RAREFACTION WAVES FOR THE TODA EQUATION VIA NONLINEAR STEEPEST DESCENT

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ABSTRACT. We apply the method of nonlinear steepest descent to compute the long-time asymptotics of the Toda lattice with steplike initial data corresponding to a rarefaction wave.

1. INTRODUCTION

In this paper we consider the doubly infinite Toda lattice

$$
\begin{aligned}
\dot{b}(n,t) &= 2(a(n,t)^2 - a(n-1,t)^2), \\
\dot{a}(n,t) &= a(n,t)(b(n+1,t) - b(n,t)),
\end{aligned}
\qquad (n,t) \in \mathbb{Z} \times \mathbb{R},
$$

(1.1)

with steplike initial profile

$$
\begin{aligned}
a(n,0) &\to a, \quad b(n,0) \to b, \quad \text{as } n \to -\infty, \\
a(n,0) &\to \frac{1}{2} \quad b(n,0) \to 0, \quad \text{as } n \to +\infty,
\end{aligned}
$$

(1.2)

where $a > 0$, $b \in \mathbb{R}$ satisfy the condition

$$
1 < b - 2a.
$$

(1.3)

This inequality implies that the spectra of the left and right background operators H_ℓ and H_r have the following mutual location:

$$
\sup \sigma(H_r) < \inf \sigma(H_\ell).
$$

In the case when $a = \frac{1}{2}$, the initial value problem (1.1)–(1.3) is called rarefaction problem. We keep this name for an arbitrary $a > 0$ and refer to the case $a = \frac{1}{2}$ as the classical rarefaction (CR) problem. The long-time asymptotics of the CR problem were studied rigorously by Deift et al. [5] in 1996 in the transitional region where $\xi := \frac{n}{t} \approx 0$ as $t \to +\infty$. To this end the authors applied the nonlinear steepest descent approach for vector Riemann–Hilbert (RH) problems. Using the same approach, our aim is to study the region $\frac{n}{t} \in (-2a + \varepsilon, -\varepsilon) \cup (\varepsilon, 1 - \varepsilon)$, where $\varepsilon > 0$ is a sufficiently small number. Note that the regions $\frac{n}{t} \in (-\infty, -2a - \varepsilon)$ and $\frac{n}{t} \in (1 + \varepsilon, +\infty)$, which are called the soliton regions, can also be studied by the vector RH approach (see [14] for decaying initial data $a = \frac{1}{2}$, $b = 0$). Although the considerations for the soliton regions in the rarefaction case are somewhat more technical than in the decaying case, they are essentially the same and lead to a sum of solitons on the respective constant background. In our opinion, the classical inverse scattering transform, with the analysis of the Marchenko equation, provides

this result easier ([2, 3, 4, 19]), and consequently we will not study the soliton regions in this paper. On the other hand, the transitional regions $\frac{n}{t} \approx 1$, $\frac{n}{t} \approx 0$ and $\frac{n}{t} \approx -2a$ require further study and are also not the subject of the present paper.

For related results using an ansatz based approach see [15]. For results on the corresponding shock problem see [11, 16, 20] and the references therein. For a general overview we refer to [17].

In summary, we will show that there are four principal sectors with the following asymptotic behavior:

- In the region $n > t$, the solution $\{a(n,t), b(n,t)\}$ is asymptotically close to the constant right background solution $\{\frac{1}{2}, 0\}$ plus a sum of solitons corresponding to the eigenvalues $\lambda_j < -1$.
- In the region $0 < n < t$, as $t \to \infty$ we have

$$(1.4) \qquad a(n,t) = \frac{n}{2t} + O\left(\frac{1}{t}\right), \quad b(n,t) = 1 + \frac{\frac{1}{2} - n}{t} + O\left(\frac{1}{t}\right).$$

- In the region $-2at < n < 0$, as $t \to \infty$ we have

$$(1.5) \qquad a(n,t) = -\frac{n+1}{2t} + O\left(\frac{1}{t}\right), \quad b(n,t) = b - 2a - \frac{n + \frac{3}{2}}{t} + O\left(\frac{1}{t}\right).$$

- In the region $n < -2at$, the solution of (1.1)–(1.3) is close to the left background solution $\{a, b\}$ plus a sum of solitons corresponding to the eigenvalues $\lambda_j > b + 2a$.

The main terms of asymptotics (1.4) and (1.5) are solutions of the Toda lattice equation. The terms $O(t^{-1})$ are uniformly bounded with respect to n, they are differentiable with respect to t, and the first derivatives are of order $O(\frac{n}{t^3})$. In the two middle regions we derive a precise formula for these error terms (see Theorem 5.1 and Proposition 6.1 below).

The following picture demonstrates the expected behavior of the Toda lattice solution in the middle regions. The numerically computed solution in Fig. 1 corresponds to "pure" steplike initial data $a(n,0) = \frac{1}{2}$, $b(n,0) = 0$ for $n \geq 0$ and $a(n,0) = 0.4$, $b(n,0) = 2$ for $n < 0$. The apparent line is due to the fact that neighboring points are very close due to the scaling. We observe that the analytically obtained asymptotics (1.4), (1.5) and the numerically computed asymptotics match well. In particular, in the analytical case the coefficient $b(n,t)$ has a jump in the transition region $\frac{n}{t} \approx 0$ as well.
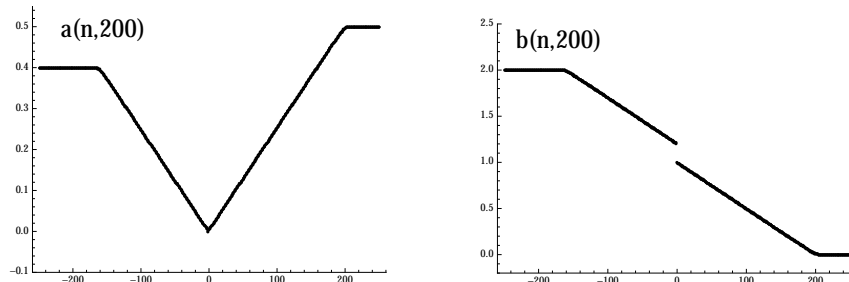

FIGURE 1.   Toda rarefaction problem with non-overlapping background spectra $\sigma(H_\ell) = [1.2, 2.8]$, $\sigma(H_r) = [-1, 1]$; $a = 0.4$, $b = 2$.

To simplify considerations we assume in addition to (1.2) that the initial data decay to their backgrounds exponentially fast

$$(1.6) \qquad \sum_{n=1}^{\infty} e^{\nu n} \left(|a(-n,0) - a| + |b(-n,0) - b| + |a(n,0) - \tfrac{1}{2}| + |b(n,0)| \right) < \infty,$$

where $\nu > 0$ is an arbitrary small number. This condition allows to continue the right reflection coefficient analytically in a small vicinity of the respective spectrum.

2. Statement of the Riemann–Hilbert problem

Let us first recall some elementary facts from scattering theory for Jacobi operators with steplike backgrounds from [7, 8, 9, 10] (see also [18] Chapter 10 for general background). The spectrum of the Jacobi operator $H(t)$ associated with the equation

$$(2.1) \quad H(t)y(n) := a(n-1,t)y(n-1) + b(n,t)y(n) + a(n,t)y(n+1) = \lambda y(n), \ \lambda \in \mathbb{C},$$

consists of two intervals $[-1,1] \cup [b-2a, b+2a]$ of continuous spectrum with multiplicity one, plus a finite number of eigenvalues, $\{\lambda_j\}_{j=1}^{N} \subset \mathbb{R} \setminus ([-1,1] \cup [b-2a, b+2a])$. In addition to the spectral parameter λ we will use two other parameters z and ζ, connected with λ by the Joukovski transformation

$$(2.2) \qquad \lambda = \frac{1}{2}\left(z + z^{-1}\right) = b + a\left(\zeta + \zeta^{-1}\right), \quad |z| \leq 1, \quad |\zeta| \leq 1.$$

Introduce the Jost solutions $\psi(z,n,t)$, $\psi_\ell(z,n,t)$ of (2.1) with asymptotic behaviour

$$\lim_{n\to\infty} z^{-n}\psi(z,n,t) = 1, \quad |z| \leq 1; \qquad \lim_{n\to-\infty} \zeta^n\psi_\ell(z,n,t) = 1, \quad |\zeta| \leq 1.$$

Denote

$$q_1 = z(b-2a), \quad q_2 = z(b+2a), \quad z_j = z(\lambda_j), \quad j = 1, \dots, N,$$

where

$$z(\lambda) = \lambda - \sqrt{\lambda^2 - 1}.$$

The points $z = -1$ and $z = 1$ correspond to the edges of the spectrum of the right background Jacobi operator, and q_2 and q_1 correspond to the respective edges of the left background operator. We will call the points z_j discrete spectrum. Denote $\mathbb{T} = \{z : |z| = 1\}$ and $\mathbb{D} = \{z : |z| < 1\}$. The map $z \mapsto \lambda$ is one-to-one between the closed domains $\mathfrak{D} := \text{clos}(\mathbb{D} \setminus [q_2, q_1])$ and $\mathcal{D} := \text{clos}(\mathbb{C} \setminus ([-1,1] \cup [b-2a, b+2a]))$. We treat closure as adding to the boundary the points of the upper and lower sides along the cuts, while considering them as distinct points. Since the function $\zeta^n \psi_\ell(z)$ is in fact an analytical function of ζ, it takes complex conjugated values on the sides of the cut along the interval $I := [q_2, q_1]$, which we denote as $I \pm i0$. Note that $z \in I - i0$ corresponds to the arc $\zeta \in \{|\zeta| = 1, \text{Im}\,\zeta < 0\}$. The Jost solution $\psi_\ell(z)$ takes equal real values at $z, z^{-1} \in \mathbb{T}$, which yields the respective properties of the scattering matrix as a function of z. This matrix consist of the right (resp., left) reflection coefficient $R(z,t)$ (resp. $R_\ell(z,t)$), defined for $|z| = 1$ (resp., $|\zeta| = 1$), and the transmission coefficients $T(z,t)$ and $T_\ell(z,t)$ defined on $\mathfrak{D}$. They are connected by the scattering relations

$$T(z,t)\psi_\ell(z,n,t) = \overline{\psi(z,n,t)} + R(z,t)\psi(z,n,t), \quad |z| = 1,$$
$$T_\ell(z,t)\psi(z,n,t) = \overline{\psi_\ell(z,n,t)} + R_\ell(z,t)\psi_\ell(z,n,t), \quad z \in I \pm i0.$$

Here we list only those properties of the scattering data which will be used in the present paper. Moreover, let z_j, $1 \leq j \leq N$, (note $|z_j| < 1$) be the eigenvalues and set

$$\gamma_j(t) := \Big(\sum_{\mathbb{Z}} \psi^2(z_j, n, t) \Big)^{-2}, \quad j = 1, \ldots, N,$$

(2.3) $$\chi(z, t) := -\lim_{\varepsilon \to 0} T(z - i\varepsilon, t) T_\ell(z - i\varepsilon, t), \quad z \in I.$$

The set of the right scattering data

(2.4) $$\{R(z, t), z \in \mathbb{T}; \ \chi(z, t), z \in I; \ (z_j, \gamma_j(t)), 1 \leq j \leq N\}$$

defines the solution of the Toda lattice uniquely. Under condition (1.6) it has the following properties ([5, 8]):

- The function $R(z, t)$ is continuous and $R(z^{-1}, t) = \overline{R(z, t)} = R^{-1}(z, t)$ for $z \in \mathbb{T}$. We have $R(-1) = -1$ if $z = -1$ is non-resonant[1] and $R(-1) = 1$ if $z = -1$ is resonant. The function $R(z)$ can be continued analytically in the annulus $e^{-\nu} < |z| < 1$.
- The right transmission coefficient $T(z, t)$ can be restored uniquely from (2.4) for $z \in \mathfrak{D}$. It is a meromorphic function with simple poles at z_j.
- The function $\chi(z, t)$ is continuous for $z \in (q_2, q_1)$ and vanishes at q_i, $i = 1, 2$, iff q_i is a non-resonant point. If q_i is a resonant point, $\chi(z) = (z - q_i)^{-1/2}$. The transmission coefficient has the same behavior at q_i as $\chi(z, t)$.

On the set $\mathfrak{D}$ we define a vector-valued function $m(z) = (m_1(z, n, t), m_2(z, n, t))$

(2.5) $$m(z, n, t) = \big(T(z, t)\psi_\ell(z, n, t)z^n, \quad \psi(z, n, t)z^{-n}\big).$$

Lemma 2.1 ([11]). *The components of $m(z, n, t)$ have the following asymptotical behavior as $z \to 0$*

(2.6)
$$m_1(z, n, t) = \prod_{j=n}^{\infty} 2a(j, t)\Big(1 + 2z \sum_{m=n}^{\infty} b(m, t)\Big) + O(z^2),$$

$$m_2(z, n, t) = \prod_{j=n}^{\infty} (2a(j, t))^{-1}\Big(1 - 2z \sum_{m=n+1}^{\infty} b(m, t)\Big) + O(z^2).$$

Evidently $m_1(z)$ is a meromorphic function with poles at z_j. Let us extend m to the set $\{z : |z| > 1\} \setminus I^*$, $I^* := [q_2^{-1}, q_1^{-1}]$, by $m(z^{-1}) = m(z)\sigma_1$, where σ_1 is the first Pauli matrix. Recall that the Pauli matrices are given by

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix},$$

we will also use them in abbreviations as for example

$$[d(z)]^{-\sigma_3} := \begin{pmatrix} d^{-1}(z) & 0 \\ 0 & d(z) \end{pmatrix}.$$

Moreover, $m_2(z)$ is a meromorphic function in $\{z : |z| > 1\} \setminus I^*$ with poles at z_j^{-1}.

By definition the vector function $m(z)$, $z \in \mathbb{C}$, has jumps along the unit circle and along the intervals I and I^*. The statement of the respective Riemann-Hilbert problem with pole conditions is given in [5]. We will not formulate this problem

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[1]The point $\hat{z} \in \{-1, 1, q_1, q_2\}$ is called a resonant point if $W(\hat{z}, t) = 0$, where $W(z, t) := a(n-1, t)(\psi_\ell(z, n-1, t)\psi(z, n, t) - \psi_\ell(z, n, t)\psi(z, n-1, t))$ is the Wronskian of the Jost solutions. If $W(\hat{z}, t) \neq 0$, then $\hat{z}$ is non-resonant.

here, but instead give an equivalent statement which is valid in the domain $\xi := \frac{n}{t} \in (-2a, 0) \cup (0, 1)$ we are interested in. In this domain we can reformulate the initial meromorphic RH problem as a holomorphic RH problem as [5, 14]. We skip the details and only provide a brief outline below.

Throughout this paper, $m_+(z)$ (resp. $m_-(z)$) will denote the limit of $m(p)$ as $p \to z$ from the positive (resp. negative) side of an oriented contour Σ. Here the positive (resp. negative) side is the one which lies to the left (resp. right) as one traverses the contour in the direction of its orientation. Using this notation implicitly assumes that these limits exist in the sense that $m(z)$ extends to a continuous function on the boundary. Moreover, all contours are symmetric with respect to the map $z \mapsto z^{-1}$, i.e., they contain with each point z also z^{-1}. The orientation on these contours should be chosen in such a way that the following symmetry is preserved for the jump matrix of the vector RH problem and for its solution.

Symmetry condition. *Let $\hat{\Sigma}$ be a symmetric oriented contour, which contains a symmetric arc $\Sigma = \hat{\Sigma} \cap \mathbb{T}$. Then the jump matrix of the vector problem $m_+(z) = m_-(z)v(z)$ satisfies*

$$(2.7) \quad (v(z))^{-1} = \sigma_1 v(z^{-1})\sigma_1, \quad as\ z \in \Sigma; \quad v(z) = \sigma_1 v(z^{-1})\sigma_1, \quad as\ z \in \hat{\Sigma} \setminus \Sigma.$$

Moreover,

$$(2.8) \qquad\qquad m(z) = m(z^{-1})\sigma_1, \quad as\ z \in \mathbb{C} \setminus \hat{\Sigma}.$$

If we reverse the orientation of the contour, we will denote this contour by $\hat{\Sigma}^{-1}$.

Most of our transformations are conjugations with diagonal matrices, so it is convenient to use the following

Lemma 2.2 (Conjugation, [13]). *Let m be the solution on $\mathbb{C}$ of the RH problem $m_+(z) = m_-(z)v(z)$, $z \in \hat{\Sigma}$, which satisfies the symmetry condition as above. Let $d : \mathbb{C} \setminus \Sigma \to \mathbb{C}$ be a sectionally analytic function. Set*

$$(2.9) \qquad\qquad \tilde{m}(z) = m(z)[d(z)]^{-\sigma_3},$$

then the jump matrix of the problem $\tilde{m}_+ = \tilde{m}_- \tilde{v}$ is given by

$$\tilde{v} = \begin{cases} \begin{pmatrix} v_{11} & v_{12}d^2 \\ v_{21}d^{-2} & v_{22} \end{pmatrix}, & p \in \hat{\Sigma} \setminus \Sigma, \\[12pt] \begin{pmatrix} \frac{d_-}{d_+}v_{11} & v_{12}d_+d_- \\ v_{21}d_+^{-1}d_-^{-1} & \frac{d_+}{d_-}v_{22} \end{pmatrix}, & p \in \Sigma. \end{cases}$$

If d satisfies $d(z^{-1}) = d(z)^{-1}$ for $z \in \mathbb{C} \setminus \Sigma$, then the transformation (2.9) respects the symmetry condition.

Recall now that the behavior of the solution of the RH problem is determined mostly by the behavior of the phase function

$$(2.10) \qquad\qquad \Phi(z) = \Phi(z, \xi) = \frac{1}{2}\left(z - z^{-1}\right) + \xi \log z.$$

In the domain $\xi \in (0, 1)$ which we study in more detail, part of the eigenvalues lie in the domain $\operatorname{Re}\Phi(z) > 0$ (namely the points $z_k \in (-1, 0)$) while the remaining eigenvalues belong to the set $\operatorname{Re}\Phi(z) < 0$. The pole conditions at the eigenvalues

are given by ([14])

$$\operatorname{Res}_{z_j} m(z) = \lim_{z \to z_j} m(z) \begin{pmatrix} 0 & 0 \\ -z_j \gamma_j \mathrm{e}^{2t\Phi(z_j)} & 0 \end{pmatrix},$$

$$\operatorname{Res}_{z_j^{-1}} m(z) = \lim_{z \to z_j^{-1}} m(z) \begin{pmatrix} 0 & z_j^{-1} \gamma_j \mathrm{e}^{2t\Phi(z_j)} \\ 0 & 0 \end{pmatrix}.$$

Let δ be a sufficiently small number such that the circles $\mathbb{T}_j = \{z : |z - z_j| = \delta\}$ around the eigenvalues do not intersect and lie away from the set $\mathbb{T} \cup I$ (the precise value of δ will be chosen later). We consider the circles as contours with counterclockwise orientation. Denote their images under the map $z \mapsto z^{-1}$ as $\mathbb{T}_j^*$. These curves are not circles, but they surround the points z_j^{-1} with minimal distance from the curve to z_j given by $\frac{\delta}{z_j(z_j - \delta)}$. The curves $\mathbb{T}_j^*$ are oriented counterclockwise as well. Now redefine $m(z)$ inside these curves according to

$$m^{\mathrm{hol}}(z) = \begin{cases} m(z)A(z), & \{|z - z_j| < \delta\} \cup \{|z^{-1} - z_j| < \delta\}\}, \\ m(z), & \text{else}, \end{cases}$$

where

$$A(z) = \begin{pmatrix} 1 & 0 \\ \frac{z_j \gamma_j \mathrm{e}^{2t\Phi(z_j)}}{z - z_j} & 1 \end{pmatrix}, \quad |z - z_j| < \delta, \text{ and } A(z^{-1}) = \sigma_1 A(z)\sigma_1.$$

Then the vector function m^{hol} is holomorphic near z_j, z_j^{-1} and has an additional jump,

$$m_+^{\mathrm{hol}}(z) = m_-^{\mathrm{hol}}(z)A(z), \quad z \in \mathbb{T}_j \cup \mathbb{T}_j^*.$$

Let

(2.11)
$$P(z) = \prod_{z_k \in (-1,0)} |z_k| \frac{z - z_k^{-1}}{z - z_k}.$$

be the Blaschke product which corresponds to the eigenvalues z_k which lie in the domain $\operatorname{Re}\Phi(z) > 0$ (if any). Introduce a second matrix

$$A_1(z) = \begin{pmatrix} 1 & \frac{z - z_k}{z_k \gamma_k \mathrm{e}^{t\Phi(z_k)}} \\ -\frac{z_k \gamma_k \mathrm{e}^{t\Phi(z_k)}}{z - z_k} & 0 \end{pmatrix} [P(z)]^{-\sigma_3}, \quad |z - z_k| < \delta,$$

for $z_k \in (-1, 0)$, and continue its definition in the interior of $\mathbb{T}_k^*$ by the symmetry condition

$$A_1(z) = \sigma_1 A_1(z^{-1})\sigma_1, \quad |z^{-1} - z_k| < \delta.$$

At all other points of $\mathbb{C}$ away from the curves $\mathbb{T}_k \cup \mathbb{T}_k^*$, $z_k \in (-1, 0)$, define $A_1(z)$ by $A_1(z) = [P(z)]^{-\sigma_3}$. Since the Blaschke product satisfies $P(z^{-1}) = P^{-1}(z)$ we have $A_1(z^{-1}) = \sigma_1 A_1(z)\sigma_1$ everywhere. Now conjugate the vector m^{hol} by setting

$$m^{\mathrm{ini}}(z) = m^{\mathrm{hol}}(z)A_1(z), \qquad z \in \mathbb{C}.$$

Then the vector $m^{\mathrm{ini}}(z)$ is a holomorphic function in

$$\mathbb{C} \setminus \{\mathbb{T} \cup I \cup I^* \cup \mathbb{T}^\delta\}, \quad \mathbb{T}^\delta := \bigcup_{j=1}^N \mathbb{T}_j \cup \mathbb{T}_j^*,$$

and solves the jump problem $m_+^{\text{ini}}(z) = m_-^{\text{ini}}(z)B(z)$, $z \in \mathbb{T}^\delta$, in neighborhoods of the discrete spectrum, where (cf. [14])

$$(2.12) \qquad B(z) = \begin{cases} \begin{pmatrix} 1 & \frac{(z-z_k)P^2(z)}{z_k \gamma_k e^{t\Phi(z_k)}} \\ 0 & 1 \end{pmatrix}, & \text{for } z \in \mathbb{T}_k, \quad z_k \in (-1,0), \\ \sigma_1 B(z^{-1})\sigma_1, & \text{for } z \in \mathbb{T}_k^*, \quad z_k \in (-1,0), \\ [P(z)]^{\sigma_3} A(z)[P(z)]^{-\sigma_3}, & \text{for } z \in \mathbb{T}_k \cup \mathbb{T}_k^*, \quad z_k \in (0,1). \end{cases}$$

We observe that the matrix $B(z)$ respects the symmetry property and

$$\|B(z) - \mathbb{I}\| \le C e^{-t|\inf_k \operatorname{Re}\Phi(z_k)|}, \quad z \in \mathbb{T}^\delta, \quad 0 < \xi < 1.$$

Here the matrix norm is to be understood as the maximum of the absolute value of its elements.

Consider the contour $\Gamma = \mathbb{T} \cup I \cup I^* \cup \mathbb{T}^\delta$, where the unit circle $\mathbb{T}$ is oriented counterclockwise and the intervals I, I^* are oriented towards the center of the circle. Continue the function (2.3) to I^* by $\chi(z) = -\chi(z^{-1})$. Then the following proposition is valid (cf. [5], [14]).

Proposition 2.3. *Suppose that the initial data of the Cauchy problem* (1.1)–(1.3), *satisfy* (1.6). *Let* $\{R(z), |z| = 1; \chi(z), z \in I; (z_j, \gamma_j), 1 \le j \le N\}$ *be the right scattering data of the operator* $H(0)$. *Suppose that the operator has no resonances at the spectral edges* $b - 2a$, $b + 2a$. *Let* $\xi = \frac{n}{t} \in (0, 1)$. *Then the vector-valued function* $m^{\text{ini}}(z) = m^{\text{ini}}(z, n, t)$, *connected with the initial function* (2.5) *by*

$$m^{\text{ini}}(z) = m(z)[P(z)]^{-\sigma_3}, \quad as \ z \in \{z : |z^{\pm 1} - z_j| > \delta, \ 1 \le j \le N\},$$

is the unique solution of the following vector Riemann–Hilbert problem: Find a vector-valued function $m^{\text{ini}}(z)$ *which is holomorphic away from* Γ, *continuous up to the boundary, and satisfies:*

I. The jump condition $m_+^{\text{ini}}(z) = m_-^{\text{ini}}(z)v(z)$, *where*

$$v(z) = \begin{cases} \begin{pmatrix} 0 & -P^2(z)\overline{R(z)}e^{-2t\Phi(z)} \\ P^{-2}(z)R(z)e^{2t\Phi(z)} & 1 \end{pmatrix}, & z \in \mathbb{T}, \\ \begin{pmatrix} 1 & 0 \\ P^{-2}(z)\chi(z)e^{2t\Phi(z)} & 1 \end{pmatrix}, & z \in I, \\ \begin{pmatrix} 1 & P^2(z)\chi(z)e^{-2t\Phi(z)} \\ 0 & 1 \end{pmatrix}, & z \in I^*, \\ B(z), & z \in \mathbb{T}^\delta. \end{cases}$$

Here the phase function $\Phi(z) = \Phi(z, n/t)$ *is given by* (2.10), *the matrix* $B(z)$ *by* (2.12), *the function* $P(z)$ *by* (2.11), *and* $\chi(z)$ *by* (2.3).
II. The symmetry condition $m^{\text{ini}}(z^{-1}) = m^{\text{ini}}(z)\sigma_1$.
III. The normalization condition

$$(2.13) \qquad m^{\text{ini}}(0) = (m_1^{\text{ini}}, m_2^{\text{ini}}), \quad m_1^{\text{ini}} \cdot m_2^{\text{ini}} = 1, \quad m_1^{\text{ini}} > 0.$$

Remark 2.4. We assume that the points $b - 2a$ and $b + 2a$ are non-resonant. This means that the initial vector function m has continuous limits on I, I^* and that $\chi(z)$ is bounded there (otherwise $T(z) \sim (z - q_j)^{-1/2}$ and both m and the jump matrix have singularities). When $\xi \in (0, 1)$ we have $I \subset \{z : \operatorname{Re}\Phi(z) = \Phi(z) < 0\}$.

Moreover, the matrix $v(z)$ satisfies the symmetry property $v(z) = \sigma_1 v(z^{-1})\sigma_1$ on I and I^*. It also admits the estimate

$$\|v(z) - \mathbb{I}\|_{L^\infty(I \cup I^*)} \leq C \mathrm{e}^{-t|\Phi(q_1)|}.$$

We omit the proof of Proposition 2.3 which is essentially the same as in [5, 14]. Note that the normalization condition $m_1^{\mathrm{ini}} > 0$ holds since it holds for the initial function m and by definition, $P(0) > 0$. Uniqueness of the solution can be proved as in [1].

3. Reduction to the model problem

In this section we perform three conjugation-deformation steps which reduce the RH problem I–III to a simple jump problem on an arc of the circle with a constant jump matrix, plus jump matrices which are small with respect to t. The jump problem with the constant matrix can be solved explicitly. Note that the property $|R(z)| = 1$ for $z \in \mathbb{T}$ does not allow us to apply a standard lower-upper factorization of the jump matrices on an arc of $\mathbb{T}$ and subsequent "lens" machinery near this arc (see [14]). This is why we first need to find a suitable g-function ([6]).

Step 1. In this step we replace the phase function by a function with "better" properties. The g-function has the same asymptotics (up to a constant term) as $\Phi(z)$ for $z \to 0$ and $z \to \infty$ and the same oddness property $g(z^{-1}) = -g(z)$. In addition, it has the convenient property that the curves separating the domains with different signs of $\operatorname{Re} g(z)$ cross at points $z_0(\xi) \in \mathbb{T}$ and $\overline{z}_0(\xi)$, where $g(z_0) = g(\overline{z}_0) = 0$. A second helpful property of the g-function is that it has a jump along the arc connecting z_0 and $\overline{z}_0$ which satisfies $g_+(z) = -g_-(z) > 0$. This simplifies further transformations, because with this property and Lemma 2.2 we do not need the lens machinery around this arc. Note that the point z_0 does not coincide with the stationary phase point of Φ. Recall that for $\xi \in (0,1)$, the curves $\operatorname{Re}\Phi(z) = 0$ intersect at the symmetric points $-\xi \pm \sqrt{\xi^2 - 1} \in \mathbb{T}$, which are the stationary points of Φ. That is, the stationary phase point corresponds to the angle $\phi_0 \in (0,\pi)$ where $\cos\phi_0 = -\xi$. Set

$$(3.1) \qquad z_0 = \mathrm{e}^{\mathrm{i}\theta_0}, \quad \text{where } \cos\theta_0 = 1 - 2\xi, \quad \theta_0 \in (0,\pi).$$

Introduce

$$(3.2) \qquad g(z) = \frac{1}{2}\int_{z_0}^{z} \sqrt{\left(1 - \frac{1}{sz_0}\right)\left(1 - \frac{z_0}{s}\right)}(1+s)\frac{ds}{s},$$

where $\sqrt{s} > 0$ for $s > 0$, and the cut of the square root in (3.2) is taken between the points z_0 and $\overline{z}_0$ along the arc

$$\Sigma = \{z \in \mathbb{T} : \operatorname{Re} z \leq \operatorname{Re} z_0 = \cos\theta_0\}.$$

We orient this arc in the same way as $\mathbb{T}$, i.e., from z_0 to $\overline{z}_0$.

Lemma 3.1. *The function $g(z)$ satisfies the following properties:*

(a) $\Phi(z) - g(z) = -K(\xi) \in \mathbb{R}$ *as $z \to \infty$;*
(b) $g(\overline{z}_0) = 0$;
(c) $g(z^{-1}) = -g(z)$ *as $z \in \mathbb{C} \setminus \Sigma$;*
(d) *It has a jump along the arc Σ with $g_+(z) = -g_-(z) > 0$ as $z \neq z_0^{\pm 1}$;*

(e) *In a vicinity of z_0,*

$$g(z) = C(\theta_0)(z - z_0)^{3/2} e^{i(\frac{\pi}{4} - \frac{3\theta_0}{2})}(1 + o(1)), \quad C(\theta_0) > 0.$$

In particular, $g(z) = C(\theta_0)(|z - z_0|)^{3/2}(1 + o(1))$ as $z \in \Sigma$.

Proof. We first prove that our choice of z_0 yields $\Phi(z) - g(z) = O(1)$ as $z \to \infty$. To match the asymptotics of $g(z)$ and $\Phi(z)$ in (2.10) we compute

$$\frac{d}{dz}g(z) = \frac{1}{2}\sqrt{\left(1 - \frac{1}{zz_0}\right)\left(1 - \frac{z_0}{z}\right)}\frac{1+z}{z} = \frac{1}{2}\left(1 - \frac{1}{2zz_0} - \frac{z_0}{2z} + \frac{1}{z}\right) + O\left(\frac{1}{z^2}\right),$$

as $z \to \infty$, and hence

$$g(z) = \frac{z}{2} + \frac{1}{2}\left(1 - \frac{z_0 + z_0^{-1}}{2}\right)\log z + O(1).$$

Now choose

$$\frac{z_0 + z_0^{-1}}{2} = 1 - 2\xi,$$

which implies that $\cos\theta_0 = 1 - 2\xi$ as desired. To prove property (b), substitute $s = e^{i\theta}$ and $z_0 = e^{i\theta_0}$ in (3.2), then

$$g(\overline{z_0}) = i\int_{\theta_0}^{2\pi - \theta_0}\sqrt{(1 - e^{-i(\theta_0 + \theta)})(e^{i\theta} - e^{i\theta_0})}\,\frac{e^{\frac{i\theta}{2}} + e^{-\frac{i\theta}{2}}}{2}d\theta$$

$$= i\sqrt{2}\int_{\theta_0}^{2\pi - \theta_0}\sqrt{\cos\theta - \cos\theta_0}\cos\tfrac{\theta}{2}d\theta.$$

For $\theta \in [\theta_0, 2\pi - \theta_0]$, we have that $f(\theta) = \cos\theta_0 - \cos\theta > 0$ and f is even with respect to α, $f(\pi - \alpha) = f(\pi + \alpha)$. Moreover, $\cos(\frac{\pi + \alpha}{2})$ is odd with respect to α, $\cos(\frac{\pi - \alpha}{2}) = -\cos(\frac{\pi + \alpha}{2})$. Hence the substitution $\theta = \pi + \alpha$ with $-\pi + \theta_0 \leq \alpha < \pi - \theta_0$ yields

$$g(\overline{z_0}) = \sqrt{2}\int_{\theta_0 - \pi}^{\pi - \theta_0}\sqrt{f(\pi + \alpha)}\cos(\tfrac{\pi + \alpha}{2})d\alpha = 0.$$

Now it is straightforward to see (c), because the integrand can be represented as

(3.3)
$$\frac{d}{ds}g(s) = \frac{1}{s}\sqrt{\frac{s + s^{-1}}{2} - \frac{z_0 + z_0^{-1}}{2}}\sqrt{\frac{s + s^{-1}}{2} + 1} =: \frac{h(s)}{s}.$$

Since $h(s^{-1}) = h(s)$, one obtains by replacing $s = t^{-1}$

$$g(z^{-1}) = \int_{z_0}^{z^{-1}}\frac{h(s)ds}{s} = -\int_{\overline{z_0}}^{z}\frac{h(t)dt}{t} = -\int_{z_0}^{z}\frac{h(t)dt}{t} - g(\overline{z_0}) = -g(z).$$

For property (d), note that due to their equal asymptotical behaviour, the signature table for g as $z \to 0$ or $z \to \infty$ is the same as the signature table for Φ (see Fig. 2). The line $\mathbb{T} \setminus \Sigma$ corresponds to $\operatorname{Re} g = 0$. Indeed, if $z_1 = e^{i\theta_1}$ and $\operatorname{Re} z_1 > \operatorname{Re} z_0$, then

$$g(z_1) = i\sqrt{2}\int_{\theta_0}^{\theta_1}\sqrt{\cos\theta - \cos\theta_0}\cos(\tfrac{\theta}{2})d\theta \in i\mathbb{R}.$$

The function $g(z)$ has a jump along the contour Σ, but the limiting values are real (compare with the proof of (b)). Thus $\operatorname{Re} g_+ = g_+ > 0$ (because this limit is taken from the domain where $\operatorname{Re} g > 0$). Respectively, $\operatorname{Re} g_- = g_- < 0$ and $g_+ = -g_-$.

Now we are ready to finish the proof of property (a). Since $g(z)$ satisfies (c) we have $g(1) = 0$. By use of (3.3) and (2.2), and taking into account that

$$\frac{dz}{z} = -\frac{d\lambda}{\sqrt{\lambda^2 - 1}}, \quad \Phi(z) = -\int_1^{\lambda(z)} \frac{x + \xi}{\sqrt{x^2 - 1}} dx,$$

we obtain for $z \to +0$, that is when $\sqrt{\lambda^2 - 1} > 0$ as $\lambda > 1$, the asymptotic behaviour

(3.4) $$\Phi(z) - g(z) = K(\xi) + k(\xi)z + O(z^2), \quad \text{as } z \to 0,$$

where

(3.5) $$K(\xi) = \int_1^\infty \frac{\sqrt{(x - 1 + 2\xi)(x + 1)} - x - \xi}{\sqrt{x^2 - 1}} dx \in \mathbb{R}, \quad k(\xi) = 1 - 2\xi + \xi^2.$$

Differentiating (3.5) we also get

(3.6) $$\frac{d}{d\xi} K(\xi) = -\log \xi.$$

To prove (e) we decompose (3.2) in a vicinity of z_0 and integrate. Since $\arg(z - z_0) = (-\frac{3\pi}{2} + \theta_0)(1 + o(1))$ for $z \in \Sigma$ in a small vicinity of z_0, then $\arg(z - z_0)^{3/2} + \frac{\pi}{4} - \frac{3\theta_0}{2} = -2\pi(1 + o(1))$, which implies the second claim of (e). □

The signature table for $g(z)$ is given in Fig. 2. With this description of the g-
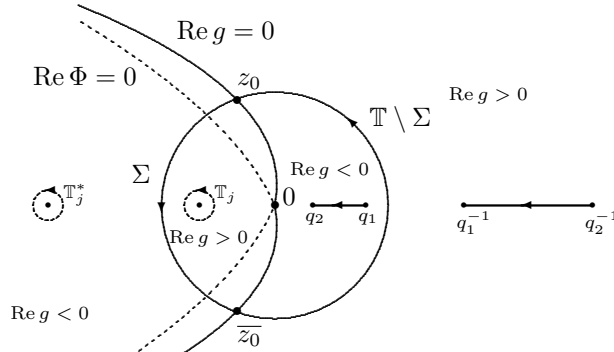

FIGURE 2. Signature table for $g(z)$.

function we introduce the function $d(z) = e^{t(\Phi(z) - g(z))}$. It satisfies the conditions of Lemma 2.2. Let $m^{\text{ini}}(z)$ be the solution of the RH problem I–III. Set

$$m^{(1)}(z) = m^{\text{ini}}(z)[d(z)]^{-\sigma_3},$$

then this vector solves the jump problem $m_+^{(1)}(z) = m_-^{(1)}(z)v^{(1)}(z)$ with

$$v^{(1)}(z) = \begin{cases} \begin{pmatrix} 0 & -\overline{\mathcal{R}(z)}e^{-2tg(z)} \\ \mathcal{R}(z)e^{2tg(z)} & 1 \end{pmatrix}, & z \in \mathbb{T} \setminus \Sigma, \\ \begin{pmatrix} 0 & -\overline{\mathcal{R}(z)} \\ \mathcal{R}(z) & e^{-2tg_+(z)} \end{pmatrix}, & z \in \Sigma, \\ E(z), & z \in \Xi, \end{cases}$$

where the following notations are introduced,

$$\mathcal{R}(z) := R(z)P^{-2}(z), \qquad \Xi := I \cup I^* \cup \mathbb{T}^\delta,$$

(3.7)
$$E(z) := \begin{cases} \begin{pmatrix} 1 & 0 \\ P^{-2}(z)\chi(z)\mathrm{e}^{2tg(z)} & 1 \end{pmatrix}, & z \in I, \\ \sigma_1 E(z^{-1})\sigma_1, & z \in I^*, \\ [d(z)]^{\sigma_3} B(z)[d(z)]^{-\sigma_3}, & z \in \mathbb{T}^\delta. \end{cases}$$

The matrix $B(z)$ was defined in (2.12). Note that in the non-resonant case for $z = q_1$ and $z = q_2$

$$\sup_{z \in I \cup I^*} \|E(z) - \mathbb{I}\| = \|E(z) - \mathbb{I}\|_{L^\infty(I \cup I^*)} \le C\mathrm{e}^{-t|g(q_1)|}.$$

To obtain an analogous estimate on $\mathbb{T}^\delta$, we now have to adjust the value for δ. Denote

$$\inf_k |\Phi(z_k)| = J > 0, \quad k = 1, \dots, N.$$

By continuity of the difference $g(z) - \Phi(z)$, choose $\delta > 0$ so small that

$$\sup_k \sup_{z \in \mathbb{T}_k} |g(z) - \Phi(z)| < \frac{J}{4}.$$

Then

$$\sup_{z \in \mathbb{T}^\delta} \|E(z) - \mathbb{I}\| = \|E(z) - \mathbb{I}\|_{L^\infty(\mathbb{T}^\delta)} \le C\mathrm{e}^{-\frac{tJ}{2}}.$$

Step 2: On $\mathbb{T} \setminus \Sigma$, the jump matrix can be factorized using the standard upper-lower factorization ([5, 14]). Let $\mathcal{C}$ be a contour close to the complementary arc $\mathbb{T} \setminus \Sigma$ with endpoints z_0 and $\overline{z}_0$ and clockwise orientation. Let $\mathcal{C}^*$ be its image under the map $z \mapsto z^{-1}$, oriented counterclockwise.
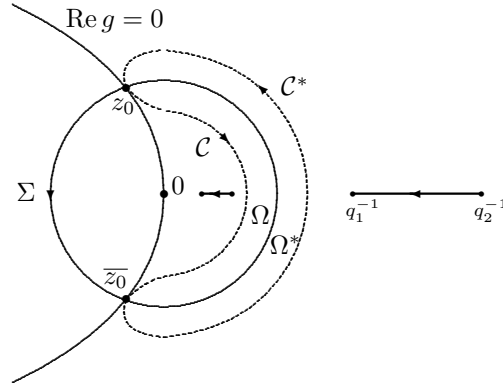

FIGURE 3. Contour deformation of Step 2.

Denote the regions adjacent to these contours by Ω and Ω^* as in Fig. 3, and set

$$W(z) := \begin{cases} \begin{pmatrix} 1 & 0 \\ -\mathcal{R}(z)\mathrm{e}^{2tg(z)} & 1 \end{pmatrix}, & z \in \Omega, \\ \sigma_1 W(z^{-1})\sigma_1, & z \in \Omega^*, \\ \mathbb{I}, & z \in \mathbb{C} \setminus (\overline{\Omega \cup \Omega^*}). \end{cases}$$

Redefine $m^{(1)}$ inside Ω and Ω^* by $m^{(2)}(z) = m^{(1)}(z)W(z)$. The new vector does not have a jump along the arc $\mathbb{T} \setminus \Sigma$, and satisfies $m^{(2)}_+(z) = m^{(2)}_-(z)v^{(2)}(z)$ with

$$v^{(2)}(z) = \begin{cases} v^{(1)}(z), & z \in \Sigma \cup \Xi, \\ W^{-1}(z), & z \in \mathcal{C} \cup \mathcal{C}^*. \end{cases}$$

The symmetry and normalization conditions are preserved in this deformation step.

Before we perform the next conjugation step, we have to study in detail the solution of the following scalar conjugation problem: *find a holomorphic function $\tilde{d}(z)$ on $\mathbb{C} \setminus \Sigma$, such that*

$$\tilde{d}_+(z)\tilde{d}_-(z) = \mathcal{R}(z)\mathcal{R}^{-1}(-1), \quad z \in \Sigma, \tag{3.8}$$

$$\text{(i) } \tilde{d}(z^{-1}) = \tilde{d}^{-1}(z), \quad z \in \mathbb{C} \setminus \Sigma; \quad \text{(ii) } \tilde{d}(\infty) \in \mathbb{R}. \tag{3.9}$$

Remark 3.2. As for any multiplicative scalar jump problem with nonvanishing jump function on a contour in $\mathbb{C}$, one can find its solution via the associated additive jump problem and the usual Cauchy integral. However, the representation via the usual Cauchy integral requires a prescribed (and known) behavior of the solutions at ∞. On the other hand, this representation cannot provide (3.9), unless the jump functions are not even on Σ. For example, condition (3.9), (i), implies $\log \tilde{d}(1) = 0$, which cannot be obtained by the usual Cauchy integral. That is why we use the Cauchy integral with kernel vanishing $z = 1$,

$$\Omega(z, s) = \frac{1}{2}\left(\frac{s+z}{s-z} - \frac{s+1}{s-1}\right). \tag{3.10}$$

In order to solve the conjugation problem (3.8)–(3.9), we first solve an auxiliary conjugation problem: find a holomorphic function $F(z)$ in $\mathbb{C} \setminus \Sigma$, bounded as $z \to \infty$ and such that

$$F_+(z) = -F_-(z), \ z \in \Sigma; \quad F(z^{-1}) = -F(z), \ z \in \mathbb{C} \setminus \Sigma; \quad F(\infty) \in \mathbb{R}. \tag{3.11}$$

Since there are two associated additive jump problems, $\log F_+(z) = \log F_-(z) \pm i\pi$, $z \in \Sigma$, and since

$$\frac{i\pi}{2\pi i}\int_{z_0}^{\overline{z_0}} \Omega(z, s)\frac{ds}{s} = \log\left(\frac{z_0 z - 1}{z_0 - z}\right)^{1/2},$$

the solution of the jump problem in (3.11) can be given by

$$F(z) = \frac{1}{2i}\left(\left(\frac{z_0 z - 1}{z_0 - z}\right)^{1/2} - \left(\frac{z_0 - z}{z_0 z - 1}\right)^{1/2}\right).$$

The symmetry $F(z^{-1}) = -F(z)$ is evident from here. It turns out that this symmetry implies that $F_+(s)$ is an even function for $s \in \Sigma$. Moreover, $F_+(s)$ is real-valued as $s = e^{i\theta}$ and $\theta \neq \theta_0$,

$$\tag{3.12}$$

$$F_+(e^{i\theta}) = \frac{1}{2i}\left(\left(\frac{\sin\frac{\theta_0+\theta}{2}}{\sin\frac{\theta_0-\theta}{2}}\right)^{1/2} - \left(\frac{\sin\frac{\theta_0-\theta}{2}}{\sin\frac{\theta_0+\theta}{2}}\right)^{1/2}\right) = \sqrt{2}\frac{\sin\frac{\theta}{2}\cos\frac{\theta_0}{2}}{\sqrt{\cos\theta_0 - \cos\theta}} \in \mathbb{R},$$

because $\cos\theta < \cos\theta_0$ on the arc Σ. On the other hand,

$$F(z) = \frac{1}{2i}\frac{(z_0+1)(z-1)}{\sqrt{(z_0-z)(z_0 z - 1)}} = \cos\frac{\theta_0}{2}\frac{(z-1)}{\sqrt{(z-z_0)(z-\overline{z}_0)}}, \tag{3.13}$$

and in particular, $F(1) = 0$ and $F(\infty) = -F(0) = \cos\frac{\theta_0}{2} \in \mathbb{R}$.

Now assume that $\tilde{d}(z)$ solves (3.8)–(3.9) and introduce the function $f(z) := F(z)\log\tilde{d}(z)$. Then (3.8), (3.9), and (3.11) imply that this function should solve the jump problem

$$(3.14) \qquad f_+(z) = f_-(z) + h(z), \quad z \in \Sigma, \quad \text{where } h(z) := F_+(z)\log\frac{\mathcal{R}(z)}{\mathcal{R}(-1)}.$$

Moreover, $f(z^{-1}) = f(z)$ should hold for $z \in \mathbb{C}\setminus\Sigma$ and $f(\infty) \in \mathbb{R}$. The jump (3.14) can be satisfied by

$$(3.15) \qquad f(z) = \frac{1}{4\pi i}\int_{z_0}^{\overline{z_0}} h(s)\left(\frac{s+z}{s-z} - \frac{s+1}{s-1}\right)\frac{ds}{s}.$$

One has to check that the other two conditions are satisfied too. The function $\log\frac{\mathcal{R}(z)}{\mathcal{R}(-1)}$ is an odd function. Indeed, recall that the reflection coefficient $R(z)$ is a continuous function which satisfies $R(z^{-1}) = R^{-1}(z)$, $|R(z)| = 1$, and so does $P^2(z)$, which implies $\mathcal{R}(z^{-1}) = \mathcal{R}^{-1}(z)$. Since $F_+(z)$ is an even function, $h(z)$ is an odd function on Σ, and the required evenness of f can be directly verified from (3.15).

To check that $f(\infty) \in \mathbb{R}$ recall that $|\mathcal{R}(z)| = 1$ yields $h(z) \in i\mathbb{R}$. When $z \to \infty$ the first integrand in (3.15) vanishes due to oddness of h. For the second we obtain

$$f(\infty) = -\frac{1}{4\pi i}\int_{z_0}^{\overline{z_0}} h(s)\frac{s+1}{s-1}\frac{ds}{s} = -\frac{1}{4\pi i}\int_{\theta_0}^{2\pi-\theta_0} h(\theta)\cot(\tfrac{\theta}{2})d\theta \in \mathbb{R}.$$

Since $1 \notin \Sigma$, the function $f(z)$ is holomorphic in a vicinity of $z = 1$ and has a zero at least of first order there. Hence $F^{-1}(z)f(z)$ is well defined, it satisfies the required additive jump problem for $\log\tilde{d}(z)$, the symmetry $F^{-1}(z)f(z) = -F^{-1}(z^{-1})f(z^{-1})$, and $F^{-1}(\infty)f(\infty) \in \mathbb{R}$. Thus, $\tilde{d}(z) = e^{F^{-1}(z)f(z)}$ solves the conjugation problem (3.8)–(3.9) and combining (3.10), (3.12)–(3.15) we obtain

$$(3.16) \quad \tilde{d}(z) = \exp\left(\frac{i\sqrt{(z-z_0)(z-\overline{z_0})}}{\sqrt{2}\pi(z-1)}\int_{\theta_0}^{2\pi-\theta_0}\arg\frac{\mathcal{R}(e^{i\theta})}{\mathcal{R}(-1)}\frac{\Omega(z,e^{i\theta})\sin\frac{\theta}{2}}{\sqrt{\cos\theta_0-\cos\theta}}d\theta\right).$$

Moreover, taking into account that

$$\frac{\Omega(z,s)}{z-1} = \frac{2s}{(s-z)(s-1)},$$

and using the first equality in (3.13) we obtain another representation for $\tilde{d}(z)$,

$$(3.17) \qquad \tilde{d}(z) = \exp\left(\frac{\sqrt{(z_0-z)(z_0 z-1)}}{2\pi i}\int_{z_0}^{\overline{z_0}}\frac{\log\frac{\mathcal{R}(s)}{\mathcal{R}(-1)}\,ds}{[\sqrt{(z_0-s)(z_0 s-1)}]_+(s-z)}\right).$$

Lemma 3.3. *The solution of the conjugation problem* (3.8)–(3.9) *given by* (3.16) *has the following asymptotic behavior as* $z \to \infty$

$$(3.18) \qquad \tilde{d}(z) = e^{A(\xi)+\frac{B(\xi)}{z}+O(z^{-2})},$$

where

(3.19)
$$A(\xi) = \frac{1}{\sqrt{2\pi}} \int_{\theta_0}^{2\pi-\theta_0} \sqrt{\cos\theta_0 - \cos\theta} \frac{2u'(\theta)\sin\frac{\theta}{2} - u(\theta)\cos\frac{\theta}{2}}{4\sin^2\frac{\theta}{2}} d\theta,$$

(3.20)
$$B(\xi) = 2\xi A(\xi) - \frac{1}{\sqrt{2\pi}} \int_{\theta_0}^{2\pi-\theta_0} \sqrt{\cos\theta_0 - \cos\theta} \left(2u'(\theta)\sin\frac{\theta}{2} + u(\theta)\cos\frac{\theta}{2}\right) d\theta,$$

and $u(\theta) := \arg\left(\frac{\mathcal{R}(e^{i\theta})}{\mathcal{R}(-1)}\right).$

Remark 3.4. The function $\mathcal{R}$, which is in fact a function of the spectral parameter λ, depends on θ via $\mathcal{R}(\cos\theta)$, therefore the derivative of $u(\theta)$ is a real valued function.

Proof. The function $u(\theta)$ is an odd function of θ on the interval of integration in the sense that $u(\frac{\pi}{2} - \alpha) = -u(\frac{\pi}{2} + \alpha)$. In the same sense, $\sin\frac{\theta}{2}(\cos\theta_0 - \cos\theta)^{-1/2}$ is an even function. Taking this into account as well as (3.10) we have for $z \to \infty$

$$i\frac{u(\theta)\Omega(z, e^{i\theta})\sin\frac{\theta}{2}}{\sqrt{\cos\theta_0 - \cos\theta}} = -i\left(\frac{e^{i\theta}}{e^{i\theta} - 1} + \frac{e^{i\theta}}{z}\right)\frac{u(\theta)\sin\frac{\theta}{2}}{\sqrt{\cos\theta_0 - \cos\theta}} + O(z^{-2})$$

$$= \left(-\frac{u(\theta)}{2\sin\frac{\theta}{2}} + \frac{2u(\theta)\sin\frac{\theta}{2}}{z}\right)\frac{d}{d\theta}\sqrt{\cos\theta_0 - \cos\theta} - i\frac{u(\theta)\sin\frac{\theta}{2}}{2\sqrt{\cos\theta_0 - \cos\theta}} + \frac{\tilde{u}(\theta)}{z} + O(z^{-2}),$$

where

$$\tilde{u}(\theta) = \frac{-iu(\theta)\cos\theta\sin\frac{\theta}{2}}{\sqrt{\cos\theta_0 - \cos\theta}}$$

is an odd function, $\tilde{u}(\frac{\pi}{2} - \alpha) = -\tilde{u}(\frac{\pi}{2} + \alpha)$. Integration by parts then yields

(3.21)
$$i\int_{\theta_0}^{2\pi-\theta_0} \frac{u(\theta)\Omega(z, e^{i\theta})\sin\frac{\theta}{2}}{\sqrt{\cos\theta_0 - \cos\theta}} d\theta = \int_{\theta_0}^{2\pi-\theta_0} \sqrt{\cos\theta_0 - \cos\theta}\frac{d}{d\theta}\left(\frac{u(\theta)}{2\sin\frac{\theta}{2}} - \frac{2u(\theta)\sin\frac{\theta}{2}}{z}\right) d\theta$$

plus the term of order z^{-2}. On the other hand,

$$\frac{\sqrt{(z - z_0)(z - \overline{z_0})}}{\sqrt{2}\pi(z - 1)} = \frac{1}{\sqrt{2}\pi}\left(1 + \frac{1 - \cos\theta_0}{z}\right) + O(z^{-2}), \quad z \to \infty.$$

Combining this with (3.21) yields (3.19) and (3.20). $\qquad\square$

Lemma 3.5. *The function* $\tilde{d}(z)$ *satisfying* (3.8)–(3.9) *has the following asymptotic behavior in a vicinity of* z_0,

(3.22)
$$\tilde{d}^{-2}(z)\mathcal{R}(z) = \mathcal{R}(-1) + O(\sqrt{z - z_0}), \ z \notin \Sigma, \quad \frac{\tilde{d}_+(z)}{\tilde{d}_-(z)} = 1 + O(\sqrt{z - z_0}), \ z \in \Sigma.$$

Proof. We will use the representation (3.17). To simplify notation set

$$r(s) = \log(\mathcal{R}(s)\mathcal{R}^{-1}(-1)), \quad q(s, z_0) = \sqrt{(z_0 - z)(z_0 z - 1)}.$$

Then

$$\log\tilde{d}(z) = \frac{q(z, z_0)}{2\pi i}\int_{z_0}^{\overline{z_0}} \frac{r(s)ds}{[q(s, z_0)]_+(s - z)} = J_1(z) + J_2(z),$$

where

$$J_1(z) = \frac{q(z, z_0)}{2\pi i} \int_{z_0}^{\bar{z}_0} \frac{(r(s) - r(z_0))ds}{[q(s, z_0)]_+ (s - z)}, \quad J_2(z) = r(z_0) \frac{q(z, z_0)}{2\pi i} \int_{z_0}^{\bar{z}_0} \frac{ds}{[q(s, z_0)]_+ (s - z)}.$$

Since $r(s) - r(z_0) \sim (z - z_0)$, the integral in $J_1(z)$ is Hölder continuous in a vicinity of z_0. Therefore,

(3.23)

$$J_1(z) = I(z_0)\sqrt{z - z_0}, \quad I(z_0) = \frac{1 - z_0^2}{2\pi i} \int_{z_0}^{\bar{z}_0} \frac{(\log \mathcal{R}(s) - \log \mathcal{R}(z_0))\, ds}{[\sqrt{(z_0 - s)(z_0 s - 1)}]_+ (s - z_0)}.$$

On the other hand, since

$$\frac{1}{2[q(z, z_0)]_+} = \frac{1}{2[q(z, z_0)]_-} + \frac{1}{[q(z, z_0)]_+}, \quad z \in \Sigma,$$

and $(q(z, z_0))^{-1} \to 0$ as $z \to \infty$, we have

$$\frac{1}{2q(z, z_0)} = \frac{1}{2\pi i} \int_{z_0}^{\bar{z}_0} \frac{ds}{[q(s, z_0)]_+ (s - z)},$$

and $J_2 = \frac{r(z_0)}{2}$. Therefore,

$$\log \tilde{d}(z) = \frac{1}{2} \log \frac{\mathcal{R}(z_0)}{\mathcal{R}(-1)} + I(z_0)\sqrt{z - z_0},$$

and (3.22) follows from (3.23) in a straightforward manner. $\qquad \square$

Step 3: Define $m^{(3)}(z) = m^{(2)}(z)[\tilde{d}(z)]^{-\sigma_3}$, then our previous considerations lead to the following statement.

Theorem 3.6. *The vector function $m^{(3)}(z)$ is the unique solution of the following RH problem: find a holomorphic vector function $\breve{m}(z)$ in $\mathbb{C} \setminus (\Sigma \cup \mathcal{C} \cup \mathcal{C}^* \cup \Xi)$, which is continuous up to the boundary and has the following properties:*

- *It solves the jump problem $\breve{m}_+(z) = \breve{m}_-(z)\breve{v}(z)$ with*

(3.24)
$$\breve{v}(z) = \begin{cases} \begin{pmatrix} 0 & -\mathcal{R}(-1) \\ \mathcal{R}(-1) & \frac{\tilde{d}_+(z)}{\tilde{d}_-(z)} e^{-2tg_+(z)} \end{pmatrix}, & z \in \Sigma, \\[12pt] \begin{pmatrix} 1 & 0 \\ -\tilde{d}^{-2}(z)\mathcal{R}(z) e^{2tg(z)} & 1 \end{pmatrix}, & z \in \mathcal{C}, \\[12pt] \begin{pmatrix} 1 & -\tilde{d}^2(z)\overline{\mathcal{R}(z)} e^{-2tg(z)} \\ 0 & 1 \end{pmatrix}, & z \in \mathcal{C}^*, \\[6pt] \breve{E}(z) := [\tilde{d}(z)]^{\sigma_3} E(z)[\tilde{d}(z)]^{-\sigma_3}, & z \in \Xi, \end{cases}$$

  where the matrix $\breve{E}(z)$ is defined by (3.7), (2.12), (3.16), and satisfies

$$\|\breve{E}(z) - \mathbb{I}\|_{L^\infty(\Xi)} \leq C e^{-t\breve{J}}, \quad \breve{J} = \min \left\{ |g(q_1)|, \min_{k=1,\ldots,N} |4^{-1}\Phi(z_k)| \right\};$$

- *$\breve{m}(z)$ satisfies symmetry and normalization conditions like (2.8) and (2.13); moreover, $\breve{v}(z)$ has the symmetry property (2.7);*
- *For small z the vector function m in (2.5) and $\breve{m}(z)$ are connected by*

(3.25)
$$\breve{m}(z) = m(z)[q(z)]^{-\sigma_3}, \quad q(z) = \tilde{d}(z)P(z)e^{t(\Phi(z) - g(z))}.$$

4. The model problem solution

Denote by $\check{\Sigma} = \Sigma \cup \mathcal{C} \cup \mathcal{C}^* \cup \Xi$ the jump contour for $\check{m}$, and let $\mathcal{O}$ and $\mathcal{O}^*$ be small and symmetric (with respect to the map $z \mapsto z^{-1}$) vicinities of z_0 and $\overline{z}_0$, which are not necessarily circles. Denote by $\Sigma_1 = \Sigma \cap \mathcal{O}$, $\Sigma_2 = \mathcal{C} \cap \mathcal{O}$, $\Sigma_3 = (\mathcal{C}^*)^{-1} \cap \mathcal{O}$ the segments of $\check{\Sigma}$ in $\mathcal{O}$ and set $\Sigma_{\mathcal{O}} = \Sigma_1 \cup \Sigma_2 \cup \Sigma_3$. Let $\Sigma_{\mathcal{O}}^*$ be the symmetric contour in a neighborhood of $\overline{z}_0$. Both contours $\Sigma_{\mathcal{O}}$ and $\Sigma_{\mathcal{O}}^*$ inherit the orientation of the respective parts of $\check{\Sigma}$, except for $\mathcal{C}^*$, where we reverse the orientation. Evidently,

$$\|\check{v}(z) - v^{\mathrm{mod}}(z)\|_{L^\infty\left(\check{\Sigma} \setminus (\Sigma_{\mathcal{O}} \cup \Sigma_{\mathcal{O}}^*)\right)} \leq C e^{-tU},$$

where

$$U = \min\{\check{J}, \inf_{z \in \partial\mathcal{O}} |g(z)|\} > 0,$$

and

$$v^{\mathrm{mod}}(z) = \begin{cases} \begin{pmatrix} 0 & -\mathcal{R}(-1) \\ \mathcal{R}(-1) & 0 \end{pmatrix}, & z \in \Sigma \\ \mathbb{I}, & z \in \check{\Sigma} \setminus \Sigma. \end{cases}$$

Thus, in a first order of approximation one can assume that the solution of the RHP (3.24) can be approximated by the solution of the following model RHP: *find a holomorphic vector function in $\mathbb{C} \setminus \Sigma$ satisfying the jump condition*

$$(4.1) \qquad\qquad m_+^{\mathrm{mod}}(z) = m_-^{\mathrm{mod}}(z) v^{\mathrm{mod}}(z), \quad z \in \Sigma,$$

the symmetry condition $m^{\mathrm{mod}}(z^{-1}) = m^{\mathrm{mod}}(z)\sigma_1$, and the normalization condition $m_1^{\mathrm{mod}}(0) > 0$, $m_1^{\mathrm{mod}}(0) m_2^{\mathrm{mod}}(0) = 1$.

Lemma 4.1. *The solution of this vector RH problem is unique.*

The proof of this Lemma is analogous to the proof of the uniqueness of the vector model problem for the KdV rarefaction problem.

For our further investigation we will also need a matrix solution $M^{\mathrm{mod}}(z) = M^{\mathrm{mod}}(z, n, t)$ of the matrix RHP: *find a holomorphic matrix function $M^{\mathrm{mod}}(z)$ on $\mathbb{C} \setminus \Sigma$ satisfying the following jump and symmetry conditions,*

$$M_+^{\mathrm{mod}}(z) = M_-^{\mathrm{mod}}(z) v^{\mathrm{mod}}(z), \quad z \in \Sigma; \qquad M^{\mathrm{mod}}(z^{-1}) = \sigma_1 M^{\mathrm{mod}}(z)\sigma_1.$$

We find a solution of the matrix problem following [12, 1]. Consider the non-resonant case, that is $\mathcal{R}(-1) = -1$. Using

$$v^{\mathrm{mod}}(z) = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ i & -i \end{pmatrix} i\sigma_3 \begin{pmatrix} 1 & 1 \\ i & -i \end{pmatrix}^{-1},$$

we first look for a holomorphic solution of the jump problem $M_+^0 = i M_-^0 \sigma_3$ satisfying the symmetry condition as above. Using (3.10) we get

$$M^0(z) = \begin{pmatrix} \beta(z) & 0 \\ 0 & \beta^{-1}(z) \end{pmatrix},$$

with

$$\beta(z) = \left(\frac{z_0 z - 1}{z_0 - z}\right)^{1/4},$$

where the branch of the fourth root is chosen with the cut along the negative half axis and $1^{1/4} = 1$. Since $\beta(z^{-1}) = \beta^{-1}(z)$, then we have the required symmetry. Concerning the original matrix solution $M^{\mathrm{mod}}(z)$, this asks for the representation

$$(4.2) \qquad M^{\mathrm{mod}}(z) = \begin{pmatrix} \frac{\beta(z)+\beta^{-1}(z)}{2} & \frac{\beta(z)-\beta^{-1}(z)}{2\mathrm{i}} \\ -\frac{\beta(z)-\beta^{-1}(z)}{2\mathrm{i}} & \frac{\beta(z)+\beta^{-1}(z)}{2} \end{pmatrix},$$

and the required symmetry condition is also fulfilled. The vector solution of the model problem is unique. Evidently, if we take $m^{\mathrm{mod}}(z) = (\alpha, \alpha)M^{\mathrm{mod}}(z)$ for some α, then (4.1) and (2.8) are fulfilled. We have to choose a suitable α to satisfy the normalization condition. Since $\beta(\infty) = \mathrm{e}^{\mathrm{i}\frac{\theta_0-\pi}{4}}$ then

$$M^{\mathrm{mod}}(\infty) = \begin{pmatrix} \cos\frac{\theta_0-\pi}{4} & \sin\frac{\theta_0-\pi}{4} \\ -\sin\frac{\theta_0-\pi}{4} & \cos\frac{\theta_0-\pi}{4} \end{pmatrix}.$$

The normalization condition for $m^{\mathrm{mod}}(\infty)$ implies

$$\alpha^2 \left(\cos^2\tfrac{\theta_0-\pi}{4} - \sin^2\tfrac{\theta_0-\pi}{4}\right)\alpha^2\cos\tfrac{\theta_0-\pi}{2} = \alpha^2\sin\tfrac{\theta_0}{2} = 1.$$

Thus,

$$(4.3) \qquad m^{\mathrm{mod}}(z) = \alpha \begin{pmatrix} 1, & 1 \end{pmatrix} M^{\mathrm{mod}}(z), \quad \alpha = \alpha(\xi) = \left(\sin\tfrac{\theta_0}{2}\right)^{-1/2}.$$

In the resonant case the matrix solution is represented by the same formula (4.2), but with $\beta^{-1}(z)$ instead of $\beta(z)$ and vice versa. In summary, we have the following

Lemma 4.2. *The solution of the vector resp. matrix model RHP, $m^{\mathrm{mod}}(z)$ resp. $M^{\mathrm{mod}}(z)$, is given by (4.3) resp. (4.2), where $\beta(z) = \left(\frac{z_0 z-1}{z_0-z}\right)^{1/4}$ in the non-resonant case and $\beta(z) = \left(\frac{z_0-z}{z_0 z-1}\right)^{1/4}$ in the resonant case.*

Lemma 4.3. *The asymptotic behavior of $m_1^{\mathrm{mod}}(z)$ as $z \to 0$ is given by*

$$(4.4) \qquad m_1^{\mathrm{mod}}(z) = \frac{\xi^{1/4}}{(1+(1-\xi)^{1/2})^{1/2}}\left(1 + \left(1 - \xi + (1-\xi)^{1/2}\right)z\right) + O(z^2)$$

in the non-resonant case. In the resonant case,

$$m_1^{\mathrm{mod}}(z) = \frac{(1+(1-\xi)^{1/2})^{1/2}}{\xi^{1/4}}\left(1 + \left(1 - \xi - (1-\xi)^{1/2}\right)z\right) + O(z^2).$$

Proof. Consider first the non-resonant case. Since

$$\beta(z) = \mathrm{e}^{\mathrm{i}\frac{\pi-\theta_0}{4}}\left(1 - \frac{\mathrm{i}\sin\theta_0}{2}z\right) + O(z^2), \quad z \to 0,$$

then

$$\frac{\beta(z)+\beta^{-1}(z)}{2} = \cos\tfrac{\pi-\theta_0}{4} + \sin\tfrac{\pi-\theta_0}{4}\frac{\sin\theta_0}{2}z + O(z^2),$$

$$\frac{\beta(z)-\beta^{-1}(z)}{2\mathrm{i}} = \sin\tfrac{\pi-\theta_0}{4} - \cos\tfrac{\pi-\theta_0}{4}\frac{\sin\theta_0}{2}z + O(z^2).$$

By (4.3),

$$m_1^{\mathrm{mod}}(z) = \frac{1}{\sqrt{\sin\frac{\theta_0}{2}}}\left(\cos\tfrac{\pi-\theta_0}{4} - \sin\tfrac{\pi-\theta_0}{4} + \left(\cos\tfrac{\pi-\theta_0}{4} + \sin\tfrac{\pi-\theta_0}{4}\right)\frac{\sin\theta_0}{2}z\right) + O(z^2)$$

$$= \frac{\sqrt{2}\sin\frac{\theta_0}{4}}{\sqrt{\sin\frac{\theta_0}{2}}} + \frac{\sqrt{2}\cos\frac{\theta_0}{4}}{\sqrt{\sin\frac{\theta_0}{2}}}\frac{\sin\theta_0}{2}z + O(z^2) = \sqrt{\tan\tfrac{\theta_0}{4}} + \frac{\sin\theta_0\, z}{2\sqrt{\tan\frac{\theta_0}{4}}} + O(z^2).$$

Next,

$$\tan\frac{\theta_0}{4} = \frac{1-\cos\frac{\theta_0}{2}}{\sin\frac{\theta_0}{2}} = \frac{\sqrt{2}-\sqrt{1+\cos\theta_0}}{\sqrt{1-\cos\theta_0}} = \frac{1-\sqrt{1-\xi}}{\sqrt{\xi}} = \frac{\sqrt{\xi}}{1+\sqrt{1-\xi}},$$

since $\cos\theta_0 = 1-2\xi$. Respectively, $\sin\theta_0 = 2\sqrt{\xi-\xi^2}$, from which (4.4) follows. The resonant case follows from the same computation using $\beta^{-1}(z)$ instead of $\beta(z)$. $\square$

5. ASYMPTOTICS IN THE REGION $\xi \in (0, 1)$.

The structure of the matrix solution (4.2) and the jump matrix (3.24) as well as the results of Lemmas 3.5, 3.1, (e), allow us to conclude that the solution of the parametrix problem (which has a local character) can be constructed as in [1] using [14], Appendix B (in particular the solution can be given in terms of Airy functions) Consequently, as $z \to 0$ (cf. [1])

$$m^{(3)}(z, \xi) = \breve{m}(z) = m^{\mathrm{mod}}(z, \xi) + \frac{f(z, \xi, t)}{t}, \quad \xi \in \mathcal{I} := [\varepsilon, 1 - \varepsilon],$$

where the vector function $f(z, \xi, t)$ is uniformly bounded with respect to $z \in \mathbb{D}_\delta$, $\xi \in \mathcal{I}$ and $t \in [T, \infty)$. Here $T > 0$ is a sufficiently large number and $\mathbb{D}_\delta$ is a circle centered at 0 with radius $\delta < 1$. Moreover, $f(z, \xi, t)$ is differentiable with respect to ξ and z, and the derivatives are also uniformly bounded in the sets mentioned above. By (3.25) we will have for $f = (f_1, f_2)$

$$(5.1) \qquad m_1(z) = \left(m_1^{\mathrm{mod}}(z) + \frac{f_1(z, \xi, t)}{t}\right)\tilde{d}(z)P(z)e^{t(\Phi(z)-g(z))}.$$

From (3.9), (3.18), (3.19), (3.20), (3.4), (3.5), and (2.11) it follows that

$$\tilde{d}(z) = e^{-A(\xi)}\left(1 - B(\xi)z\right) + O(z^2), \quad z \to 0,$$

$$e^{t(\Phi(z)-g(z))} = e^{tK(\xi)}\left(1 + tk(\xi)z\right) + O(z^2), \quad z \to 0,$$

$$P(z) = \left(1 - 2z\sum_{z_j \in (-1,0)}\sqrt{\lambda_j^2 - 1}\right)\prod_{z_j \in (-1,0)}\frac{1}{|z_j|} + o(z^2).$$

Combining this with (4.4), (5.1), and $f_1(z, \xi, t) = f_1(0, \xi, t) + zf_1^{(1)}(0, \xi, t) + O(z^2)$ we get

$$(5.2) \qquad m_1(z) = \mathcal{A}(\xi)\left\{1 + 2z\left(\mathcal{B}(\xi) + \frac{h_3(\xi, t)}{t}\right) + O(z^2) + \frac{h_2(\xi, t)}{t}\right\},$$

where

$$\mathcal{A}(\xi, t) = e^{tK(\xi)-A(\xi)}\frac{\xi^{1/4}}{(1+(1-\xi)^{1/2})^{1/2}}\prod_{z_j \in (-1,0)}\frac{1}{|z_j|}$$

$$= S\exp\left(tK(\xi) - A(\xi) + \frac{1}{4}\log\xi - \frac{1}{2}\log(1 + (1-\xi)^{1/2})\right),$$

$$\mathcal{B}(\xi, t) = -S_1 + \frac{tk(\xi)}{2} - \frac{B(\xi)}{2} + \frac{1-\xi+(1-\xi)^{1/2}}{2},$$

$$S = \prod_{z_j \in (-1,0)}\frac{1}{|z_j|}, \quad S_1 = \sum_{z_j \in (-1,0)}\sqrt{\lambda_j^2 - 1}.$$

The coefficients $h_2(\xi, t)$ and $h_3(\xi, t)$ given by

$$h_2(\xi, t) = f_1(0, \xi, t)\frac{(1 + (1 - \xi)^{1/2})^{1/2}}{\xi^{1/4}}, \quad h_3(\xi, t) = f_1^{(1)}(0, \xi, t)\frac{(1 + (1 - \xi)^{1/2})^{1/2}}{\xi^{1/4}},$$

and are bounded with respect to $\xi \in \mathcal{I}$ and $t \geq T$, together with their derivatives with respect to ξ. Comparing (5.2) with (2.6) implies the asymptotics

$$\prod_{j=n}^{\infty} 2a(j, t) = \mathcal{A}(\xi)\left(1 + \frac{h_2(\xi, t)}{t}\right), \quad \sum_{m=n}^{\infty} b(m, t) = \mathcal{B}(\xi) + \frac{h_3(\xi, t)}{t}, \quad \xi = \frac{n}{t}.$$

Set $\hat{\xi} = \frac{n+1}{t}$ and note that $\xi - \hat{\xi} = -\frac{1}{t}$. Then

$$2a(n, t) = \frac{\mathcal{A}(\xi)}{\mathcal{A}(\hat{\xi})} = e^{-K'(\xi) + \frac{1}{t}L(\xi) + O(t^{-2})},$$

where

$$L(\xi) = \frac{-K''(\xi)}{2} + A'(\xi) - \frac{1}{4\xi} - \frac{1}{4(1 - \xi + (1 - \xi)^{1/2})},$$

and the prime denotes the derivative with respect to ξ. Analogously,

$$b(n, t) = \mathcal{B}(\xi) - \mathcal{B}(\hat{\xi}) = -\frac{k'(\xi)}{2} + \frac{1}{t}\left(-\frac{k''(\xi)}{4} + \frac{B'(\xi)}{2} + \frac{1}{2} + \frac{1}{4(1 - \xi)^{1/2}}\right) + O(t^{-2}).$$

By (3.19) and (3.20)

$$A'(\xi) = -\frac{1}{4\sqrt{2}\,\pi}\int_{\theta_0}^{2\pi - \theta_0}\frac{2u'(\theta)\sin\frac{\theta}{2} - u(\theta)\cos\frac{\theta}{2}}{\sqrt{\cos\theta_0 - \cos\theta}\sin^2\frac{\theta}{2}}\,d\theta,$$

$$(5.3) \qquad B'(\xi) = 2\xi A'(\xi) + 2A(\xi) + \frac{1}{\sqrt{2}\,\pi}\int_{\theta_0}^{2\pi - \theta_0}\frac{2u'(\theta)\sin\frac{\theta}{2} + u(\theta)\cos\frac{\theta}{2}}{\sqrt{\cos\theta_0 - \cos\theta}}\,d\theta,$$

with (recall that we consider the non-resonant case at $z = -1$)

$$(5.4) \qquad\qquad u(\theta) = \arg\left(\mathcal{R}(e^{i\theta}) + 1\right).$$

Taking into account (3.5) and (3.6) we get

$$a(n, t) = \frac{\xi}{2}\left(1 + \frac{L(\xi)}{t} + O(t^{-2})\right),$$

where

$$L(\xi) = -\frac{1 + 2\xi + (1 - \xi)^{1/2}}{4\xi(1 - \xi + (1 - \xi)^{1/2})} + A'(\xi).$$

Analogously, by (3.5)

$$b(n, t) = 1 - \xi + \frac{1}{4t}\left((1 - \xi)^{-1/2} + 2B'(\xi)\right) + O(t^{-2}).$$

In summary, we proved the following

Theorem 5.1. *Suppose that the initial operator $H(0)$ associated with the sequences $\{a(n, 0), b(n, 0)\}$ has no resonance at the edges of spectrum $-1, b - 2a, b + 2a$. Let $R(z)$, $z = \lambda - \sqrt{\lambda^2 - 1}$, be its right reflection coefficient and $\lambda_j = \frac{1}{2}(z_j + z_j^{-1})$ its eigenvalues. Let $\varepsilon > 0$ be an arbitrary small number. Then in the sector*

$$\varepsilon t \leq n \leq (1 - \varepsilon t)$$

the following asymptotics are valid for the solution $\{a(n,t), b(n,t)\}$ of the Toda lattice as $t \to +\infty$:

$$(5.5) \qquad a(n,t) = \frac{n}{2t} - \frac{n}{8t^2} \left(\frac{\sqrt{1 - \frac{n}{t}} + 1 + \frac{2n}{t}}{\frac{n}{t}\left(1 - \frac{n}{t} + \sqrt{1 - \frac{n}{t}}\right)} \right.$$

$$\left. + \frac{1}{\sqrt{2}\pi} \int_{\theta_0}^{2\pi - \theta_0} \frac{2u'(\theta)\sin\frac{\theta}{2} - u(\theta)\cos\frac{\theta}{2}}{\sqrt{1 - \frac{2n}{t} - \cos\theta \, \sin^2\frac{\theta}{2}}} d\theta \right) + O(t^{-2}),$$

$$(5.6) \qquad b(n,t) = 1 - \frac{n}{t} + \frac{1}{4t}\left(\frac{1}{\sqrt{1 - \frac{n}{t}}} + 2B'\left(\frac{n}{t}\right)\right) + O(t^{-2}),$$

where the term $O(t^{-2})$ is uniformly bounded with respect to n. The function $B'(\xi)$ is defined by (5.3), (3.20), (3.19), (5.4); $\theta_0 = \arccos(1 - \frac{2n}{t}) \in (0, \pi)$ and

$$u(\theta) = \arg\left(R(e^{i\theta})P^{-2}(e^{i\theta}) - \arg R(e^{i\pi})P^{-2}(e^{i\pi})\right)$$

where P is defined in (2.11).

6. Discussion of the region $-2at < n < 0$.

To obtain the asymptotic of the solution in the region $\frac{n}{t} \in (-2a, 0)$ we could study an analogous RH problem connected with the left scattering data and the variable ζ (cf. (2.2)). Instead of this extensive analysis let us consider the Toda lattice associated with the functions

$$(6.1) \qquad \hat{a}(n,t) = \frac{1}{2a}a(-n-1, \tfrac{t}{2a}), \quad \hat{b}(n,t) = \frac{1}{2a}\big(b - b(-n, \tfrac{t}{2a})\big).$$

It is straightforward to check that $\hat{a}(n,t), \hat{b}(n,t)$ satisfy the Toda equations (1.1) associated with the initial profile

$$\hat{a}(n,0) \to \frac{1}{2a}, \quad \hat{b}(n,0) \to \frac{b}{2a}, \quad \text{as } n \to -\infty,$$

$$\hat{a}(n,0) \to \frac{1}{2}, \quad \hat{b}(n,0) \to 0, \quad \text{as } n \to +\infty.$$

For this solution we obtain by our previous results in the region $\frac{n}{t} \in (0, 1)$ that

$$(6.2) \qquad \hat{a}(n,t) = \frac{n}{2t} + \frac{f(n,t)}{t}, \quad \hat{b}(n,t) = 1 - \frac{n}{t} + \frac{g(n,t)}{t},$$

where $f(n,t)$ (resp., $g(n,t)$) are the same as the second order terms in (5.5) (resp., (5.6)), but the function u corresponds to a new Jacobi operator (2.1) with coefficients $\hat{H} \sim \{\hat{a}(n,0), \hat{b}(n,0)\}$. Set $\tau = \frac{t}{2a}$ and $m = -n - 1$. If $\frac{n}{t} \in (0,1)$, then $\frac{m}{\tau} \in (-2a, 0)$. From (6.1) and (6.2) we get

$$a(m, \tau) = 2a\hat{a}(n,t) = \frac{2an}{2t} + \frac{2af(n,t)}{t} = \frac{-m-1}{2\tau} + \frac{f(-m-1, 2a\tau)}{\tau},$$

$$b(m, \tau) = b - 2a\hat{b}(n+1, t) = b - 2a - \frac{m}{\tau} - \frac{g(-m-1, 2a\tau)}{\tau}.$$

An elementary analysis shows that the right reflection coefficient of $\hat{H}$, given in the variable ζ, is the same as the left reflection coefficient $R_1(\zeta^{-1})$ of the initial

operator $H(0)$, and the discrete spectra of both operators are the same. In terms of ζ we denote the eigenvalues by $\zeta_1, \ldots, \zeta_N$ and set

$$P_1(\zeta) = \prod_{\zeta_j \in (0,1)} |\zeta_j| \frac{\zeta - \zeta_j^{-1}}{\zeta - \zeta_j}.$$

Note that $P(z)$ corresponds to the eigenvalues λ_j of $H(0)$, which satisfy $\lambda_j < -1$, and $P_1(\zeta)$ corresponds to $\lambda_j > b + 2a$. Now set $\zeta = \mathrm{e}^{\mathrm{i}\theta}$ and

$$U(\theta) = \arg\left(R_1(\mathrm{e}^{\mathrm{i}\theta}) P_1^{-2}(\mathrm{e}^{\mathrm{i}\theta}) - R_1(1) P_1^{-2}(1)\right), \quad \theta \in [-\theta_0, \theta_0].$$

Theorem 6.1. *Let $\varepsilon > 0$ be an arbitrary small number. Suppose that the initial operator $H(0)$ has no resonances at the points $-1, 1, b + 2a$, and let $R_1(\zeta)$ be its left reflection coefficient, $\lambda - b = a(\zeta + \zeta^{-1})$. Then in the domain $-(2a - \varepsilon)t \leq n \leq -\varepsilon t$ the following asymptotic is valid as $t \to +\infty$ in the non-resonant case,*

$$a(n,t) = -\frac{n+1}{2t} - \frac{n+1}{16at^2} \left(\frac{\sqrt{1 + \frac{n+1}{2at}} + 1 - \frac{n+1}{at}}{\frac{n+1}{2at}\left(1 + \frac{n+1}{2at} + \sqrt{1 + \frac{n+1}{2at}}\right)} \right.$$

$$\left. -\frac{1}{\sqrt{2\pi}} \int_{-\theta_0}^{\theta_0} \frac{2U'(\theta)\sin\frac{\theta}{2} - U(\theta)\cos\frac{\theta}{2}}{\sqrt{1 + \frac{n+1}{at} - \cos\theta \sin^2\frac{\theta}{2}}} d\theta \right) + O(t^{-2}),$$

$$b(n,t) = b - 2a - \frac{n}{t} - \frac{1}{2t}\left(\frac{1}{2\sqrt{1 + \frac{n+1}{2at}}} + \tilde{B}(\tfrac{n}{t}) \right) + O(t^{-2}),$$

where $\theta_0 = \arccos(1 + \frac{n+1}{at}) \in (0,\pi)$ and

$$\tilde{B}(\tfrac{n}{t}) = \frac{1}{\sqrt{2\pi}} \int_{-\theta_0}^{\theta_0} \left(\frac{n+1}{4at} \frac{2U'(\theta)\sin\frac{\theta}{2} - U(\theta)\cos\frac{\theta}{2}}{\sqrt{1 + \frac{n+1}{at} - \cos\theta \sin^2\frac{\theta}{2}}} + \frac{2U'(\theta)\sin\frac{\theta}{2} + U(\theta)\cos\frac{\theta}{2}}{\sqrt{1 + \frac{n+1}{at} - \cos\theta}} \right.$$

$$\left. + \sqrt{1 + \frac{n+1}{at} - \cos\theta}\, \frac{2U'(\theta)\sin\frac{\theta}{2} - U(\theta)\cos\frac{\theta}{2}}{2\sin^2\frac{\theta}{2}} \right) d\theta.$$

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