)}\right) \tag{3.52}
\end{equation*}
$$

We now estimate the right-hand side of (3.52) in a slightly different way from the previous step. Namely, let us set $\alpha \triangleq 1+1 / l_{0}$. By iterating (3.51), we obtain that

$$
\left|I_{m}\right| \leq\left(C_{V, l_{0}}\left|I_{1}\right|\right)^{\alpha^{m-1}} \quad \forall m \geq 1
$$

Therefore, we can bound the two terms on the right-hand side of (3.52) as follows:

$$
\begin{equation*}
\sum_{k=1}^{m}\left|I_{k}\right| \leq C_{V, l_{0}}\left|I_{1}\right| \cdot\left(\sum_{k=1}^{m}\left(C_{V, l_{0}}\left|I_{1}\right|\right)^{\alpha^{k-1}-1}\right) \tag{3.53}
\end{equation*}
$$

and

$$
\begin{equation*}
\sum_{k=1}^{m} k^{2 H+1}\left|I_{k}\right|^{2(1-H)} \leq\left(C_{V, l_{0}}\left|I_{1}\right|\right)^{2(1-H)} \cdot\left(\sum_{k=1}^{m} k^{2 H+1}\left(C_{V, l_{0}}\left|I_{1}\right|\right)^{2(1-H)\left(\alpha^{k-1}-1\right)}\right) \tag{3.54}
\end{equation*}
$$

To estimate $\left|I_{1}\right|$, recall that $u_{1}=u$. We further choose $\kappa_{2}>0$ so that

$$
\begin{equation*}
\|u\|_{\mathrm{CC}}<\kappa_{2} \quad \Longrightarrow \quad C_{V, l_{0}}\left|I_{1}\right|=C_{V, l_{0}}\left\|u_{1}\right\|_{\mathrm{CC}} \leq \frac{1}{2} \tag{3.55}
\end{equation*}
$$

By taking $\kappa=\kappa_{1} \wedge \kappa_{2}$, we can assume that both (3.50) and (3.55) hold. Also, note that both of the series

$$
\begin{equation*}
\sum_{m=1}^{\infty}\left(\frac{1}{2}\right)^{\alpha^{m-1}-1} \quad \text { and } \quad \sum_{m=1}^{\infty} m^{2 H+1}\left(\frac{1}{2}\right)^{2(1-H)\left(\alpha^{m-1}-1\right)} \tag{3.56}
\end{equation*}
$$

are convergent. Therefore, by plugging (3.56) and (3.55) and then (3.53) and (3.54) into (3.52), we have

$$
d\left(x, x+F_{l_{0}}(u, x)\right)^{2} \leq C_{H, V, l_{0}}\left|I_{1}\right|^{2 H} \cdot\left|I_{1}\right|^{2(1-H)}=C_{H, V, l_{0}}\left|I_{1}\right|^{2}=C_{H, V, l_{0}}\|u\|_{\mathrm{CC}}^{2}
$$

provided that $\|u\|_{\mathrm{CC}}<\kappa$. Our result (3.49) thus follows.
It is an immediate consequence of Lemma 3.7 that

$$
d(x, y) \leq C g(x, y)
$$

provided that $|x-y|$ is small so that the infimum in (3.48) can be taken over those $u$ 's with $\|u\|_{\mathrm{CC}}<\kappa$. To complete the proof of Theorem 1.3, it remains to establish the other direction of the above inequality. We now provide the remaining details by using some pathwise estimates from rough path theory.

Proof of Theorem 1.3. Suppose that $x, y \in \mathbb{R}^{N}$, and let $h \in \Pi_{x, y}$ be a CameronMartin path that joins $x$ to $y$ in the sense of differential equations such that $\|h\|_{\overline{\mathcal{H}}} \leq 2 d(x, y)$. Set $u \triangleq \log S_{l_{0}}(h) \in \mathfrak{g}^{\left(l_{0}\right)}$. Note that $u$ does not join $x$ to $y$ in the sense of Taylor approximation. However, we can use the map $\Psi_{l_{0}}$ to find a $w \in \mathfrak{g}^{\left(l_{0}\right)}$ which does this, and, more precisely, $w$ is given by

$$
w \triangleq \Psi_{l_{0}}\left(u, x, y-x-F_{l_{0}}(u, x)\right)
$$

When $|x-y|$ is small, the quantities $d(x, y),\|h\|_{\overline{\mathcal{H}}}, u$ are all small and $w$ is well defined.
According to Proposition 2.9, we have

$$
\left|\|w\|_{\mathrm{CC}}-\|u\|_{\mathrm{CC}}\right| \leq \rho_{\mathrm{CC}}(w, u) \leq C_{1}\|w-u\|_{\mathrm{HS}}^{1 / l_{0}},
$$

and, thus,

$$
\|w\|_{\mathrm{CC}} \leq C_{1}\|w-u\|_{\mathrm{HS}}^{1 / l_{0}}+\|u\|_{\mathrm{CC}}
$$

We now estimate the two terms on the right-hand side separately.
For the first term, according to Lemma 3.1 (ii), we have

$$
\|w-u\|_{\mathrm{HS}} \leq A\left|y-x-F_{l_{0}}(u, x)\right|=A\left|\Phi_{1}(x ; h)-x-F_{l_{0}}(u, x)\right|
$$

with some constant $A$. By the rough path estimates of [13], Corollary 10.15, and Proposition 2.3, we know that

$$
\left|\Phi_{1}(x ; h)-x-F_{l_{0}}(u, x)\right| \leq C_{V, l_{0}}\|h\|_{q-\mathrm{var}}^{\bar{I}_{0}} \leq C_{H, V, l_{0}}\|h\|_{\overline{\mathcal{H}}}^{\bar{L}_{0}},
$$

where $q \in[1,2)$ and $\bar{l}_{0}$ is an arbitrary number in $\left(l_{0}, l_{0}+1\right)$. Therefore, we have

$$
\|w-u\|_{\mathrm{HS}} \leq A C_{H, V, l_{0}}\|h\|_{\overline{\mathcal{H}}}^{\overline{\overline{\mathcal{H}}}_{0}}
$$

For the second term, we claim that $\|u\|_{\mathrm{CC}} \leq C_{H, l_{0}}\|h\|_{\overline{\mathcal{H}}}$. Indeed, recall that $u=$ $\log \left(S_{l_{0}}(h)\right)$, and set $S_{l_{0}}(h)=g$. Since $h \in C^{q-\mathrm{var}}$ with $q \in[1,2)$, Lyons' extension theorem (cf. [22], Theorem 2.2.1) implies that, for all $i=1, \ldots, l_{0}$, we have

$$
\left\|g_{i}\right\|_{\mathrm{HS}} \leq C_{H, l_{0}}\|h\|_{q-\mathrm{var}}^{i},
$$

where $g_{i}$ is the $i$ th component of $g$. If we define the homogeneous norm $\left\|\|\cdot\| \mid\right.$ on $G^{\left(l_{0}\right)}$ by

$$
\|\xi\|\left\|\max _{1 \leq i \leq l_{0}}\right\| \xi_{i} \|_{\mathrm{HS}}^{1 / i}, \xi \in G^{\left(l_{0}\right)}
$$

we get the following estimate:

$$
\begin{equation*}
\|g\|\left\|\leq C_{H, l_{0}}\right\| h \|_{q-\mathrm{var}} \quad \text { and } \quad\|u\|_{\mathrm{CC}} \leq C_{H, l_{0}}\|g\| \|, \tag{3.57}
\end{equation*}
$$

where the second inequality stems from the equivalence of homogeneous norms in $G^{\left(l_{0}\right)}$ (cf. [13], Theorem 7.44). Now, combining the two inequalities in (3.57) and the variation estimate for Cameron-Martin paths (cf. [12], Corollary 1), we end up with

$$
\begin{equation*}
\|u\|_{\mathrm{CC}} \leq C_{H, l_{0}}\|h\|_{q-\mathrm{var}} \leq C_{H, l_{0}}\|h\|_{\overline{\mathcal{H}}} \tag{3.58}
\end{equation*}
$$

Therefore, we arrive at

$$
\|w\|_{\mathrm{CC}} \leq C_{H, V, l_{0}} \cdot\left(\|h\|_{\overline{\mathcal{H}}}^{\bar{l}_{0} / l_{0}}+\|h\|_{\overline{\mathcal{H}}}\right) \leq C_{H, V, l_{0}}\|h\|_{\overline{\mathcal{H}}}
$$

By the fact that $w$ joins $x$ to $y$ in the sense of Taylor approximation as well as the choice of $h$, we conclude that

$$
g(x, y) \leq C_{H, V, l_{0}} d(x, y)
$$

This completes the proof of Theorem 1.3.

REMARK 3.8. The Cameron-Martin space $\overline{\mathcal{H}}$ and control distance function $d_{H}(x, y)$ are well defined, even when $H \leq 1 / 4$. A careful examination of the proofs shows that Theorem 1.2 and Theorem 1.3 remain valid in this case. We do not emphasise this in the main context since the SDE (1.1) does not make sense in this case.
4. Local lower estimate for the density of solution. In this section we develop the proof of Theorem 1.6 under the uniform hypoellipticity assumption (1.2). Comparing with the elliptic case in [15], one faces a much more complex situation here. More specifically, the deterministic Malliavin covariance matrix of $X_{t}^{x}$ will not be uniformly nondegenerate (i.e., Lemma 3.6 in [15] is no longer true). Without this key ingredient the entire elliptic argument will break down, and one needs new approaches. Our strategy follows the main philosophy of Kusuoka-Stroock [20] in the diffusion case. However, there are nontrivial challenges in several key steps for the fractional Brownian setting which require new ideas and methods.

To increase readability, we first summarize the main strategy of the proof. Our analysis starts from the existence of the truncated signature of order $l$ for the fractional Brownian motion (cf. Proposition 2.10). Specifically, with our notation (2.11) in mind we write

$$
\begin{equation*}
\Gamma_{t} \triangleq S_{l}(B)_{0, t}=1+\sum_{i=1}^{l} \int_{0<t_{1}<\cdots<t_{i}<t} d B_{t_{1}} \otimes \cdots \otimes d B_{t_{i}} \tag{4.1}
\end{equation*}
$$

In the sequel we will also use the truncated $\mathfrak{g}^{(l)}$-valued log-signature of $B$, defined by

$$
\begin{equation*}
U_{t}^{(l)} \triangleq \log S_{l}(B)_{0, t} \tag{4.2}
\end{equation*}
$$

Notice that $U_{t}^{(l)}$ features in relation (2.17), and, more precisely, the process

$$
\begin{equation*}
X_{l}(t, x) \triangleq x+F_{l}\left(U_{t}^{(l)}, x\right) \tag{4.3}
\end{equation*}
$$

is the Taylor approximation of order $l$ for the solution of the rough equation (1.1) in small time (cf. relation (2.18)).

With those preliminary notation in hand, we decompose the strategy toward the proof of Theorem 1.6 into three major steps:

Step one. According to the scaling property of fractional Brownian motion, a precise local lower estimate on the density of $U_{t}^{(l)}$ can be easily obtained from a general positivity property.
Step two. When $l \geq l_{0}$, the hypoellipticity of the vector fields allows us to obtain a precise local lower estimate on the density of the process $X_{l}(t, x)$, defined by (4.3) from the estimate on $U_{t}^{(l)}$ derived in step one.
Step three. When $t$ is small, the density of $X_{l}(t, x)$ is close to the density of the actual solution in a reasonable sense, and the latter inherits the lower estimate obtained in step two.

The above philosophy was first proposed by Kusuoka-Stroock [20] in the diffusion case. However, in the fractional Brownian setting, there are several difficulties when implementing these steps precisely. Conceptually, the main challenge arises from the need of respecting the fractional Brownian scaling and the Cameron-Martin structure in each step in order to obtain sharp estimates. More specifically, for step one we need a new idea to prove the positivity for the density of $U_{t}^{(l)}$ when the Markov property is not available. For step two we rely on Theorem 1.3 that we have proven in the Section 3.3 which yields sharp estimates for the density of $X_{l}(t, x)$. In step three a new ingredient is needed to prove uniformity for an upper estimate for the density of $X_{l}(t, x)$ with respect to the degree $l$ of expansion. In the following sections, we develop the above three steps mathematically.
4.1. Step one: Local lower estimate for the signature density of fractional Brownian motion. We fix $l \geq 1$. Recall that the truncated signature process $\Gamma_{t}$ is defined by (4.1). Let $\left\{\mathrm{e}_{1}, \ldots, \mathrm{e}_{d}\right\}$ be the standard basis of $\mathbb{R}^{d}$. By viewing this family as vectors in $\mathfrak{g}^{(l)} \cong T_{1} G^{(l)}$, we denote the associated left invariant vector fields on $G^{(l)}$ by $\left\{\tilde{W}_{1}, \ldots, \tilde{W}_{d}\right\}$. It is standard (cf. [13], Remark 7.43) that $\Gamma_{t}$ satisfies the following intrinsic rough differential equation on $G^{(l)}$ :

$$
\left\{\begin{array}{l}
d \Gamma_{t}=\sum_{\alpha=1}^{d} \tilde{W}_{\alpha}\left(\Gamma_{t}\right) d B_{t}^{\alpha},  \tag{4.4}\\
\Gamma_{0}=\mathbf{1}
\end{array}\right.
$$

Let $U_{t} \triangleq \log \Gamma_{t} \in \mathfrak{g}^{(l)}$ be the truncated log-signature path, as defined in (4.2). Since $\left\{\tilde{W}_{1}, \ldots, \tilde{W}_{d}\right\}$ satisfies Hörmander's condition by the definition of $\mathfrak{g}^{(l)}$, we know that $U_{t}$ admits a smooth density with respect to the Lebesgue measure $d u$ on $\mathfrak{g}^{(l)}$. Denote this density by $\rho_{t}(u)$.

A key ingredient in this step is to show that the density function $\rho_{t}$ is everywhere strictly positive. In the Brownian case, this was proved in [20] by using support theorem and the Markov property. In the fractional Brownian setting, the argument breaks down, although general support theorems for Gaussian rough paths are still available. It turns out that there is a simple neat proof based on Sard's theorem and a general positivity criterion established Baudoin-Nualart-Ouyang-Tindel [3]. We mention that Baudoin-Feng-Ouyang [1] also contains a proof of this property using a different approach.

We first recall the classical Sard's theorem, and we refer the reader to [24] for a beautiful presentation. Let $f: M \rightarrow N$ be a smooth map between two finite dimensional differentiable manifolds $M$ and $N$. A point $x \in M$ is said to be a critical point of $f$ if the differential $d f_{x}: T_{x} M \rightarrow T_{f(x)} N$ is not surjective. A critical value of $f$ in $N$ is the image of a critical point in $M$. Also, recall that a subset $E \subseteq N$ is a Lebesgue null set if its intersection with any coordinate chart has zero Lebesgue measure in the corresponding coordinate space.

THEOREM 4.1 (Sard's theorem). Let $f: M \rightarrow N$ be a smooth map between two finite dimensional differentiable manifolds. Then, the set of critical values of $f$ is a Lebesgue null set in $N$.

We now prove the positivity result announced above which will be important for our future considerations.

LEmmA 4.2. For each $t>0$, the density $\rho_{t}$ of the truncated signature path $U_{t}$ is everywhere strictly positive.

Proof. We only consider the case when $t=1$. The general case follows from the scaling property (4.11) below. Our strategy relies on the fact that $\Gamma_{t}=\exp \left(U_{t}\right)$ solves equation (4.4). In addition, recall our Convention 3.6 about the identification of $\mathfrak{g}^{(l)}$ and $G^{(l)}$. Therefore, we can get the desired positivity by applying [3], Theorem 1.4. To this aim, recall that the standing assumptions in [3], Theorem 1.4, are the following:
(i) The Malliavin covariance matrix of $U_{t}$ is invertible with inverse in $L^{p}(\Omega)$ for all $p>1$;
(ii) for any $u \in \mathfrak{g}^{(l)}$, there exists $h \in \overline{\mathcal{H}}$ such that $\log S_{l}(h)=u$ and

$$
\begin{equation*}
\left(d \log S_{l}\right)_{h}: \overline{\mathcal{H}} \rightarrow \mathfrak{g}^{(l)} \quad \text { is surjective }, \tag{4.5}
\end{equation*}
$$

where $S_{l}(h) \triangleq S_{l}(h)_{0,1}$ is the truncated map.

Notice that item (i) is proved in [1], Theorem 3.3. We will thus focus on condition (ii) in the remainder of the proof.

In order to prove relation (4.5) in item (ii) above, let us introduce some additional notation. First, we shall write $G \triangleq G^{(l)}$ for the sake of simplicity. Then, for all $n \geq 1$, we introduce a linear map $H_{n}:\left(\mathbb{R}^{d}\right)^{n} \rightarrow \overline{\mathcal{H}}$ in the following way. Given $y=\left(y_{1}, \ldots, y_{n}\right)$, the function $H_{n}(y)$ is defined to be the piecewise linear path obtained by concatenating the vectors $y_{1}, \ldots, y_{n}$ successively. We also define a set $\overline{\mathcal{H}}_{0}$ of piecewise linear paths by

$$
\overline{\mathcal{H}}_{0} \triangleq \bigcup_{n=1}^{\infty} H_{n}\left(\left(\mathbb{R}^{d}\right)^{n}\right) \subseteq \overline{\mathcal{H}}
$$

Note that $\overline{\mathcal{H}}_{0}$ is closed under concatenation, and $S_{l}\left(\overline{\mathcal{H}}_{0}\right)=G$ by the Chow-Rashevskii theorem (cf. Remark 2.8). Now, we claim that
$(\mathbf{P})$ for any $g \in G$, there exists $h \in \overline{\mathcal{H}}_{0}$ such that $S_{l}(h)=g$ and the differential $\left.\left(d S_{l}\right)_{h}\right|_{\overline{\mathcal{H}}_{0}}$ : $\overline{\mathcal{H}}_{0} \rightarrow T_{g} G$ is surjective.
The property $(\mathbf{P})$ is clearly stronger than the original desired claim (4.5). To prove $(\mathbf{P})$, let $\mathcal{P}$ be the set of elements in $G$ which satisfy $(\mathbf{P})$. We first show that $\mathcal{P}$ is either $\varnothing$ or $G$. The main idea behind our strategy is that if there exists $g_{0} \in \mathcal{P}$ such that $\left(d S_{l}\right)_{h_{0}}$ is a submersion for some $h_{0} \in \overline{\mathcal{H}}_{0}$ satisfying $S_{l}\left(h_{0}\right)=g_{0}$, then one can obtain every point $g \in G$ by a left translation $L_{a}$, since $d L_{a}$ is an isomorphism. To be more precise, suppose that $g_{0} \in G$ is an element satisfying $(\mathbf{P})$. By definition, there exists a path $h_{0} \in \overline{\mathcal{H}}_{0}$ such that $S_{l}\left(h_{0}\right)=g_{0}$ and $\left.\left(d S_{l}\right)_{h_{0}}\right|_{\mathcal{H}_{0}}$ is surjective. Now, pick a generic element $a \in G$, and choose a path $\alpha \in \overline{\mathcal{H}}_{0}$ so that $S_{l}(\alpha)=a$. Then, $S_{l}\left(\alpha \sqcup h_{0}\right)=a \otimes g_{0}$. We want to show that $\left(d S_{l}\right)_{\alpha \sqcup h_{0}}: \overline{\mathcal{H}}_{0} \rightarrow T_{a \otimes g_{0}} G$ is surjective. For this purpose, let $\xi \in T_{a \otimes g_{0}} G$, and set

$$
\xi_{0} \triangleq d L_{a^{-1}}(\xi) \in T_{g_{0}} G
$$

By the surjectivity of $\left.\left(d S_{l}\right)_{h_{0}}\right|_{\overline{\mathcal{H}}_{0}}$, there exists $\gamma \in \overline{\mathcal{H}}_{0}$ such that $\left(d S_{l}\right)_{h_{0}}(\gamma)=\xi_{0}$. It follows that, for $\varepsilon>0$, we have

$$
S_{l}\left(\alpha \sqcup\left(h_{0}+\varepsilon \cdot \gamma\right)\right)=a \otimes S_{l}\left(h_{0}+\varepsilon \cdot \gamma\right)
$$

By differentiation with respect to $\varepsilon$ at $\varepsilon=0$, we obtain that

$$
\left(d S_{l}\right)_{\alpha \sqcup h_{0}}(0 \sqcup \gamma)=\left(d L_{a}\right)_{S_{l}\left(h_{0}\right)} \circ\left(d S_{l}\right)_{h_{0}}(\gamma)=\left(d L_{a}\right)_{g_{0}}\left(\xi_{0}\right)=\xi
$$

Therefore, $\left.\left(d S_{l}\right)_{\alpha \sqcup h_{0}}\right|_{\overline{\mathcal{H}}_{0}}$ is surjective. Since $a$ is arbitrary, we conclude that if $\mathcal{P}$ is nonempty, then $\mathcal{P}=G$.

To complete the proof, it remains to show that $\mathcal{P} \neq \varnothing$. This will be a simple consequence of Sard's theorem. Indeed, for each $n \geq 1$, define

$$
\begin{equation*}
f_{n} \triangleq S_{l} \circ H_{n}:\left(\mathbb{R}^{d}\right)^{n} \rightarrow G \tag{4.6}
\end{equation*}
$$

where we recall that $H_{n}(y)$ is the piecewise linear path obtained by concatenating $y_{1}, \ldots, y_{n}$. The map $f_{n}$ is simply given by

$$
f_{n}\left(y_{1}, \ldots, y_{n}\right)=\exp \left(y_{1}\right) \otimes \cdots \otimes \exp \left(y_{n}\right)
$$

where we recall that the exponential maps is defined by (2.6). It is readily checked that $f_{n}$ is a smooth map. According to Sard's theorem (cf. Theorem 4.1), the set of critical values of $f_{n}$, denoted as $E_{n}$, is a Lebesgue null set in $G$. It follows that $E \triangleq \bigcup_{n=1}^{\infty} E_{n}$ is also a Lebesgue null set in $G$. We have thus obtained that,

$$
G \backslash E=\left(\bigcup_{n=1}^{\infty} f_{n}\left(\left(\mathbb{R}^{d}\right)^{n}\right)\right) \backslash E \neq \varnothing
$$

where the first equality is due to the fact that $S_{l}\left(\overline{\mathcal{H}}_{0}\right)=G$ by the Chow-Rashevskii theorem. Pick any element $g \in G \backslash E$. Then, for some $n \geq 1$, we have $g \in f_{n}\left(\left(\mathbb{R}^{d}\right)^{n}\right) \backslash E_{n}$. In particular, there exists $y \in\left(\mathbb{R}^{d}\right)^{n}$ such that $f_{n}(y)=g$ and $\left(d f_{n}\right)_{y}$ is surjective. We claim that $g \in \mathcal{P}$ with $h \triangleq H_{n}(y) \in \overline{\mathcal{H}}_{0}$ being the associated path. Indeed, it is apparent that $S_{l}(h)=g$. In addition, let $\xi \in T_{g} G$ and $w \in\left(\mathbb{R}^{d}\right)^{n}$ be such that $\left(d f_{n}\right)_{y}(w)=\xi$. The existence of $w$ follows from the surjectivity of $\left(d f_{n}\right)_{y}$. Since $H_{n}$ is linear, we obtain that

$$
\begin{aligned}
\left(d S_{l}\right)_{h}\left(H_{n}(w)\right) & =\left.\frac{d}{d \varepsilon}\right|_{\varepsilon=0} S_{l}\left(H_{n}(y)+\varepsilon \cdot H_{n}(w)\right)=\left.\frac{d}{d \varepsilon}\right|_{\varepsilon=0} S_{l}\left(H_{n}(y+\varepsilon \cdot w)\right) \\
& =\left.\frac{d}{d \varepsilon}\right|_{\varepsilon=0} f_{n}(y+\varepsilon \cdot w)=\left(d f_{n}\right)_{y}(w)=\xi
\end{aligned}
$$

Therefore, the pair $(h, g)$ satisfies property $(\mathbf{P})$, and, thus, $\mathcal{P}$ is nonempty.
REMARK 4.3. Theorem 1.4 in [3] was stated for SDEs in which the vector fields are of class $C_{b}^{\infty}$. However, the uniform boundedness assumption was not relevant since the argument was essentially local. To elaborate this, first note that Condition (b) in [3], Theorem 3.1, which is the key ingredient for the proof of [3], Theorem 1.4, was stated under convergence in probability. In addition, Condition (b) and the result of Theorem 3.3 in the same paper were both local since the path $h$ is fixed. Therefore, one can localize the relevant probabilities in a large compact subset and apply local continuity theorems in rough path theory, even though the underlying vector fields are not bounded. In particular, the result applies to our truncated signature process $U_{t}$.

REMARK 4.4. When $H>1 / 2$, it is not clear whether $\overline{\mathcal{H}}$ contains the space of piecewise linear paths. It is though true from $\overline{\mathcal{H}}=I_{0+}^{H+1 / 2}\left(L^{2}([0,1])\right)$ that it contains all smooth paths. To fix this issue, one can reparametrize the piecewise linear path $y_{1} \sqcup \cdots \sqcup y_{n}$ in an obvious way such that the resulting path is smooth but the trajectory remains unchanged. This does not change the truncated signature, as it is invariant under reparametrization.

Essentially, the same amount of effort allows us to adapt the above argument to establish the general positivity result for hypoelliptic SDEs, as stated in Theorem 1.5, which is of independent interest. This complements the result of [3], Theorem 1.4, by affirming that Hypothesis 1.2 in that theorem is always verified under hypoellipticity.

Proof of Theorem 1.5. Without loss of generality, we only consider $t=1$. Continuing to denote by $\Phi_{t}(x ; h)$ the skeleton of equation (1.1), defined by (1.4), let $F: \overline{\mathcal{H}} \rightarrow \mathbb{R}^{N}$ be the end-point map defined by $F(h) \triangleq \Phi_{1}(x ; h)$. As in the proof of Lemma 4.2, the key ingredient that needs to be established is the following property: for any $y \in \mathbb{R}^{N}$, there exists $h \in \overline{\mathcal{H}}$ such that

$$
\begin{equation*}
F(h)=y \quad \text { and } \quad(d F)_{h}: \overline{\mathcal{H}} \rightarrow \mathbb{R}^{N} \quad \text { is surjective. } \tag{4.7}
\end{equation*}
$$

Along the same lines as in the proof of Lemma 4.2, we define $\mathcal{P}$ to be the set of points $y \in \mathbb{R}^{N}$ satisfying (4.7) for some $h \in \overline{\mathcal{H}}$. We first show that $\mathcal{P}$ is nonempty which then implies $\mathcal{P}=\mathbb{R}^{N}$ again by a translation argument.

To show that $\mathcal{P}$ is nonempty, we first define $H_{n}:\left(\mathbb{R}^{d}\right)^{n} \rightarrow \overline{\mathcal{H}}$ and $\overline{\mathcal{H}}_{0} \subseteq \overline{\mathcal{H}}$ in the same way, as in the proof of Lemma 4.2. Also, define a map $F_{n}$ by

$$
F_{n} \triangleq F \circ H_{n}:\left(\mathbb{R}^{d}\right)^{n} \rightarrow \mathbb{R}^{N}
$$

According to Sard's theorem, the set of critical values of $F_{n}$, again denoted as $E_{n}$, is a Lebesgue null set in $\mathbb{R}^{N}$, and so is $E \triangleq \bigcup_{n} E_{n}$.

Next, consider a given $q \in \mathbb{R}^{N}$. Thanks to the hypoellipticity assumption (1.2), we can equip a neighborhood $U_{q}$ of $q$ with a sub-Riemannian metric, by requiring that a certain subset of $\left\{V_{1}, \ldots, V_{d}\right\}$ is an orthonormal frame near $q$. Then, according to the Chow-Rashevskii theorem (cf. [25], Theorem 2.1.2), every point in $U_{q}$ is reachable from $q$ by a horizontal path. And if one examines the proof of the theorem in Section 2.4 of [25] carefully, this horizontal path is controlled by a piecewise linear path in $\mathbb{R}^{d}$, that is, $U_{q} \subseteq \bigcup_{n} \Phi_{1}\left(q ; H_{n}\left(\left(\mathbb{R}^{d}\right)^{n}\right)\right)$. Now, for given $y \in \mathbb{R}^{N}$, choose an arbitrary continuous path $\gamma$ joining $x$ to $y$. By compactness, we can cover the image of $\gamma$ by finitely many open sets of the form $U_{q_{i}}$ such that $U_{q_{i}} \cap U_{q_{i+1}} \neq \varnothing$ for all $i$ where $q_{i} \in \operatorname{Im}(\gamma)$. It follows that $y$ can be reached from $x$ by a horizontal path controlled by a piecewise linear path in $\mathbb{R}^{d}$. In other words, we have $y \in F_{n}\left(\left(\mathbb{R}^{d}\right)^{n}\right)$ for some $n$. This establishes the property that $\mathbb{R}^{N}=\bigcup_{n} F_{n}\left(\left(\mathbb{R}^{d}\right)^{n}\right)$.

Now, the same argument, as in the proof of Lemma 4.2, allows us to conclude that

$$
\mathbb{R}^{N} \backslash E=\bigcup_{n=1}^{\infty} F_{n}\left(\left(\mathbb{R}^{d}\right)^{n}\right) \backslash E \subseteq \mathcal{P}
$$

showing that $\mathcal{P}$ is nonempty since $E$ is a Lebesgue null set.
Finally, we show that $\mathcal{P}=\mathbb{R}^{N}$. To this aim, first note that, for any $h_{0}, \gamma, \alpha \in \overline{\mathcal{H}}$ and $\varepsilon>0$, we have

$$
\Phi_{1}\left(x ;\left(h_{0}+\varepsilon \cdot \gamma\right) \sqcup \alpha\right)=\Phi_{1}\left(\Phi_{1}\left(x ; h_{0}+\varepsilon \cdot \gamma\right) ; \alpha\right)
$$

where paths are always assumed to be parametrized on $[0,1]$. Therefore, by differentiating with respect to $\varepsilon$ at $\varepsilon=0$, we obtain that

$$
(d F)_{h_{0} \sqcup \alpha}(\gamma \sqcup 0)=J_{1}\left(F\left(h_{0}\right) ; \alpha\right) \circ(d F)_{h_{0}}(\gamma),
$$

where recall that $J_{t}(\cdot ; \cdot)$ is the Jacobian of the flow $\Phi_{t}$. This shows that

$$
\begin{equation*}
(d F)_{h_{0} \sqcup \alpha}=J_{1}\left(F\left(h_{0}\right) ; \alpha\right) \circ(d F)_{h_{0}} . \tag{4.8}
\end{equation*}
$$

Now, pick any fixed $y_{0} \in \mathcal{P}$ with an associated $h_{0} \in \overline{\mathcal{H}}$ satisfying (4.7). For any $\eta \in \mathbb{R}^{N}$, choose $\alpha \in \overline{\mathcal{H}}$ such that $F(\alpha)=\eta$. Then, $F\left(h_{0} \sqcup \alpha\right)=y+\eta$ and the surjectivity of $(d F)_{h_{0} \sqcup \alpha}$ follows from (4.8), the surjectivity of $(d F)_{h_{0}}$ and the invertibility of the Jacobian. In particular, $y+\eta \in \mathcal{P}$. Since $\eta$ is arbitrary, we conclude that $\mathcal{P}=\mathbb{R}^{N}$.

REMARK 4.5. A general support theorem for hypoelliptic SDEs allows one to show that the support of the density $p_{t}(x, y)$ is dense. In the diffusion case, together with the semigroup property

$$
p(s+t, x, y)=\int_{\mathbb{R}^{N}} p(s, x, z) p(t, z, y) d z
$$

one immediately sees that $p(t, x, y)$ is everywhere strictly positive. This argument clearly breaks down in the fractional Brownian setting.

REMARK 4.6. In Theorem 1.5, although the time horizon is taken to be [0, 1], the proof clearly applies to arbitrary time horizon [ $0, T$ ] without difficulty. As for the study of the control distance function, relation (3.3) also allows us to restrict on the time horizon $[0,1]$.

Finally, we present the main result in this part which gives a precise local lower estimate for the density $\rho_{t}(u)$. This will be a consequence of Lemma 4.2 together with the scaling property and left invariance for the process $\Gamma_{t}$.

Proposition 4.7. For each $M>0$, define $\beta_{M} \triangleq \inf \left\{\rho_{1}(u):\|u\|_{\mathrm{CC}} \leq M\right\}$. Then, $\beta_{M}$ is strictly positive, and, for all $(u, t) \in \mathfrak{g}^{(l)} \times(0,1]$ with $\|u\|_{\mathrm{CC}} \leq M t^{H}$, we have

$$
\begin{equation*}
\rho_{t}(u) \geq \beta_{M} t^{-H v} \tag{4.9}
\end{equation*}
$$

where the constant $v$ is given by $v \triangleq \sum_{k=1}^{l} k \operatorname{dim} \mathcal{L}_{k}$ and $\mathcal{L}_{k}$ is the space of homogeneous Lie polynomials of degree $k$.

Proof. First, observe that the strict positivity of $\beta_{M}$ is a direct consequence of Lemma 4.2 and the compactness of the set $\left\{u \in \mathfrak{g}^{(l)} ;\|u\|_{\text {CC }} \leq M\right\}$. Next, using the scaling property of fractional Brownian motion and the left invariance of the vector fields $\tilde{W}_{\alpha}$ defining the equation (4.4) for $\Gamma_{t}$, it is not hard to see that

$$
\left(\delta_{\lambda} \Gamma_{t}\right)_{0 \leq t \leq 1} \stackrel{\text { law }}{=}\left(\Gamma_{\lambda^{1 / H}}\right)_{0 \leq t \leq 1},
$$

where recall that $\delta_{\lambda}$ is the dilation operator on $T^{(l)}$. As a result, if we define $Q_{t}$ to be the law of $U_{t}$ on $\mathfrak{g}^{(l)}$, then $Q_{s} \circ \delta_{\lambda}^{-1}=Q_{\lambda^{1 / H_{s}}}$ for all $s \in[0,1]$. In particular, by setting $s=1$ and $\lambda=t^{H}$, we obtain that

$$
\begin{equation*}
Q_{t}=Q_{1} \circ \delta_{t^{H}}^{-1} \tag{4.10}
\end{equation*}
$$

It follows that, for any $f \in C_{b}^{\infty}\left(\mathfrak{g}^{(l)}\right)$, we have

$$
\begin{aligned}
\int_{\mathfrak{g}^{(l)}} f(u) \rho_{t}(u) d u & =\int_{\mathfrak{g}^{(l)}} f(u) Q_{t}(d u)=\int_{\mathfrak{g}^{(l)}} f\left(\delta_{t^{H}} u\right) Q_{1}(d u) \\
& =\int_{\mathfrak{g}^{(l)}} f\left(\delta_{t^{H}} u\right) \rho_{1}(u) d u=\int_{\mathfrak{g}^{(l)}} t^{-H v} f(u) \rho_{1}\left(\delta_{t^{-H}} u\right) d u,
\end{aligned}
$$

where the last equality follows from the change of variables $u \leftrightarrow \delta_{t^{-H}} u$ and the fact that $d u \circ \delta_{t^{H}}^{-1}=t^{-H v} d u$ (cf. relation (2.8)). Therefore, we conclude that

$$
\begin{equation*}
\rho_{t}(u)=t^{-H v} \rho_{1}\left(\delta_{t^{H}}^{-1} u\right) \quad \text { for all }(u, t) \in \mathfrak{g}^{(l)} \times(0,1] \tag{4.11}
\end{equation*}
$$

from which our result (4.9) follows.
4.2. Step two: Local lower estimate for the density of the Taylor approximation process. Recall from (4.3) that $X_{l}(t, x)=x+F_{l}\left(U_{t}^{(l)}, x\right)$ is the Taylor approximation process of order $l$ for the actual solution of the $\operatorname{SDE}$ (1.1). Due to hypoellipticity, it is natural to expect that $F_{l}$ is "nondegenerate" in a suitable sense, provided $l \geq l_{0}$. In addition, a precise local lower estimate for the density of $X_{l}(t, x)$ should naturally follow from Proposition 4.7 in step one, combined with such "nondegeneracy" property of $F_{l}$. Here, the main subtlety lies in finding a way of respecting the fractional Brownian scaling and Cameron-Martin structure so that the estimate we obtain on $X_{l}(t, x)$ is sharp. However, this will be a direct consequence of Theorem 1.3 together with a result for the diffusion case in [20]. In this part we always fix $l \geq l_{0}$.

We first give the precise meaning of the nondegeneracy of $F_{l}$. Let $J F_{l}(u, x): \mathfrak{g}^{(l)} \rightarrow \mathbb{R}^{N}$ be the Jacobian of $F_{l}$ with respect to $u$. Since $\mathfrak{g}^{(l)}$ has a canonical Hilbert structure, induced from $T^{(l)}\left(\mathbb{R}^{d}\right)$, we can also consider the adjoint map $J F_{l}(u, x)^{*}: \mathbb{R}^{N} \rightarrow \mathfrak{g}^{(l)}$. The nondegeneracy of $J F_{l}$ is summarized in the following lemma, which was proved in [20], Lemma 3.13.

Lemma 4.8. Let $F_{l}$ be the approximation map given in Definition 2.11, and let $J F_{l}(u, x): \mathfrak{g}^{(l)} \rightarrow \mathbb{R}^{N}$ be its Jacobian. Then, there exists a constant $c>0$, depending only on $l_{0}$ and the vector fields, such that

$$
J F_{l}(0, x) \cdot J F_{l}(0, x)^{*} \geq c \cdot \operatorname{Id}_{\mathbb{R}^{N}}
$$

for all $l \geq l_{0}$ and $x \in \mathbb{R}^{N}$.

An immediate corollary of Lemma 4.8 is the following.
Corollary 4.9. Given $l \geq l_{0}$, there exists $r>0$, depending on $l$ and the vector fields, such that $\operatorname{det}\left(J F_{l}(u, x) \cdot J F_{l}(u, x)^{*}\right)$ is uniformly positive on $\left\{u \in \mathfrak{g}^{(l)}:\|u\|_{\mathrm{HS}}<r\right\} \times \mathbb{R}^{N}$. In particular, the map

$$
\left\{u \in \mathfrak{g}^{(l)}:\|u\|_{\mathrm{HS}}<r\right\} \rightarrow \mathbb{R}^{N}, \quad u \mapsto x+F_{l}(u, x)
$$

is a submersion in the sense of differential geometry.

REmark 4.10. Note that the map $F_{l}$ and the constant $r$ in Corollary 4.9 depend on $l$. For technical reasons we will assume that $r$ is chosen (still depending on $l$ ) so that, for all $l_{0} \leq l^{\prime} \leq l$, the map $J F_{l^{\prime}}\left(\pi^{\left(l^{\prime}\right)}(u), x\right)$ has full rank whenever $(u, x) \in \mathfrak{g}^{(l)} \times \mathbb{R}^{N}$ with $\|u\|_{\text {HS }}<r$, where $\pi^{\left(l^{\prime}\right)}: \mathfrak{g}^{(l)} \rightarrow \mathfrak{g}^{\left(l^{\prime}\right)}$ is the canonical projection. This property will be used in the proof of Lemma 4.20 in step three below.

Now, let $r$ be the constant given in Remark 4.10. It is standard from differential geometry that, for each $x \in \mathbb{R}^{N}$ and $y \in\left\{x+F_{l}(u, x):\|u\|_{\mathrm{HS}}<r\right\}$, the "bridge space"

$$
\begin{equation*}
M_{x, y} \triangleq\left\{u \in \mathfrak{g}^{(l)}:\|u\|_{\mathrm{HS}}<r \text { and } x+F_{l}(u, x)=y\right\} \tag{4.12}
\end{equation*}
$$

is a submanifold of $\left\{u \in \mathfrak{g}^{(l)}:\|u\|_{\text {HS }}<r\right\}$ with dimension $\operatorname{dim} \mathfrak{g}^{(l)}-N$. In addition, since both of $\mathfrak{g}^{(l)}$ and $\mathbb{R}^{N}$ are oriented Riemannian manifolds, we know from differential topology that $M_{x, y}$ carries a natural orientation and hence a volume form which we denote as $m_{x, y}$. The following result is the standard disintegration formula in Riemannian geometry (cf. [6], equation (0.3)).

Proposition 4.11. For any $\varphi \in C_{c}^{\infty}\left(\left\{u \in \mathfrak{g}^{(l)}:\|u\|_{\mathrm{HS}}<r\right\}\right)$, we have

$$
\begin{equation*}
\int_{\mathfrak{g}^{(l)}} \varphi(u) d u=\int_{\mathbb{R}^{N}} d y \int_{M_{x, y}} K(v, x) \varphi(v) m_{x, y}(d v) \tag{4.13}
\end{equation*}
$$

where the kernel $K$ is given by

$$
\begin{equation*}
K(v, x) \triangleq\left(\operatorname{det}\left(J F_{l}(v, x) \cdot J F_{l}(v, x)^{*}\right)\right)^{-\frac{1}{2}} \tag{4.14}
\end{equation*}
$$

and we define $m_{x, y} \triangleq 0$ if $M_{x, y}=\varnothing$.
The disintegration formula (4.13) immediately leads to a formula for the (localized) density of the Taylor approximation process $X_{l}(t, x)$.

Proposition 4.12. Let $\eta \in C_{c}^{\infty}\left(\left\{u \in \mathfrak{g}^{(l)}:\|u\|_{\mathrm{HS}}<r\right\}\right)$ be a bump function so that $0 \leq \eta \leq 1$ and $\eta=1$ when $\|u\|_{\mathrm{HS}}<r / 2$, where $r$ is the constant featuring in Proposition 4.11. Define $\mathbb{P}_{l}^{\eta}(t, x, \cdot)$ to be the measure

$$
\mathbb{P}_{l}^{\eta}(t, x, A) \triangleq \mathbb{E}\left[\eta\left(U_{t}\right) \mathbf{1}_{\left\{X_{l}(t, x) \in A\right\}}\right], \quad A \in \mathcal{B}\left(\mathbb{R}^{N}\right),
$$

where $U_{t}=\log \Gamma_{t}, \Gamma_{t}$ is defined by (4.4) and $X_{l}(t, x)=x+F_{l}\left(U_{t}, x\right)$ is the approximation given by (2.19). Then, the measure $\mathbb{P}_{l}^{\eta}(t, x, \cdot)$ is absolutely continuous with respect to the Lebesgue measure, and its density is given by

$$
\begin{equation*}
p_{l}^{\eta}(t, x, y) \triangleq \int_{M_{x, y}} \eta(u) K(u, x) \rho_{t}(u) m_{x, y}(d u) \tag{4.15}
\end{equation*}
$$

where $\rho_{t}$ is the density of $U_{t}$ and $K$ is given by (4.14).

To obtain a sharp lower estimate on $p_{l}^{\eta}(t, x, y)$ from formula (4.15) and the lower estimate of $\rho_{t}(u)$ given by (4.9), one needs to estimate the volume form $m_{x, y}$ precisely. For this purpose we resort to a change of variables $S_{x, y}$ from the "bridge space" $M_{x, y}$ to the "loop space" $M_{x, x}$, introduced in [20]. The construction of this map $S_{x, y}$ is based on the simple idea that, in order to transform an arbitrary loop $\alpha$ from $x$ to $x$ into a path from $x$ to $y$, we just concatenate the loop $\alpha$ to a fixed path from $x$ to $y$. However, this idea does not project to a map from $M_{x, y}$ to $M_{x, x}$ in a simple way, and one needs to use the function $\Psi_{l}$, defined in Lemma 3.1, to make it work at the level of Taylor approximation. We summarize the technical formulation of this map $S_{x, y}$ in the following lemma which was proved by [20] in the diffusion case (i.e., $H=1 / 2$ ). Thanks to our Theorem 1.3, the lemma (in particular the estimate (4.19) below) holds in exactly the same way in the fractional Brownian context for $H>1 / 4$.

Lemma 4.13. Recall that the function $\Psi_{l}$ is defined in Lemma 3.1. We define the operation $\times$ to be the multiplication induced from $G^{(l)}$ through the exponential map, namely,

$$
\begin{equation*}
v \times u \triangleq \log (\exp (v) \otimes \exp (u)), \quad v, u \in \mathfrak{g}^{(l)} \tag{4.16}
\end{equation*}
$$

Then, the following statements hold true:
(i) There exist $\varepsilon, \rho_{1}>0$ and $\rho_{2} \in(0, r)$ such that, for any given $x \in \mathbb{R}^{N}$ and $h \in \overline{\mathcal{H}}$ with $\|h\|_{\overline{\mathcal{H}}}<\rho_{1}$, the map

$$
v \mapsto \tilde{\Psi}_{x, h}(v) \triangleq \Psi_{l}\left(v \times u, x, y-x-F_{l}(v \times u, x)\right)
$$

where $y \triangleq \Phi_{1}(x ; h)$ and $u \triangleq \log S_{l}(h)$, defines a diffeomorphism from an open neighbourhood $V_{x, h} \subseteq \mathfrak{g}^{(l)}$ of 0 containing the ball $\left\{v \in \mathfrak{g}^{(l)}:\|v\|_{\mathrm{HS}}<\varepsilon\right\}$ onto $W \triangleq\left\{w \in \mathfrak{g}^{(l)}\right.$ : $\left.\|w\|_{\mathrm{HS}}<\rho_{2}\right\}$ such that

$$
v \in V_{x, h} \cap M_{x, x} \quad \Longleftrightarrow \quad w \triangleq \tilde{\Psi}_{x, h}(v) \in W \cap M_{x, y}
$$

(ii) Given $x, y \in \mathbb{R}^{N}$ with $d(x, y)<\rho_{1} / 2$, we choose $h \in \Pi_{x, y}$ satisfying

$$
d(x, y) \leq\|h\|_{\overline{\mathcal{H}}} \leq 2 d(x, y)<\rho_{1}
$$

and define

$$
\begin{equation*}
\left.S_{x, y} \triangleq \tilde{\Psi}_{x, h}^{-1}\right|_{W \cap M_{x, y}}: W \cap M_{x, y} \rightarrow V_{x, h} \cap M_{x, x} \tag{4.17}
\end{equation*}
$$

Then, there exists a constant $\Lambda>0$ such that

$$
\begin{equation*}
\frac{1}{\Lambda} \cdot m_{x, x}(\cdot) \leq m_{x, y} \circ S_{x, y}^{-1}(\cdot) \leq \Lambda \cdot m_{x, x}(\cdot) \quad \text { on } V_{x, h} \cap M_{x, x} \tag{4.18}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{1}{\Lambda} \cdot\left(\|v\|_{\mathrm{CC}}+d(x, y)\right) \leq\left\|S_{x, y}^{-1}(v)\right\|_{\mathrm{CC}} \leq \Lambda \cdot\left(\|v\|_{\mathrm{CC}}+d(x, y)\right) \tag{4.19}
\end{equation*}
$$

for any $v \in V_{x, h} \cap M_{x, x}$.
Now, we apply the change of variables involving $S_{x, y}$ in Lemma 4.13 to establish a lower estimate of the density $p_{l}^{\eta}(t, x, y)$ in terms of the measure $m_{x, x}$ that does not depend on $y$.

LEMMA 4.14. Let $p_{l}^{\eta}(t, x, y)$ be the density defined by (4.15), and recall that the exponent $v$ is defined by $v \triangleq \sum_{k=1}^{l} k \operatorname{dim} \mathcal{L}_{k}$. Then, there exist constants $C, \tau>0$ such that, for all $x, y$, $t$ with $d(x, y) \leq t^{H}$ and $0<t<\tau$, we have

$$
\begin{equation*}
p_{l}^{\eta}(t, x, y) \geq C t^{-H v} m_{x x}\left(\left\{v \in M_{x, x}:\|v\|_{\mathrm{CC}} \leq t^{H}\right\}\right), \tag{4.20}
\end{equation*}
$$

where $m_{x, x}$, is the volume form on $M_{x, x}$ given by (4.12)

Proof. Lemma 4.13 asserts that there exists $\rho_{1}>0$ such that if $d(x, y)<\rho_{1} / 2$, then $S_{x, y}$ given by (4.17) defines a change of variables (i.e., a diffeomorphism) for (4.15). Specifically, we have

$$
\begin{aligned}
p_{l}^{\eta}(t, x, y) & \geq \int_{M_{x, y} \cap W} \eta(u) K(u, x) \rho_{t}(u) m_{x, y}(d u) \\
& =\int_{M_{x, x} \cap V_{x, h}} \eta\left(S_{x, y}^{-1} v\right) K\left(S_{x, y}^{-1} v, x\right) \rho_{t}\left(S_{x, y}^{-1} v\right) m_{x, y} \circ S_{x, y}^{-1}(d v)
\end{aligned}
$$

In addition, since $V_{x, h}$ contains the ball $\left\{v \in \mathfrak{g}^{(l)}:\|v\|_{\text {HS }}<\varepsilon\right\}$, owing to relation (4.18) and thanks to the fact that $K$ defined by (4.14) is bounded below, we obtain

$$
p_{l}^{\eta}(t, x, y) \geq C_{H, V, l} \int_{M_{x, x} \cap\left\{v \in \mathfrak{g}^{(l)}:\|v\|_{\mathrm{HS}}<\varepsilon\right\}} \rho_{t}\left(S_{x, y}^{-1} v\right) m_{x, x}(d v) .
$$

Now, choose $\tau<\left(\rho_{1} / 2\right)^{\frac{1}{H}}$ to be such that

$$
0<t<\tau \quad \Longrightarrow \quad\left\{v \in \mathfrak{g}^{(l)}:\|v\|_{\mathrm{CC}} \leq t^{H}\right\} \subseteq\left\{v \in \mathfrak{g}^{(l)}:\|v\|_{\mathrm{HS}}<\varepsilon\right\} .
$$

We will thus lower bound $p_{l}^{\eta}(t, x, y)$ as follows:

$$
\begin{equation*}
p_{l}^{\eta}(t, x, y) \geq C_{H, V, l} \int_{M_{x, x} \cap\left\{v \in \mathfrak{g}^{(l)}:\|v\|_{\left.\mathrm{CC} \leq t^{H}\right\}}\right.} \rho_{t}\left(S_{x, y}^{-1} v\right) m_{x, x}(d v) . \tag{4.21}
\end{equation*}
$$

Next, according to the second inequality of (4.19), if $d(x, y) \leq t^{H}$ and $t<\tau$ (so that $\left.d(x, y)<\rho_{1} / 2\right)$, then

$$
\left\|S_{x, y}^{-1} v\right\|_{\mathrm{CC}} \leq 2 C t^{H}
$$

provided that $v \in M_{x, x}$ with $\|v\|_{\mathrm{CC}} \leq t^{H}$. For such $x, y, t, v$, by Proposition 4.7 we have

$$
\rho_{t}\left(S_{x, y}^{-1} v\right) \geq \beta_{2 C} t^{-H v}
$$

Plugging this inequality into (4.21), we arrive at

$$
p_{l}^{\eta}(t, x, y) \geq C_{H, V, l} \beta_{2 C} t^{-H v} m_{x, x}\left(\left\{v \in M_{x, x}:\|v\|_{\mathrm{CC}} \leq t^{H}\right\}\right)
$$

The next lemma relates the measure $m_{x, x}$ with the volume of the ball $B_{d}\left(x, t^{H}\right)$ defined with respect to the control distance function.

Lemma 4.15. Let $M_{x, x}$ be the set defined by (4.12), and recall that $m_{x, x}$ is the volume measure on $M_{x, x}$. There exist constants $C, \tau>0$ such that

$$
\begin{equation*}
t^{-H v} m_{x, x}\left(\left\{v \in M_{x, x}:\|v\|_{\mathrm{CC}} \leq t^{H}\right\}\right) \geq \frac{C}{\left|B_{d}\left(x, t^{H}\right)\right|} \tag{4.22}
\end{equation*}
$$

for all $x \in \mathbb{R}^{N}$ and $0<t<\tau$.
Proof. We know from [20], Lemma 3.31, that, for any $\beta \in(0,1)$, there exists a constant $C_{\beta}$ such that

$$
\begin{equation*}
t^{-v / 2} m_{x x}\left(\left\{v \in M_{x, x}:\|v\|_{\mathrm{CC}} \leq t^{1 / 2}\right\}\right) \geq \frac{C_{\beta}}{\left|B_{d_{\mathrm{BM}}}\left(x, \beta t^{1 / 2}\right)\right|} \tag{4.23}
\end{equation*}
$$

for all $(t, x) \in(0,1] \times \mathbb{R}^{N}$, where $d_{\mathrm{BM}}$ is the control distance function for the diffusion case. On the other hand, according to Theorem 1.3, there exist constants $C_{H}$ and $\delta$ such that

$$
d(y, x) \leq C_{H} d_{\mathrm{BM}}(y, x)
$$

when $|y-x|<\delta$. If we choose $\beta=\frac{1}{C_{H}}$, when $t$ is small, we have

$$
B_{d_{\mathrm{BM}}}\left(x, \beta t^{1 / 2}\right) \subseteq B_{d}\left(x, t^{1 / 2}\right) \subseteq\{y:|y-x|<\delta\},
$$

and, thus, the right-hand side of (4.23) is further bounded below by $\frac{C_{\beta}}{\left|B_{d}\left(x, t^{1 / 2}\right)\right|}$. The desired inequality (4.22) follows by changing $t \mapsto t^{2 H}$.

Summarizing the contents of Lemma 4.14 and Lemma 4.15, we have obtained the following lower bound on $p_{l}^{\eta}(t, x, y)$ which finishes the second step of the main strategy.

COROLLARY 4.16. Let $p_{l}^{\eta}(t, x, y)$ be the density given by (4.15), and recall the notations of Lemma 4.14. Then, there exist constants $C, \tau>0$, depending only on $H, l$ and the vector fields, such that

$$
\begin{equation*}
p_{l}^{\eta}(t, x, y) \geq \frac{C}{\left|B_{d}\left(x, t^{H}\right)\right|} \tag{4.24}
\end{equation*}
$$

for all $x, y, t$ satisfying $d(x, y) \leq t^{H}$ and $0<t<\tau$.
4.3. Step three: Comparing approximating and actual densities. The last step toward the proof of Theorem 1.6 will be to show that the approximating density $p_{l}^{\eta}(t, x, y)$ and the actual density $p(t, x, y)$ of $X_{t}^{x}$ are close to each other when $t$ is small. For this part we combine the Fourier transform approach developed in [20] with general estimates for Gaussian rough differential equations. As before, we assume that $l \geq l_{0}$.

Recall that the Fourier transform of a function $f(y)$ on $\mathbb{R}^{N}$ is defined by

$$
\mathcal{F} f(\xi)=\hat{f}(\xi) \triangleq \int_{\mathbb{R}^{N}} f(y) \mathrm{e}^{2 \pi i\langle\xi, y\rangle} d y, \quad \xi \in \mathbb{R}^{N}
$$

In the sequel we will consider the Fourier transform $\hat{p}(t, x, \xi)$ (respectively, $\hat{p}_{l}^{\eta}(t, x, \xi)$ ) of the density $p(t, x, y)$ (respectively, $p_{l}^{\eta}(t, x, y)$ ) with respect to the $y$-variable. We will invoke the following trivial bound on $p-p_{l}$ in terms of $\hat{p}$ and $\hat{p}_{l}^{\eta}$ :

$$
\begin{equation*}
\left|p(t, x, y)-p_{l}^{\eta}(t, x, y)\right| \leq \int_{\mathbb{R}^{N}}\left|\hat{p}(t, x, \xi)-\hat{p}_{l}^{\eta}(t, x, \xi)\right| d \xi \tag{4.25}
\end{equation*}
$$

Therefore, our aim in this section will be to estimate the right-hand side of (4.25) by considering two regions $\{|\xi| \leq R\}$ and $\{|\xi|>R\}$ separately in the integral, where $R$ is some large number to be chosen later on.
4.3.1. Integrating relation (4.25) in a neighborhood of the origin. We first integrate our Fourier variable $\xi$ in (4.25) over the region $\{|\xi| \leq R\}$. In this case we make use of a tail estimate for the error of the Taylor approximation of $X_{t}^{x}$ which is provided below.

LEMMA 4.17. Let $X_{t}^{x}$ be the solution to the $\operatorname{SDE}$ (1.1), and consider its approximation $X_{l}(t, x)$ of order $l \geq \max \left\{l_{0}, H^{-1}\right\}$, as defined by (4.3). Fix $\bar{l} \in(l, l+1)$. Then, there exist constants $C_{1}, C_{2}$, depending only on $H, l$ and the vector fields, such that, for all $t \in(0,1]$ and $x, y \in \mathbb{R}^{N}$, we have

$$
\begin{equation*}
\mathbb{P}\left(\left|X_{t}^{x}-X_{l}(t, x)\right| \geq \lambda\right) \leq C_{1} \exp \left(-\frac{C_{2} \lambda^{2 / \bar{l}}}{t^{2 H}}\right) \quad \text { for all } \lambda>0 \tag{4.26}
\end{equation*}
$$

Proof. According to [13], Corollary 10.15, we have the following pathwise estimate:

$$
\begin{equation*}
\left|X(t, x)-X_{l}(t, x)\right| \leq C \cdot\|\mathbf{B}\|_{p-\operatorname{var} ;[0, t]}^{\bar{l}}, \tag{4.27}
\end{equation*}
$$

where $C=C_{H, V, l}>0$ and $\mathbf{B}$ is the rough path lifting of $B$ alluded to in Proposition 2.10. In equation (4.27) the parameter $p$ is any number greater that $1 / H$, and the $p$-variation norm is defined with respect to the CC-norm. It follows from (4.27) that, for any $\lambda>0$ and $\eta>0$, we have

$$
\mathbb{P}\left(\left|X_{t}^{x}-X_{l}(t, x)\right| \geq \lambda\right) \leq \mathbb{P}\left(\|\mathbf{B}\|_{p-\text { var; }[0, t]}^{\bar{l}} \geq \lambda / C\right)
$$

In addition, the fBm signature satisfies the identity in law

$$
\left(\mathbf{B}_{s}\right)_{0 \leq s \leq t} \stackrel{d}{=}\left(\delta_{t} \circ \mathbf{B}_{s / t}\right)_{0 \leq s \leq t} .
$$

Owing to the scaling properties of the CC-norm, we thus get that, for an arbitrary $\zeta>0$, we have

$$
\begin{align*}
\mathbb{P}\left(\left|X_{t}^{x}-X_{l}(t, x)\right| \geq \lambda\right) & \leq \mathbb{P}\left(\|\mathbf{B}\|_{p \text {-var; }[0,1]} \geq \frac{(\lambda / C)^{1 / \bar{l}}}{t^{H}}\right) \\
& \leq \exp \left(-\frac{\zeta(\lambda / C)^{2 / \bar{l}}}{t^{2 H}}\right) \cdot \mathbb{E}\left[\mathrm{e}^{\left.\zeta\|\mathbf{B}\|_{p-\text { var; }[0,1]}^{2}\right]}\right. \tag{4.28}
\end{align*}
$$

According to a Fernique-type estimate for the fractional Brownian rough path (cf. [13], Theorem 15.33), we know that there exists $\zeta=\zeta_{H}>0$ such that

$$
\mathbb{E}\left[\mathrm{e}^{\zeta\|\mathbf{B}\|_{p-\text { var; }[0,1]}^{2}}\right]<\infty
$$

Plugging this inequality into (4.28), our conclusion (4.26) is easily obtained.
We are now ready to derive a Fourier transform estimate for small values of $\xi$.
LEMMA 4.18. Under the same notation as in Lemma 4.17, let $p(t, x, \cdot)$ be the density of the random variable $X_{t}^{x}$, and let $p_{l}^{\eta}(t, x, \cdot)$ be the approximating density given by (4.15). Then, their Fourier transforms satisfy the following inequality over the region $\{|\xi| \leq R\}$ :

$$
\begin{equation*}
\left|\hat{p}(t, x, \xi)-\hat{p}_{l}^{\eta}(t, x, \xi)\right| \leq C_{H, V, l}(1+|\xi|) t^{H \bar{l}} \tag{4.29}
\end{equation*}
$$

provided that $t<\tau_{1}$ for some constant $\tau_{1}$, depending on $H, l$ and the vector fields.
Proof. Notice that, according to our definition (4.15) of $p_{l}^{\eta}$, we have

$$
\hat{p}(t, x, \xi)=\mathbb{E}\left[e^{2 \pi i\left\langle\xi, X_{t}^{x}\right\rangle}\right] \quad \text { and } \quad \hat{p}_{l}^{\eta}(t, x, \xi)=\mathbb{E}\left[\eta\left(U_{t}\right) e^{2 \pi i\left\langle\xi, X_{l}(t, x)\right\rangle}\right] .
$$

Hence, it is easily seen that

$$
\begin{align*}
\left|\hat{p}(t, x, \xi)-\hat{p}_{l}^{\eta}(t, x, \xi)\right| & \leq \mathbb{E}\left[\left|\mathrm{e}^{-2 \pi i\left\langle\xi, X_{t}^{x}\right\rangle}-\mathrm{e}^{-2 \pi i\left\langle\xi, X_{l}(t, x)\right\rangle}\right|\right]+\mathbb{E}\left[1-\eta\left(U_{t}\right)\right] \\
& \leq 2 \pi|\xi| \cdot \mathbb{E}\left[\left|X_{t}^{x}-X_{l}(t, x)\right|\right]+\mathbb{E}\left[1-\eta\left(U_{t}\right)\right] \tag{4.30}
\end{align*}
$$

Now, in order to bound the right-hand side of (4.30), we first invoke Lemma 4.17. This yields

$$
\begin{align*}
\mathbb{E}\left[\left|X_{t}^{x}-X_{l}(t, x)\right|\right] & =\int_{0}^{\infty} \mathbb{P}\left(\left|X_{t}^{x}-X_{l}(t, x)\right| \geq \lambda\right) d \lambda \\
& \leq C_{1} \int_{0}^{\infty} \exp \left(-\frac{C_{2} \lambda^{2 / \bar{l}}}{t^{2 H}}\right) d \lambda=C_{3} t^{H \bar{l}} \tag{4.31}
\end{align*}
$$

On the other hand, using a similar argument to the proof of Lemma 4.17, we see that there exists a strictly positive exponent $\alpha_{H, l}$ such that

$$
\begin{equation*}
\mathbb{E}\left[1-\eta\left(U_{t}\right)\right] \leq \mathbb{P}\left(\left\|U_{t}\right\|_{\mathrm{HS}} \geq \frac{r}{2}\right) \leq C_{4} \cdot \exp \left(-\frac{C_{5}}{t^{\alpha_{H, l}}}\right) \tag{4.32}
\end{equation*}
$$

By taking $t$ small enough, we can make the right-hand side of (4.32) smaller than $C_{6} t^{H \bar{l}}$. Hence, there exists $\tau_{1}>0$ such that if $t \leq \tau_{1}$, we have

$$
\begin{equation*}
\mathbb{E}\left[1-\eta\left(U_{t}\right)\right] \leq C_{6} t^{H \bar{l}} \tag{4.33}
\end{equation*}
$$

Now, combining (4.31) and (4.33), we easily get our conclusion (4.29).
4.3.2. Integrating relation (4.25) for large Fourier modes. We now integrate the Fourier variable $\xi$ over the region $\{|\xi|>R\}$. In this case we make use of certain upper estimates for $p(t, x, y)$ and $p_{l}^{\eta}(t, x, y)$. We start with a bound on the density of $X_{t}^{x}$ which is also of independent interest. The main ingredients of the proof are basically known in the literature, but, to our best knowledge, the result (for the hypoelliptic case) has not been formulated elsewhere.

Proposition 4.19. Let $p(t, x, y)$ be the density of the random variable $X_{t}^{x}$. As in Lemma 4.18, we assume that the uniform hypoellipticity condition (1.2) is satisfied. Then, for each $n \geq 1$, there exist constants $C_{1, n}, C_{2, n}, v_{n}>0$, depending on $n, H$ and the vector fields, such that

$$
\begin{equation*}
\left|\partial_{y}^{n} p(t, x, y)\right| \leq C_{1, n} t^{-v_{n}} \exp \left(-\frac{C_{2, n}|y-x|^{2 \wedge(2 H+1)}}{t^{2 H}}\right) \tag{4.34}
\end{equation*}
$$

for all $(t, x, y) \in(0,1] \times \mathbb{R}^{N} \times \mathbb{R}^{N}$, where $\partial_{y}^{n}$ denotes the $n$th order derivative operator with respect to the $y$ variable.

Proof. Elaborating on the integration by parts invoked, for example, in [3], Relation (24), there exist exponents $\alpha, \beta, p, q>1$ such that

$$
\begin{equation*}
\left|\partial_{y}^{n} p(t, x, y)\right| \leq C_{1, n} \mathbb{P}\left(\left|X_{t}^{x}-x\right| \geq|y-x|\right)^{\frac{1}{2}} \cdot\left\|\gamma_{X_{t}^{x}}^{-1}\right\|_{\alpha, p}^{\alpha} \cdot\left\|\mathbf{D} X_{t}^{x}\right\|_{\beta, q}^{\beta}, \tag{4.35}
\end{equation*}
$$

where $\gamma_{X_{t}^{x}}$ denotes the Malliavin covariance matrix and $\|\cdot\|_{k, p}$ denotes the Gaussian-Sobolev norm. Then, with (4.35) in hand we proceed in the following way:
(i) An exponential tail estimate for $X_{t}^{x}$ yield the exponential term in (4.34). This step is achieved as in [3], Relation (25).
(ii) The Malliavin derivatives of $X_{t}^{x}$ are estimated as in [4], Lemma 3.5(1). This produces some positive powers of $t$ in (4.34).
(iii) The inverse of the Malliavin covariance matrix is bounded as in [4], Lemma 3.5(2). It gives some negative powers of $t$ in (4.34).

For the sake of conciseness, we will not detail the steps outlined as above. We refer the reader to [3, 4] for the details.

The following lemma parallels Proposition 4.19 for the approximating process $X_{l}(t, x)$.
LEMmA 4.20. Assume the same hypothesis as in Proposition 4.19. Recall that the approximating density $p_{l}^{\eta}$ is defined by (4.15). Fix $l \geq l_{0}$. Then, for each $n \geq 1$, there exist constants $C_{n}=C_{n}(H, l)$ and $\gamma_{n}=\gamma_{n}\left(H, l_{0}\right)$ such that, for all $(t, x) \in(0,1] \times \mathbb{R}^{N}$, the following bound holds true:

$$
\begin{equation*}
\left\|\partial_{y}^{n} p_{l}^{\eta}(t, x, \cdot)\right\|_{C_{b}^{n}\left(\mathbb{R}^{N}\right)} \leq C_{n} \cdot t^{-\gamma_{n}} \tag{4.36}
\end{equation*}
$$

Moreover, the function $\partial_{y}^{n} p_{l}^{\eta}(t, x, \cdot)$ is compactly supported in $\mathbb{R}^{N}$.

Proof. Recall that, for a differentiable random vector $Z=\left(Z^{1}, \ldots, Z^{n}\right)$ in the sense of Malliavin, we use the notation $\gamma_{Z} \triangleq\left(\left\langle D Z^{i}, D Z^{j}\right\rangle_{\overline{\mathcal{H}}}\right)_{1 \leq i, j \leq n}$ to denote its Malliavin covariance matrix. By the definition (4.3) of $X_{l}(t, x)$, it is immediate that

$$
\begin{equation*}
\gamma_{X_{l}(t, x)}=J F_{l}\left(U_{t}^{(l)}, x\right) \cdot \gamma_{U_{t}^{(l)}} \cdot J F_{l}\left(U_{t}^{(l)}, x\right)^{*} \tag{4.37}
\end{equation*}
$$

It follows from [1], Theorem 3.3, and Lemma 4.8 that $\gamma_{X_{l}(t, x)}^{-1} \in L^{q}$ for all $q>1$. Now, the uniform upper bound for the derivatives of $p_{l}^{\eta}(t, x, y)$ follows from the same lines as in the proof of Proposition 4.19 (with the same three main ingredients (i), (ii), (iii)), based on the integration by parts formula.

REMARK 4.21. We must point out that the exponent $\gamma_{n}$ in (4.36) depends on $l_{0}$ but not on $l$. This subtle technical point is critical for us, and its proof requires a nontrivial amount of analysis. For the sake of conciseness, we refer the reader to [14], Section 5.3 IV, for a detailed discussion on this matter.

We now return to the Fourier estimate (4.25) for the region $|\xi|>R$. In particular, we have the following result.

LEMmA 4.22. Using the same notation and hypothesis as in Lemma 4.18, the Fourier transforms $\hat{p}$ and $\hat{p}_{l}^{\eta}$ are such that, for all $|\xi|>R$, we have

$$
\begin{equation*}
|\xi|^{N+2}\left(|\widehat{p}(t, x, \xi)|+\left|\widehat{p}_{l}^{\eta}(t, x, \xi)\right|\right) \leq C \cdot t^{-\mu} \tag{4.38}
\end{equation*}
$$

for some strictly positive constants $C=C_{N, H, V, l}$ and $\mu=\mu_{N, H, V, l_{0}}$.
Proof. According to standard compatibility rules between Fourier transform and differentiation, we have (recall that $\mathcal{F} f$ and $f$ are both used to designate the Fourier transform of a function $f$ )

$$
\begin{aligned}
& |\xi|^{N+2}\left(|\hat{p}(t, x, \xi)|+\left|\hat{p}_{l}^{\eta}(t, x, \xi)\right|\right) \\
& \quad \leq C_{N}\left(\left|\mathcal{F}\left(\partial_{y}^{N+2} p(t, x, y)\right)\right|+\left|\mathcal{F}\left(\partial_{y}^{N+2} p_{l}^{\eta}(t, x, y)\right)\right|\right)
\end{aligned}
$$

Plugging (4.34) and (4.36) into this relation and using the fact that $\partial_{y}^{N+2} p_{l}^{\eta}(t, x, \cdot)$ is compactly supported, our claim (4.38) is easily proved.
4.3.3. Comparison of the densities. Combining the previous preliminary results on Fourier transforms, we get the following uniform bound on the difference $p-p_{l}^{\eta}$.

Proposition 4.23. We still keep the same notation and assumptions of Lemma 4.18. Then, there exists $\tau>0$ such that, for all $t \leq \tau$ and $x, y \in \mathbb{R}^{N}$, we have

$$
\begin{equation*}
\left|p(t, x, y)-p_{l}^{\eta}(t, x, y)\right| \leq C_{H, V, l} t . \tag{4.39}
\end{equation*}
$$

Proof. Thanks to (4.25), we can write

$$
\left|p(t, x, y)-p_{l}^{\eta}(t, x, y)\right| \leq\left(\int_{|\xi| \leq \mathbb{R}}+\int_{|\xi|>\mathbb{R}}\right)\left|\hat{p}(t, x, \xi)-\hat{p}_{l}^{\eta}(t, x, \xi)\right| d \xi
$$

Next, we invoke the bounds (4.29) and (4.38), which allows us to write

$$
\left|p(t, x, y)-p_{l}^{\eta}(t, x, y)\right| \leq C_{1}\left(R^{N+1} t^{H \bar{l}}+t^{-\mu} \int_{|\xi|>R}|\xi|^{-N-2} d \xi\right), \quad t \in[0,1]
$$

where we recall that $\bar{l}$ is a fixed number in $[l, l+1]$ introduced in Lemma 4.17. Now, an elementary change of variable yields

$$
\begin{align*}
& \left|p(t, x, y)-p_{l}^{\eta}(t, x, y)\right| \\
& \quad \leq C_{1}\left(R^{N+1} t^{H \bar{l}}+t^{-\mu} R^{-1} \int_{|\xi| \geq 1}|\xi|^{-(N+1)} d \xi\right)  \tag{4.40}\\
& \quad \leq C_{2}\left(R^{N+1} t^{H \bar{l}}+t^{-\mu} R^{-1}\right)
\end{align*}
$$

We can easily optimize expression (4.40) with respect to $R$ by choosing $R=t^{-(\mu+1)}$. It follows that

$$
\begin{equation*}
\left|p(t, x, y)-p_{l}^{\eta}(t, x, y)\right| \leq C_{2} t^{-(N+1)(\mu+1)+H \bar{l}}+t \tag{4.41}
\end{equation*}
$$

for all $t \in(0,1]$. In addition, recall that a crucial point in our approach is that the exponent $\mu$ in (4.41) does not depend on $l$. Therefore, we can choose $l \geq l_{0}$ large enough so that

$$
-(N+1)(\mu+1)+H \bar{l} \geq 1
$$

For this value of $l$, the upper bound (4.39) is easily deduced from (4.41).
4.4. Completing the proof of Theorem 1.6. Finally, we are in a position to complete the proof of Theorem 1.6. Indeed, recall from (4.24) and (4.39) that, for $x$ and $y$ such that $d(x, y) \leq t^{H}$ and $t<\tau$, we have

$$
\begin{equation*}
p^{\eta}(t, x, y) \geq \frac{C_{1}}{\left|B_{d}\left(x, t^{H}\right)\right|} \quad \text { and } \quad\left|p(t, x, y)-p_{l}^{\eta}(t, x, y)\right| \leq C_{H, V, l} t \tag{4.42}
\end{equation*}
$$

In addition, owing to (1.6), for small $t$ we have

$$
\begin{equation*}
\frac{1}{\left|B_{d}\left(x, t^{H}\right)\right|} \geq \frac{C}{t^{H N / l_{0}}} \tag{4.43}
\end{equation*}
$$

Putting together (4.42) and (4.43), it is easily seen that, when $t$ is small enough, we have

$$
p(t, x, y) \geq \frac{C_{2}}{\left|B_{d}\left(x, t^{H}\right)\right|}
$$

Therefore, the proof of Theorem 1.6 is complete.
REMARK 4.24. With mild assumptions on the Cameron-Martin space, the proof of Theorem 1.5 can be adapted to the case when $B$ is a general Gaussian rough path. Theorem 1.6 does rely on the fractional Brownian scaling in a crucial way. Although it is plausible that Theorem 1.2 may hold for more generic Gaussian rough paths $B$, our proof is based on the fine analysis of the Cameron-Martin structure represented in terms of fractional calculus. It is not clear whether the argument can be adapted to a more generic situation. It is also not clear if there is a natural class of Gaussian rough paths (described in terms of the CameronMartin structure), whose associated control distance functions are equivalent to the intrinsic one associated with the Brownian motion. These are interesting directions to be further investigated.

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