

Spectral loci of Sturm–Liouville operators with polynomial potentials

Alexandre Eremenko and Andrei Gabrielov*

September 27, 2013

Abstract

We consider differential equations $-y'' + P(z, a)y = \lambda y$, where P is a polynomial of the independent variable z depending on a parameter a . The spectral locus is the set of (a, λ) such that the equation has a non-trivial solution tending to zero on two fixed rays in the complex plane. We study the topology of the spectral loci for polynomials P of degree 3 or 4 with respect to z .

MSC: 81Q05, 34M60, 34A05.

Keywords: one-dimensional Schrödinger operators, quasi-exact solvability, PT-symmetry.

We consider differential operators

$$L : y \mapsto -y'' + Py,$$

where the potential P is a polynomial of degree d . When $d \in \{0, 1, 2\}$, the general solution of $Ly = 0$ can be expressed in terms of special functions (elementary, Airy or Weber functions, respectively). The eigenvalue problem for $d = 2$ (harmonic oscillator) plays an important role in quantum mechanics.

We mostly consider the cases $d = 3$ and $d = 4$.

Cubic and quartic oscillators were studied extensively from the very beginning of quantum mechanics, mostly by perturbative methods. Cubic oscillator arises in quantum field theory [24] and in the theory of Painlevé equations [15, 16].

*Both authors are supported by NSF grant DMS-1067886.

We normalize the equation by an affine transformation of the independent variable to obtain: $P(z) = z^d + O(z^{d-2})$, $z \rightarrow \infty$.

Every solution y of $Ly = \lambda y$, for $\lambda \in \mathbf{C}$ is an entire function. Consider the Stokes sectors

$$S_j = \left\{ z : \left| \arg z - \frac{2\pi j}{d+2} \right| < \frac{\pi}{d+2} \right\}, \quad 0 \leq j \leq d+1.$$

For each $j \in \{0, \dots, d+1\}$ a solution y either grows exponentially along all rays from the origin in S_j or tends to zero exponentially along every such ray in S_j .

We choose two non-adjacent sectors, and impose the boundary condition that a solution y of $Ly = \lambda y$ tends to 0 in these sectors. Those λ for which such $y \neq 0$ exists are called eigenvalues, and the corresponding y 's eigenfunctions.

The interest of physicists to such eigenvalue problems originates from the paper [4] where a systematic study of analytic continuation of eigenvalues as parameters vary in the complex plane was made for the first time.

With such boundary conditions, the problem has an infinite discrete spectrum and eigenvalues tend to infinity [20]. To each eigenvalue corresponds a one-dimensional eigenspace.

Suppose that the polynomial potential depends analytically on a parameter $a \in \mathbf{C}^n$. Then the spectral locus Z is defined as the set of all pairs $(a, \lambda) \in \mathbf{C}^{n+1}$ such that the differential equation

$$-y'' + P(z, a)y = \lambda y$$

has a solution y satisfying the boundary conditions. The spectral locus is an analytic hypersurface in \mathbf{C}^{n+1} ; it is the zero-set of an entire function $F(a, \lambda)$ which is called the spectral determinant [20].

The multi-valued function $\lambda(a)$ defined by $F(a, \lambda) = 0$ has the following property: its only singularities are algebraic ramification points, and there are finitely many of them over every compact set in the a -space [8].

Next we discuss connectedness of the spectral locus.

Theorem 1 [9] *For the cubic oscillator*

$$-y'' + (z^3 - az)y = -\lambda y, \quad y(\pm i\infty) = 0, \quad (1)$$

the spectral locus is a smooth irreducible curve in \mathbf{C}^2 .

Theorem 2 [8] *For the even quartic oscillator*

$$-y'' + (z^4 + az^2)y = \lambda y, \quad y(\pm\infty) = 0,$$

the spectral locus consists of two disjoint smooth irreducible curves in \mathbf{C}^2 , one corresponding to even eigenfunctions, another to odd ones.

These theorems can be generalized to polynomials of arbitrary degree if we use all coefficients as parameters [14, 1, 2].

However if we consider a subfamily of all polynomials of given degree, then the spectral locus can be reducible in an interesting way [3].

One example is the family of quasi-exactly solvable quartics L_J

$$-y'' + (z^4 - 2bz^2 + 2Jz)y = \lambda y, \quad y(re^{\pm\pi i/3}) \rightarrow 0, \quad r \rightarrow +\infty. \quad (2)$$

When J is a positive integer, this problem has J elementary eigenfunctions of the form $p(z) \exp(z^3/3 - bz)$, with a polynomial p . The (b, λ) corresponding to these eigenfunctions form the quasi-exactly solvable part Z_J^{QES} of the spectral locus Z_J , and Z_J^{QES} is an algebraic curve.

Theorem 3 [9] *For every positive integer J , the QES spectral locus Z_J^{QES} is a smooth irreducible curve in \mathbf{C}^2 .*

Thus the curve Z_J^{QES} is an irreducible component of Z_J . It is not known whether $Z_J \setminus Z_J^{QES}$ is irreducible.

A similar phenomenon occurs in degree 6: there are one-parametric families of even quasi-exactly solvable sextics [22, 23], and for each such family the quasi-exactly solvable part of the spectral locus is a smooth irreducible algebraic curve [9].

When $J \rightarrow \infty$, an appropriate rescaling of Z_J^{QES} tends to the spectral locus of one-parametric cubic family (1), and a rescaling of the sextic QES spectral locus tends to the spectral locus of the even quartic family [23].

Now we consider Hermitian and PT-symmetric operators. An eigenvalue problem may be invariant under a reflection with respect to a line in the complex z -plane. Without loss of generality, we can take this line to be the real line, and the reflection to be the complex conjugation. Then the coefficients of the potential are real. For the boundary conditions, two cases are possible:

a) Each of the two boundary conditions is preserved by the reflection. In this case the problem is Hermitian.

b) The two boundary conditions are interchanged by the reflection. Such problems are called PT-symmetric. (Physicists prefer to choose the reflection with respect to the imaginary line in this case. PT stands for “parity and time”.)

For example, there is a one-parametric family of PT-symmetric cubics (1), with real parameter a . There are two different two-parametric families of PT-symmetric quartics which we call type I and II:

$$-y'' + (-z^4 + az^2 + cz)y = -\lambda y, \quad y(\pm i\infty) = 0, \quad (3)$$

and

$$-y'' + (z^4 - 2bz^2 + 2Jz)y = \lambda y, \quad y(re^{\pm\pi i/3}) \rightarrow 0, \quad r \rightarrow +\infty. \quad (4)$$

The real spectral locus $Z(\mathbf{R})$ is the subset of the spectral locus Z which consists of points with real coordinates.

We begin with the cubic PT-symmetric spectral locus (1).

Theorem 4 [10] *For every integer $n \geq 0$, there exists a simple curve $\Gamma_n \subset \mathbf{R}^2$, that is the image of a proper analytic embedding of a line, and which has these properties:*

(i) *For every $(a, \lambda) \in \Gamma_n$ problem (1) has an eigenfunction with $2n$ non-real zeros.*

(ii) *The curves Γ_n are disjoint and the real spectral locus of (1) is $\bigcup_{n \geq 0} \Gamma_n$*

(iii) *The map*

$$\Gamma_n \cap \{(a, \lambda) : a \geq 0\} \rightarrow \mathbf{R}_{\geq 0}, \quad (a, \lambda) \mapsto a$$

is a 2-to-1 covering.

(iv) *For $a \geq 0$, $(a, \lambda) \in \Gamma_n$ and $(a, \mu) \in \Gamma_{n+1}$ imply $\mu > \lambda$.*

The following computer-generated plot of the real spectral locus of (1) is taken from Trinh’s thesis [21]. Theorem 4 rigorously establishes some features of this picture. At the points whose abscissas are marked in Fig. 1, pairs of real eigenvalues collide and escape to the complex plane. Theorem 4 proves rigorously that infinitely many such points exist.

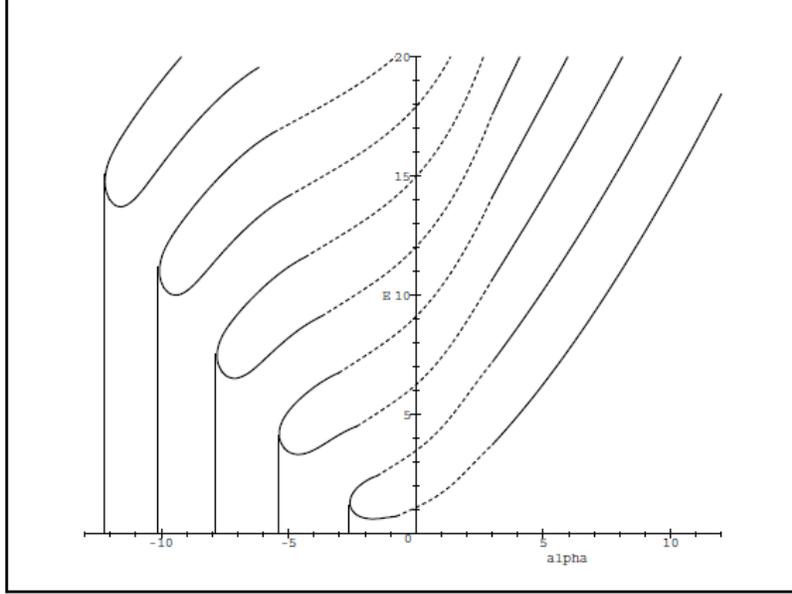


Fig. 1. Real spectral locus for PT -symmetric cubic.

Consider the PT -symmetric quartic family (3) of type I: It is equivalent to the PT -symmetric family

$$-w'' + (z^4 + az^2 +icz)w = \lambda w, \quad w(\pm\infty) = 0,$$

studied by Bender, *et al* [3] and Delabaere and Pham [7].

Theorem 5 [10] *The real spectral locus of (3) consists of disjoint smooth analytic properly embedded surfaces $S_n \subset \mathbf{R}^3$, $n \geq 0$, homeomorphic to a punctured disk. For $(a, c, \lambda) \in S_n$, the eigenfunction has exactly $2n$ non-real zeros. For large a , projection of S_n on the (a, c) plane approximates the region $9c^2 - 4a^3 \leq 0$.*

Numerical computation suggests that the surfaces have the shape of infinite funnels with the sharp end stretching towards $a = -\infty$, $c = 0$, and that the section of S_n by every plane $a = a_0$ is a closed curve.

Theorem 5 implies that this section is compact for large a_0 . The following computer-generated plot is taken from Trinh's thesis:

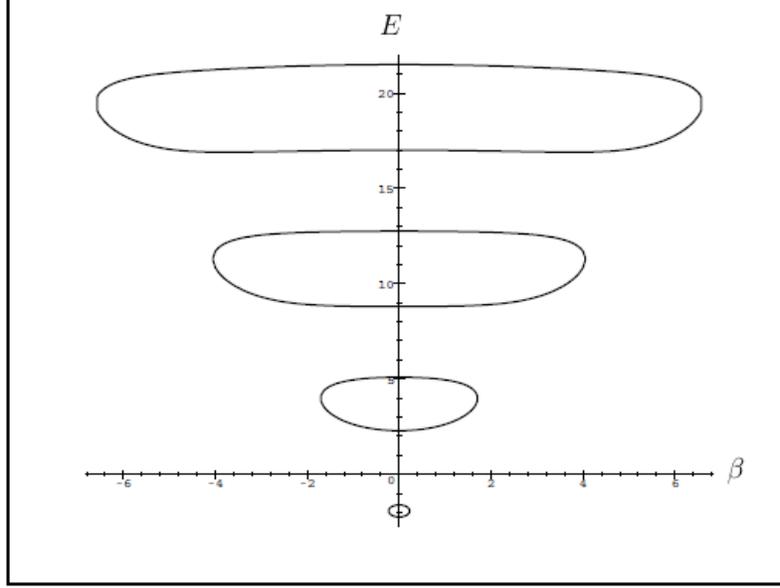


Fig 2. Section of the surfaces S_0, \dots, S_3 by the plane $a = -9$.

The PT-symmetric quartic family of the type II is more complicated, due to the presence of the QES spectrum. Let $Z_J^{QES}(\mathbf{R})$ be the real QES spectral locus of the operator L_J , defined in (2).

Theorem 6 [12] *For $J = n + 1 > 0$, $Z_{n+1}^{QES}(\mathbf{R})$ consists of $[n/2] + 1$ disjoint analytic curves $\Gamma_{n,m}$, $0 \leq m \leq [n/2]$.*

For $(b, \lambda) \in \Gamma_{n,m}$, the eigenfunction has n zeros, $n - 2m$ of them real.

If n is odd, then $b \rightarrow +\infty$ on both ends of $\Gamma_{m,n}$. If n is even, the same holds for $m < n/2$, but on the ends of $\Gamma_{n,n/2}$ we have $b \rightarrow \pm\infty$.

If $(b, \lambda) \in \Gamma_{n,m}$ and $(b, \mu) \in \Gamma_{n,m+1}$ and b is sufficiently large, then $\mu > \lambda$.

It follows from these theorems that in each family, there are infinitely many parameter values where pairs of real eigenvalues collide and escape from the real line to the complex plane.

In the quartic family of the type II, another interesting feature of the real spectral locus is present: for some parameter values the QES spectral locus crosses the rest of the spectral locus. This is called “level crossing”.

Theorem 7 [11] *The points $(b, \lambda) \in Z_J^{QES}$ where the level crossing occurs are the intersection points of Z_J^{QES} with Z_{-J} . For each $J \geq 1$ there are*

infinitely many such points, in general, complex. When J is odd, there are infinitely many level crossing points with $b_k < 0$ and real λ_k . We have

$$b_k \sim -((3/4)\pi k)^{2/3}, \quad k \rightarrow \infty.$$

The only known general result of reality of eigenvalues of PT -symmetric operators is a theorem of K. Shin [19], which for our quartic of type II implies that all eigenvalues are real if $J \leq 0$.

We have the following extensions of this result.

Theorem 8 [13] *For every positive integer J , all non-QES eigenvalues of L_J are real.*

Theorem 9 [13] *All eigenvalues of L_J are real for every real $J \leq 1$ (not necessarily integer).*

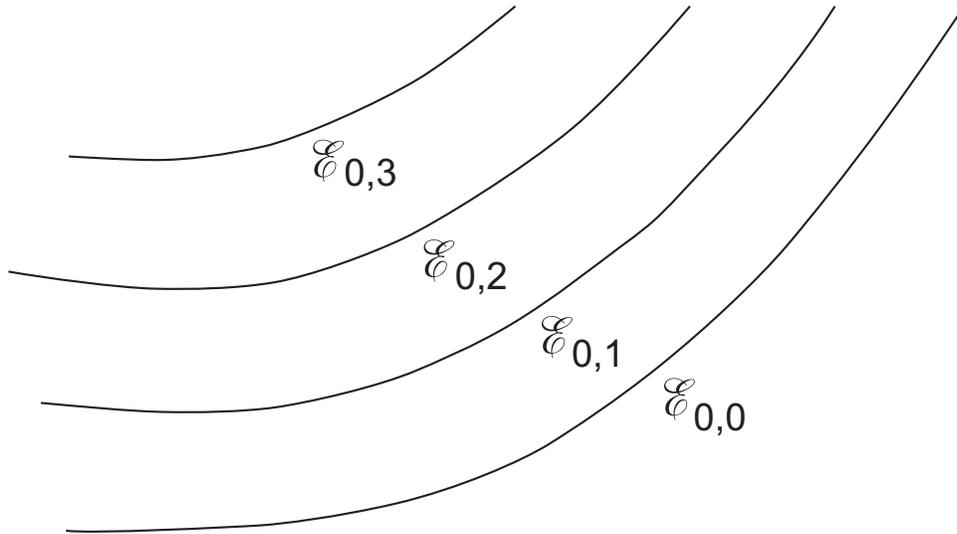


Fig. 3. $Z_0(\mathbf{R})$.

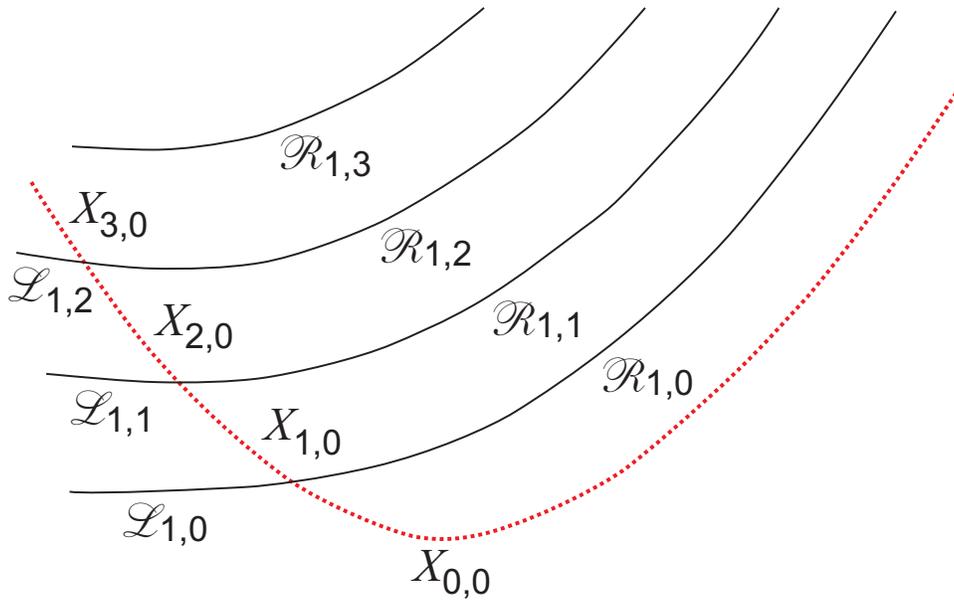


Fig. 4. $Z_1(\mathbf{R})$.

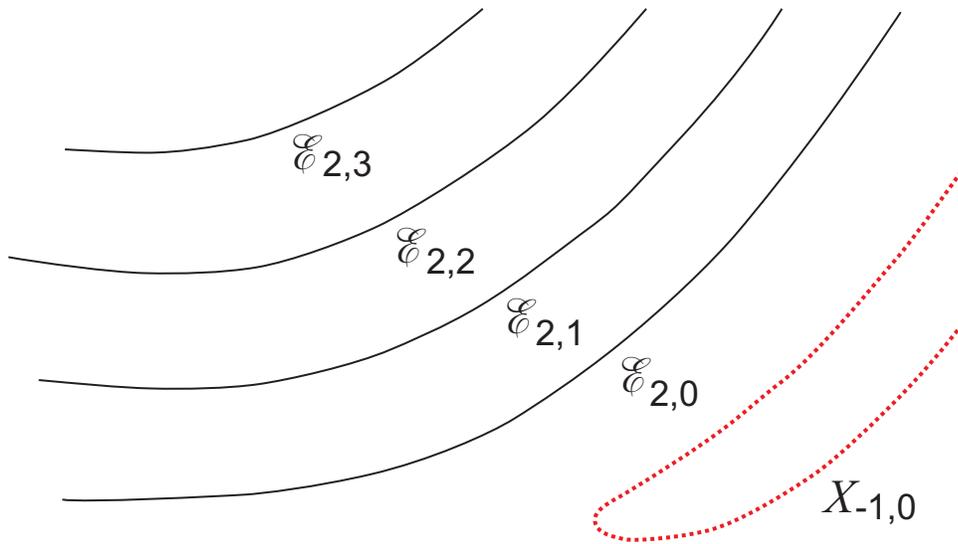


Fig. 5. $Z_2(\mathbf{R})$.

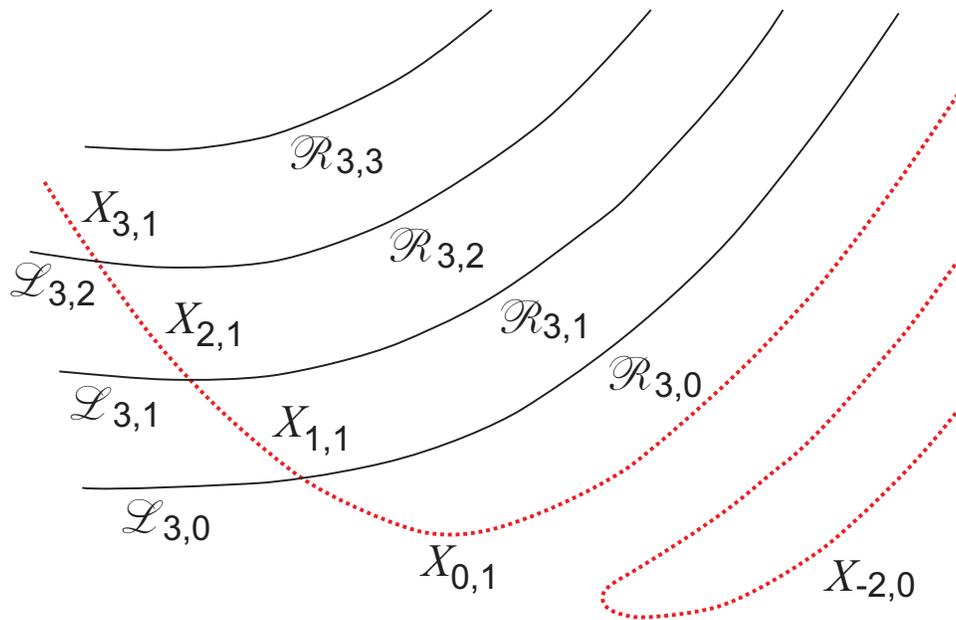


Fig. 6. $Z_3(\mathbf{R})$.

We briefly describe the methods of proofs. There are three main ingredients:

- a) Nevanlinna parametrization of the spectral locus.
- b) Asymptotics at infinity (singular perturbation theory).
- c) Darboux transform of the QES quartic.

Nevanlinna parametrization is the following. Let Z be the spectral locus of the problem

$$-y'' + P(z, a)y = \lambda y, \quad y(z) \rightarrow 0, \quad z \in S_j \cup S_k,$$

where S_j and S_k are non-adjacent Stokes sectors. Let $(a, \lambda) \in Z$, and y_0 an eigenfunction. Let y_1 be a second linearly independent solution. Then $f = y_0/y_1$ satisfies the Schwarz differential equation

$$\frac{f'''}{f'} - \frac{3}{2} \left(\frac{f''}{f'} \right)^2 = -2(P - \lambda).$$

This function f is meromorphic in \mathbf{C} , has no critical points and has $d + 2$ asymptotic values (radial limits along rays), one asymptotic value in each

Stokes sector. Asymptotic values in S_j and S_k are 0. Asymptotic values in adjacent sectors are distinct.

In the opposite direction: if we have a meromorphic function in \mathbf{C} without critical points and with finitely many asymptotic values, then it satisfies a Schwarz equation whose RHS is a polynomial [18]. The degree of this polynomial is the number of asymptotic tracts minus 2. (Asymptotic tracts are in one-to-one correspondence with the logarithmic singularities of the inverse function).

Asymptotic values are meromorphic functions on Z which serve as local parameters. These are called the Nevanlinna parameters; they are simply related to the Stokes multipliers of the linear ODE [20, 15, 16].

Functions f of the above type with given set of asymptotic values $A = \{a_0, \dots, a_{d+1}\}$ have the property that

$$f : \mathbf{C} \setminus f^{-1}(A) \rightarrow \overline{\mathbf{C}} \setminus A$$

is a covering map. For a fixed A such covering map can be completely described by certain combinatorial information, a cell decomposition of the plane. These cell decompositions label the charts of our description of the spectral locus.

It is important that we know exactly which cell decompositions can occur and how the cell decomposition changes when the point (a_0, \dots, a_{d+1}) describes a closed loop in \mathbf{C}^{d+2} . This induces an action of the braid group on the set of special cell decompositions of the plane which can be explicitly computed.

This reduces the problem of parametrization of a spectral locus to combinatorial topology.

For QES operators we use the Darboux transform. Let $-D^2 + V$ be a second order linear differential operator with potential V . Let ϕ_0, \dots, ϕ_n be some eigenfunctions with eigenvalues $\lambda_0, \dots, \lambda_n$. The transformed operator is

$$-D^2 + V - 2 \frac{d^2}{dz^2} \log W(\phi_0, \dots, \phi_n),$$

where W is the Wronski determinant. The eigenvalues of the transformed operator are exactly those eigenvalues of $-D^2 + V$ which are *distinct* from $\lambda_0, \dots, \lambda_n$.

We use the Darboux transform to kill the QES part of the spectrum of L_J and it turns out that the transformed operator is L_{-J} ! Our study of the

QES locus of the quartic family gives the following interesting identities with elementary intergals.

To describe these identities, we address the question studied by Heine and Stieltjes in XIX century. Let h and p be polynomials. When does $y = pe^h$ satisfy a linear differential equation $y'' + Py = 0$ with a polynomial P ?

Theorem 10 *The following conditions are equivalent:*

- a) $p'' + 2p'h'$ is divisible by p ,
- b) $p^{-2}e^h$ has no residues,
- c) zeros of p satisfy the system of equations

$$\sum_{j:j \neq k} \frac{1}{z_k - z_j} = -h'(z_k), \quad 1 \leq k \leq \deg p.$$

Now take $h(z) = z^3/3 - bz$.

Theorem 11 *Let p be a polynomial. All residues of $y = p^{-2}e^{-2h}$ vanish if and only if there exists a constant C and a polynomial q such that*

$$\left(p^2(-z) - \frac{C}{p^2(z)} \right) e^{-2h(z)} = \frac{d}{dz} \left(\frac{q(z)}{p(z)} e^{-2h(z)} \right).$$

Moreover, if this happens then

$$C = (-1)^n 2^{-2n} \frac{\partial}{\partial \lambda} Q_{n+1},$$

where $\lambda = y''/y - z^4 + 2bz^2 - 2(n+1)z$, and $Q_{n+1}(b, \lambda) = 0$ is the equation of the QES spectral locus of L_{n+1} .

This was conjectured in [11] on the basis of calculations with Darboux transform of L_{n+1} and proved by E. Mukhin and V. Tarasov [17].

References

- [1] P. Alexandersson, On eigenvalues of the Schrödinger operator with an even complex-valued polynomial potential, *Comput. Methods Funct. Theory* 12 (2012), no. 2, 465–481.

- [2] P. Alexandersson and A. Gabrielov, On eigenvalues of the Schrödinger operator with a complex-valued polynomial potential, *Comput. Methods Funct. Theory* 12 (2012), no. 1, 119–144.
- [3] C. Bender and S. Boettcher, Quasi-exactly solvable quartic potential, *J. Phys. A* 31 (1998) L273–L277.
- [4] C. Bender and T.-T. Wu, Anharmonic oscillator, *Phys. Rev.*, 184 (1969) 1231–1260.
- [5] M. Crum, Associated Sturm–Liouville systems, *Quart. J. Math.*, 6 (1955) 121–127.
- [6] G. Darboux, Sur la représentation sphérique des surfaces, *C. R. Acad. Sci.*, XCIV, No. 20 (1882) 1343–1345.
- [7] E. Delabaere, F. Pham, Eigenvalues of complex Hamiltonians with PT-symmetry. I,II. *Phys. Lett. A* 250 (1998), no. 1-3, 25–28, 29–32.
- [8] A. Eremenko and A. Gabrielov, Analytic continuation of eigenvalues of a quartic oscillator, *Comm. Math. Phys.*, 287 (2009) 431–457.
- [9] A. Eremenko and A. Gabrielov, Irreducibility of some spectral determinants, arXiv:0904.1714.
- [10] A. Eremenko and A. Gabrielov, Singular perturbation of polynomial potentials in the complex domain with applications to PT-symmetric families, *Moscow Math. J.*, 11 (2011) 473–503.
- [11] A. Eremenko and A. Gabrielov, Quasi-exactly solvable quartic: elementary integrals and asymptotics, *J. Phys. A: Math. Theor.* 44 (2011) 312001.
- [12] A. Eremenko and A. Gabrielov, Quasi-exactly solvable quartic: real algebraic spectral locus, *J. Phys. A: Math. Theor.* 45 (2012) 175205.
- [13] A. Eremenko and A. Gabrielov, Two-parametric PT-symmetric quartic family, *J. Phys. A: Math. Theor.* 45 (2012) 175206.
- [14] H. Habsch, Die Theorie der Grundkurven und Äquivalenzproblem bei der Darstellung Riemannscher Flächen, *Mitt. Math. Seminar Giessen*, 42, 1952.

- [15] D. Masoero, Poles of Integrable Trigonometric and Anharmonic Oscillators. A WKB Approach, *J. Phys. A: Math. Theor.* 43 (2010) 095201
- [16] D. Masoero, Y-System and Deformed Thermodynamic Bethe Ansatz, *Lett. Math. Phys.* 94 (2010) 151–164.
- [17] E. Mukhin and V. Tarasov, On conjectures of A. Eremenko and A. Gabrielov, *Lett. Math. Phys.* 103 (2013), no. 6, 653–663.
- [18] R. Nevanlinna, Über Riemannsche Flächen mit endlich vielen Windungspunkten, *Acta Math.*, 58 (1932) 295–373.
- [19] K. Shin, On the reality of eigenvalues for a class of PT-symmetric oscillators, *Comm. Math. Phys.*, 229 (2002) 543–564.
- [20] Y. Sibuya, Global theory of a second order linear ordinary differential equation with a polynomial coefficient, North-Holland, Amsterdam, 1975.
- [21] D. T. Trinh, Asymptotique et analyse spectrale de l'oscillateur cubique, These, 2002.
- [22] A. Turbiner Quasi-exactly-solvable problems and $sl(2)$ algebra, *Comm. Math. Phys.* 118 (1988), p. 467–474.
- [23] A. Ushveridze, Quasi-exactly solvable models, Inst. of Physics, Bristol and Philadelphia, 1994.
- [24] J. Zinn-Justin and U. Jentschura, Imaginary cubic perturbation: numerical and analytic study, *J. Phys. A*, 43 (2010) 425301.

Purdue University
www.math.purdue.edu/~eremenko
www.math.purdue.edu/~agabriel