High-frequency imaging of a moving object

Clifford Nolan University of Limerick

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- Joint work with Felea and Gaburro.
- Motivated by papers of Cheney and Borden.
- Seek to invert RADAR data for a time-dependent reflectivity function.
- We will see that backprojected RADAR images have artifacts that can be reduced by pre-processing the data.

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Static Synthetic Aperture RADAR (SAR)

• In static RADAR we have the following set up:



- The object to be imaged (v(x) := c⁻²(x) c₀⁻²) is assumed static and the RADAR flies, makes measurements, creating a synthetic aperture
- Backprojection produces an image with standard mirror-artifacts.

- Now consider an object which moves as time elapses
- Scalar wave equation model for radio waves emitted from location y at time T_y and measured at (x, t):

$$(\Delta - \frac{1}{c^2(x,t)}\partial_t^2)u(y;x,t) = \delta(t+T_y)\delta(x-y)$$

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- Notice the time-dependent speed c(x, t).
- Corresponding to this, we build in the flexibility of the initiation time T_y of our RADAR at location y - makes it possible to see multiple facets of object, as it moves.

Linearization

- x is to be thought of as a location of a scatterer in the vicinity of the ground - we consider x ∈ ℝ² or x ∈ ℝ³.
- y is to be thought of as the location of the source we consider y ∈ ℝ² or y ∈ ℝ³.
- Linearization: $u(y; x, t) = u^{in}(y; x, t) + u^{sc}(y; x, t)$ where

$$(\Delta - \frac{1}{c_0^2}\partial_t^2)u^{in}(y; x, t) = \delta(t + T_y)\delta(x - y),$$

$$(\Delta - \frac{1}{c_0^2} \partial_t^2) u^{sc}(y; x, t) = v(x, t) \frac{\partial^2 u^{in}}{\partial^2 t}(y; x, t)$$

and

$$v(x,t) := \frac{1}{c^2(x,t)} - \frac{1}{c_0^2}$$

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• Convolving the source for u^{sc} with the Green's function gives

$$u^{sc}(y;z,t) = \int \frac{\delta(t-t'-\frac{|x-z|}{c_0})}{4\pi|x-z|} \frac{\partial_{t'}^2 \delta(t+T_y-\frac{|x-y|}{c_0})}{4\pi|x-y|} v(x,t') dx dt'$$

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- We make a simple and concrete choice T_y = αy₁, so that sources are fired at different times for different locations along the Y₁ axis.
- We also make a technical but reasonable assumption that $c_0 \alpha \neq 1$.

FIO Representation of Scattering Operator F

 Writing δ as an oscillatory integral we arrive at the forward modeling scattering operator F which maps v to u^{sc} as follows

$$Fv(y,z,t) = u^{sc}(y;z,t) = \int e^{i[(t-t'-|x-z|)\omega+(t'+c_0\alpha y_1-|x-y|)\omega']} \frac{c_0^2\omega^2}{|x-y||x-z|}v(x,t')d\omega d\omega' dx dt'$$

• The operator *F* can easily be verified to be a Fourier Integral Operator (FIO) of order *q*, say.

Summary of Results

- We consider the model $v(x, t') \in \mathcal{E}'(X)$ where $X := \mathbb{R}^m$ and the data $Fv(y, z, t) \in \mathcal{E}'(Y)$ where $Y := \mathbb{R}^d$.
- Since F is a FIO, it has a canonical relation
 C ⊂ T*Y × T*X, which describes how F maps singularities in v to singularities in u^{sc}.
- The acquisition geometry determined in part by *m*, *d* strongly influences the structure of the relation *C* and we now consider some explicit geometries ...

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• In all cases we consider the natural projections $\pi_L: C \longrightarrow T^*Y, \ \pi_R: C \longrightarrow T^*X$

- Case 1: The "deluxe" data case: y, z, x belong to bounded subsets of ℝ³
- In this case, d = 7, m = 4.
- $\Rightarrow \pi_L$ an embedding (and π_R a submersion).
- $\Rightarrow F^*F$ is a microlocally elliptic Ψ DO (by Guillemin's result)

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• Therefore, *F* has a left-parametrix and this is the nicest possible case - but expensive to collect and invert data.

Summary of Results

- Case 2: Same as case 1 but the source and receiver are coincident *y* = *z*.
- This is a formally determined case (d = m = 4).
- π_L and π_R both have blowdown singularities (more on this later).
- If we apply *F*^{*} to the data, we show that artifacts appear which **can be more singular** than the true singularities.
- This follows from K_{F*F} belonging to the class I^{5/2,-1/2}(△, Λ) (more on this later).

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- Case 3: Same as case 2 but y₃ = z₃ are constant and the scatterer is assumed to be on the ground x₃ = 0.
- This is a formally determined case (d = m = 3).
- π_L and π_R both have blowdown singularities.
- If we apply *F*^{*} to the data, we show that artifacts appear which **can be just as singular** as the true singularities.
- This follows from K_{F^*F} belonging to the class $I^{3,0}(\triangle, \Lambda)$.

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- Case 4: y₃ = z₃ are constant. y₂ = z₂, y₁ ≠ z₁, no restriction on scatterer location.
- This is a formally determined case (d = m = 4).
- π_L and π_R both have blowdown singularities.
- If we apply *F*^{*} to the data, we show that artifacts appear which **can be just as singular** as the true singularities.
- This follows from K_{F^*F} belonging to the class $I^{2,0}(\triangle, \Lambda)$.

Ingredients of Analysis

- The results for case 1 are easily understood.
- For cases 2-4 we find that π_L, π_R have a blowdown singularity along a submanifold $\Sigma \subset C$.
- For π_L for example, this means that (i) π_L drops rank by k > 0 at Σ, (ii) Ker(Dπ_L|_Σ) ⊂ TΣ and (iii) det(Dπ_L) vanishes to order k at Σ.
- Canonical form of a blowdown: $f(x_1, \dots, x_{n-k}, x_{n-k+1}, \dots, x_n) = (x_1, x_2, \dots, x_{n-k}, x_{n-k+1}x_1, \dots, x_nx_1)$
- We are able to apply a theorem of Marhuenda which states that if π_L, π_R only have blowdown singularities at Σ and π_L(Σ), π_R(Σ) are involutive and non-radial, then the distribution kernel K_{F*F} ∈ I^{2q+(k-1)/2,-(k-1)/2}(△, Λ_{π_R(Σ)}).

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• The 'artifact' manifold $\Lambda_{\pi_R(\Sigma)}$ is as follows . . .

- As stated π_R(Σ) is involutive, which roughly means its homogeneous defining functions are given by equations of the form ξ_i = 0, i = 1,...r for some r > 0 with {ξ_i, ξ_j} = 0, i, j = 1,..., r.
- The artifact submanifold Λ_{π_R(Σ)} is the joint flow-out from Σ by the Hamiltonian vector fields {H_{ξi}}^r_{i=1}.
- Note that u ∈ I^{p,I}(Δ, Λ) ⇒ u ∈ I^{p+I}(Δ \ Λ) and u ∈ I^p(Λ \ Δ) which leads to the results about the strength of the artifact singularities that we quoted.

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Ingredients of Analysis

• For example, in case 4, Σ is defined by the single equation:

$$\frac{\omega'}{|x-y|} + \frac{\omega}{|x-z|} = 0$$

$$C = \left\{ y_1, y_2, z_1, t, c_0 \alpha \omega' + \omega' \frac{x_1 - y_1}{|x - y|}, \omega' \frac{x_2 - y_2}{|x - y|} + \omega \frac{x_2 - y_2}{|x - z|}, \\ \omega \frac{x_1 - z_1}{|x - z|}, \omega; x_1, x_2, x_3, t', \omega' \frac{x_1 - y_1}{|x - y|} + \omega \frac{x_1 - z_1}{|x - z|}, \\ \omega' \frac{x_2 - y_2}{|x - y|} + \omega \frac{x_2 - y_2}{|x - z|}, \omega' \frac{x_3 - h}{|x - y|} + \omega \frac{x_3 - h}{|x - z|}, \omega - \omega', \right\}$$

with the travel time conditions

$$t' = -c_0 \alpha y_1 + |x - y|$$
; $t = -c_0 \alpha y_1 + |x - y| + |x - z|$.

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• In this case,
$$K_{F^*F} \in I^{2,0}(\Delta, \Lambda_{\pi_R(\Sigma)}) \Rightarrow K_{F^*F} \in I^2(\Delta \setminus \Lambda_{\pi_R(\Sigma)}), \ K_{F^*F} \in I^2(\Lambda_{\pi_R(\Sigma)} \setminus \Delta)$$

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Reduction of Artifacts

- The previous slide indicates that artifacts will flow out from Σ.
- However, we can preprocess the data and reduce the strength of the artifacts as follows.
- We construct a zero order ΨDO, namely Q such that it's principal symbol σ_Q vanishes to order s ≥ 1 on π_L(Σ).
- It then follows from the $I^{p,l}$ calculus (Greenleaf, Uhlmann, Marhuenda) that

$$F^*Qd = F^*QFv \in I^{2q-s+(k-1)/2,s-(k-1)/2}(\cdot, \cdot)$$
.

• In case 4, $K_{F^*QF} \in I^{2-s,2}(\Delta, \Lambda_{\pi_R(\Sigma)}) \Rightarrow K_{F^*QF} \in I^2(\Delta \setminus \Lambda_{\pi_R(\Sigma)}), K_{F^*QF} \in I^{2-s}(\Lambda_{\pi_R(\Sigma)} \setminus \Delta)$

• For example in case 3,
$$Q=(\partial_{y_2}^2+(\partial_{y_1}+\partial_{y_3})^2)(-\Delta)^{-1}$$

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• We have shown that it is possible to image moving objects and that unless a high-demensional data set is used (case 1), or else a filtered backprojection is employed, strong artifacts can occur in the back projected image.

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• A more practical consideration / implementation of this method would be useful.