

A Remark on Nonlinear Dirac Equations

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Abstract

For a n -dimensional spin manifold M with a fixed spin structure and a spinor bundle ΣM , we prove an ϵ -regularity theorem for weak solutions to the nonlinear Dirac equation of cubic nonlinearity. This, in particular, answers a regularity question raised by Chen-Jost-Wang [5] when $n = 2$.

1 Introduction

Linear Dirac type equations, including the Cauchy-Riemann equation in dimension two, are the most fundamental first order system of elliptic equations. During the course to study *Dirac-harmonic maps with curvature term* from a Riemann surface into a Riemannian manifold, Chen-Jost-Wang [4, 5] introduced the nonlinear Dirac equation with cubic nonlinearity:

$$\not{D}\psi^i = \sum_{j,k,l=1}^N H_{jkl}^i \langle \psi^j, \psi^k \rangle \psi^l, \quad 1 \leq i \leq N. \quad (1)$$

In dimension two, an interesting feature of this nonlinear Dirac equation is that it is conformally invariant and has critical nonlinearity, where the classical methods fail to apply. Thus it is an interesting question to study the regularity of weak solutions of (2). The aim of this short note is to provide an elementary proof of a general regularity criterion for (1).

In order to describe the results, we briefly review some background materials on spin manifolds. The interested readers can consult with Lawson-Michelsohn [6], Chen-Jost-Li-Wang [2, 3] for more details. For $n \geq 2$, let (M, g) be a spin manifold with a given spin structure and an associated spinor bundle Σ . Let $\langle \cdot, \cdot \rangle$ be a Hermitian metric on Σ and ∇ be the Levi-Civita connection on Σ compatible with both $\langle \cdot, \cdot \rangle$ and g . The Dirac operator on M is defined by $\not{D} = e_\alpha \circ \nabla_{e_\alpha}$, where $\{e_\alpha\}_{\alpha=1}^n$ is a local orthonormal frame on M , and $\circ : TM \otimes_{\mathbb{C}} \Sigma \rightarrow \Sigma$ is the Clifford multiplication.

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Now let's write (2) into the form

$$\not\partial\psi = H_{jkl}\langle\psi^j, \psi^k\rangle\psi^l, \quad (2)$$

where $\psi = (\psi^1, \dots, \psi^N) \in (\Gamma\Sigma)^N$, $N \geq 1$, $H_{jkl} = (H_{jkl}^1, \dots, H_{jkl}^N) \in C^\infty(M, \mathbb{R}^N)$.

We refer the readers to [5] §1, where the authors discussed two interesting examples in which (2) arises naturally. The first example is the Dirac-harmonic map (ϕ, ψ) associated with the Dirac-harmonic energy functional with curvature term, a nonlinear σ -model in the superstring theory, in which the nonlinear Dirac equation for ψ reduces to (2) when ϕ is a constant map. The second example is the Weierstrass representation formula for minimal surfaces X immersed in \mathbb{R}^3 by holomorphic 1-forms and meromorphic functions, in which an equation of the form (2) appears naturally.

It turns out that the underlying function space for the equation (2) is $L^4(M)$. As pointed out by [5] that any weak solution ψ of (2) is smooth provide $\psi \in L^p(M)$ for some $p > 4$. In [5], the authors proved three interesting analytic properties of (2) for $n = 2$: (i) the gradient estimate for *smooth* solutions ψ of (2) under the smallness condition of L^4 -norm of ψ , (ii) the isolated singularity removable theorem, and (iii) the energy identity theorem for sequentially weak convergent smooth solutions of (2). At the end of §1 in [5], the authors asked

Conjecture 1.1 *For $n = 2$, any weak solution $\psi \in L^4(M)$ of (2) is smooth.*

In this short note, we answer Conjecture 1.1 affirmatively. In fact, we prove a general regularity theorem for weak solutions of (2) in any dimensions. The ideas is based on an application of the estimate of Reisz potentials between Morrey spaces, due to Adams [1]. Similar techniques have been employed in the proof of higher order regularity of Dirac-harmonic maps by Wang-Xu [7]. The proof turns out to be very elementary, and may be applicable to other similar problems.

Before stating our results, let's first recall the definition of weak solutions of (2).

Definition 1.2 A section $\psi \in L^4((\Gamma\Sigma)^N)$ is a weak solution of (2) if

$$\int_M \langle\psi, \not\partial\eta\rangle = \int_M H_{jkl} \langle\psi^j, \psi^k\rangle \langle\psi^l, \eta\rangle \quad (3)$$

holds for any smooth section $\eta \in C^\infty((\Gamma\Sigma)^N)$.

Denote by $i_M > 0$ the injectivity radius of M . For $0 < r < i_M$ and $x \in M$, denote by $B_r(x)$ the geodesic ball in M with center x and radius r . Now we state our theorems.

Theorem 1.3 *For any $n \geq 2$, there exists $\epsilon_0 > 0$ depending on n such that if $\psi \in L^4((\Gamma\Sigma)^N)$ is a weak solution of the Dirac equation (2) and satisfies, for some $x_0 \in M$ and $0 < r_0 \leq \frac{1}{2}i_M$,*

$$\sup_{x \in B_{r_0}(x_0), 0 < r \leq r_0} \left\{ \frac{1}{r^{n-2}} \int_{B_r(x)} |\psi|^4 \right\} \leq \epsilon_0^4, \quad (4)$$

then $\psi \in C^\infty(B_{\frac{r_0}{2}}(x_0))$.

Note that by Hölder inequality, we have for $n \geq 2$,

$$\frac{1}{r^{n-2}} \int_{B_r(x)} |\psi|^4 \leq \left(\int_{B_r(x)} |\psi|^{2n} \right)^{\frac{2}{n}}.$$

Thus, as an immediate consequence of Theorem 1.3, we obtain

Corollary 1.4 *For $n \geq 2$, if $\psi \in L^{2n}((\Gamma\Sigma)^N)$ is a weak solution of the Dirac equation (2), then $\psi \in C^\infty((\Gamma\Sigma)^N)$.*

It is clear that when $n = 2$, Corollary 1.4 implies Conjecture 1.1.

2 Proof of Theorem 1.3

This section is devoted to the proof of Theorem 1.3. Since the regularity is a local property, we assume, for simplicity of presentation, that for $x_0 \in M$, the geodesic ball $B_{i_M}(x_0) \subset M$ with the metric g is identified by (B_2, g_0) . Here $B_2 \subset \mathbb{R}^n$ is the ball with center 0 and radius 2, and g_0 is the Euclidean metric on \mathbb{R}^n . We also assume that $\Sigma|_{B_2} = B_2 \times \mathbb{C}^L$, with $L = \text{rank}_{\mathbb{C}} \Sigma$.

Let's also recall the definition of Morrey spaces.

Definition 2.1 For $1 \leq p \leq n$, $0 < \lambda \leq n$, and a domain $U \subseteq \mathbb{R}^n$, the Morrey space $M^{p,\lambda}(U)$ is defined by

$$M^{p,\lambda}(U) := \left\{ f \in L^p_{\text{loc}}(U) : \|f\|_{M^{p,\lambda}(U)} < +\infty \right\},$$

where

$$\|f\|_{M^{p,\lambda}(U)}^p = \sup \left\{ r^{\lambda-n} \int_{B_r} |f|^p : B_r \subseteq U \right\}.$$

It is easy to see that for $1 \leq p \leq n$, $M^{p,\lambda}(U) \subset L^p(U)$, $M^{p,n}(U) = L^p(U)$ and $M^{p,p}(U)$ behaves like $L^n(U)$ from the view of scalings.

It is clear that the condition (4) in Theorem 1.3 is equivalent to

$$\|\psi\|_{M^{4,2}(B_{r_0}(x_0))} \leq \epsilon_0.$$

Thus Theorem 1.3 follows from the following Lemma.

Lemma 2.2 *For any $4 < p < +\infty$ and $n \geq 2$, there exists $\epsilon_0 > 0$ depending only on p and n such that if $\psi \in M^{4,2}(B_1)$ is a weak solution of (2) and*

$$\|\psi\|_{M^{4,2}(B_1)} \leq \epsilon_0,$$

then $\psi \in L^p(B_{\frac{1}{16}}, \mathbb{C}^{NL})$. Furthermore, $\psi \in C^\infty(B_{\frac{1}{16}}, \mathbb{C}^{NL})$ and the estimate

$$\left\| \nabla^l \psi \right\|_{C^0(B_{\frac{1}{16}})} \leq C(\epsilon_0, l), \quad \forall l \geq 1 \tag{5}$$

holds

Proof. Applying ∂ to (2), we have, for $1 \leq i \leq N$,

$$\partial^2 \psi^i = \partial \left(H_{jkl}^i \langle \psi^j, \psi^k \rangle \psi^l \right) \quad (6)$$

in the sense of distributions. By Lichnerowitz's formula (cf. [6]), we have

$$-\Delta \psi^i = \partial^2 \psi^i.$$

Hence we have

$$-\Delta \psi^i = \partial \left(H_{jkl} \langle \psi^j, \psi^k \rangle \psi^l \right) \quad (7)$$

in the sense of distributions.

For $m = 1, 2$, let $\eta_m \in C_0^\infty(B_1)$ be such that $0 \leq \eta_m \leq 1$, $\eta_m \equiv 1$ on $B_{2^{1-2m}}$. For $1 \leq i \leq N$, define $f_m^i : \mathbb{R}^n \rightarrow \mathbb{C}^L$ by letting

$$f_m^i(x) = \int_{\mathbb{R}^n} \frac{\partial G(x, y)}{\partial y_\alpha} \frac{\partial}{\partial y_\alpha} \circ \left(\eta_m^3 H_{jkl} \langle \psi^j, \psi^k \rangle \psi^l \right) (y) dy, \quad (8)$$

where $G(x, y)$ is the fundamental solution of Δ on \mathbb{R}^n . For $1 \leq i \leq N$, define $g_m^i : B_1 \rightarrow \mathbb{C}$ by letting

$$\psi^i = f_m^i + g_m^i. \quad (9)$$

Direct calculations imply that for $m = 1, 2$ and $1 \leq i \leq N$,

$$\begin{aligned} -\Delta f_m^i &= \partial \left(\eta_m^3 H_{jkl} \langle \psi^j, \psi^k \rangle \psi^l \right) \\ &= \partial \left(H_{jkl} \langle \psi^j, \psi^k \rangle \psi^l \right) \quad \text{in } B_{2^{1-2m}}. \end{aligned} \quad (10)$$

This and (7) imply

$$\Delta g_m^i = 0 \quad \text{in } B_{2^{1-2m}}. \quad (11)$$

It follows from (8) that for $m = 1, 2$ and $1 \leq i \leq N$,

$$|f_m^i|(x) \leq C \int_{\mathbb{R}^n} |x - y|^{1-n} (\eta_m(y) |\psi(y)|)^3 dy = CI_1(\eta_m^3 |\psi|^3)(x), \quad (12)$$

where

$$I_1(f)(x) = \int_{\mathbb{R}^n} |x - y|^{1-n} |f(y)| dy, \quad f : \mathbb{R}^n \rightarrow \mathbb{R},$$

is the Riesz potential of order one. Let's recall Adams' inequality on Morrey spaces (cf. [1]):

$$\|I_1(f)\|_{M^{\frac{\lambda q}{\lambda - q}, \lambda}(\mathbb{R}^n)} \leq C \|f\|_{M^{q, \lambda}(\mathbb{R}^n)}, \quad \forall 1 \leq q < \lambda \leq n. \quad (13)$$

Step 1 ($m = 1$). Since $(\eta_1 |\psi|)^3 \in M^{\frac{4}{3}, 2}(\mathbb{R}^n)$, (13) implies that for $1 \leq i \leq N$,

$$\begin{aligned} \|f_1^i\|_{M^{4, 2}(\mathbb{R}^n)} &\leq C \|\eta_1^3 |\psi|^3\|_{M^{\frac{4}{3}, 2}(\mathbb{R}^n)} = C \|\eta_1 |\psi|\|_{M^{4, 2}(\mathbb{R}^n)}^3 \\ &\leq C \|\psi\|_{M^{4, 2}(B_1)}^3 \leq C \epsilon_0^2 \|\psi\|_{M^{4, 2}(B_1)}. \end{aligned} \quad (14)$$

On the other hand, by the standard estimate for harmonic functions, we have that for any $\theta \in (0, \frac{1}{4})$ and $x_0 \in B_{\frac{1}{4}}$

$$\|g_1^i\|_{M^{4,2}(B_\theta(x_0))} \leq C\theta^{\frac{1}{2}}\|g_1^i\|_{M^{4,2}(B_{\frac{1}{2}})}, \quad \forall 1 \leq i \leq N. \quad (15)$$

Putting (14) and (15) into (9) yields that for $1 \leq i \leq N$,

$$\begin{aligned} \|\psi^i\|_{M^{4,2}(B_\theta(x_0))} &\leq C\theta^{\frac{1}{2}}\|g_1^i\|_{M^{4,2}(B_{\frac{1}{2}})} + C\epsilon_0^2\|\psi\|_{M^{4,2}(B_1)} \\ &\leq C\theta^{\frac{1}{2}}\left[\|\psi^i\|_{M^{4,2}(B_{\frac{1}{2}})} + \|f_1^i\|_{M^{4,2}(B_{\frac{1}{2}})}\right] + C\epsilon_0^2\|\psi\|_{M^{4,2}(B_1)} \\ &\leq C\left(\epsilon_0^2 + \theta^{\frac{1}{2}}\right)\|\psi\|_{M^{4,2}(B_1)}. \end{aligned} \quad (16)$$

This clearly implies that for any $\theta \in (0, \frac{1}{4})$ and $x_0 \in B_{\frac{1}{4}}$,

$$\|\psi\|_{M^{4,2}(B_\theta(x_0))} \leq C\left(\epsilon_0^2 + \theta^{\frac{1}{2}}\right)\|\psi\|_{M^{4,2}(B_1)}. \quad (17)$$

For any $\alpha \in (0, \frac{1}{3})$, first choose $\theta \in (0, \frac{1}{4})$ be such that $C\theta^{\frac{1}{2}} \leq \theta^{\frac{\alpha}{2}}$ and then choose $\epsilon_0 > 0$ be such that $C\epsilon_0^2 \leq \theta^{\frac{\alpha}{2}}$. Then we have

$$\|\psi\|_{M^{4,2}(B_\theta(x_0))} \leq \theta^{\frac{\alpha}{2}}\|\psi\|_{M^{4,2}(B_1)}, \quad \forall x_0 \in B_{\frac{1}{4}}. \quad (18)$$

Iteration of (18) yields

$$\|\psi\|_{M^{4,2}(B_r(x_0))} \leq Cr^{\frac{\alpha}{2}}\|\psi\|_{M^{4,2}(B_1)}, \quad \forall x_0 \in B_{\frac{1}{4}} \text{ and } 0 \leq r < \frac{1}{4}. \quad (19)$$

In particular, we have for any $0 < \alpha < \frac{1}{3}$,

$$r^{2(1-\alpha)-n} \int_{B_r(x_0)} |\psi|^4 \leq C \int_{B_1} |\psi|^4, \quad \forall x_0 \in B_{\frac{1}{4}} \text{ and } 0 < r < \frac{1}{4}. \quad (20)$$

Thus $\psi \in M^{4,2(1-\alpha)}(B_{\frac{1}{4}})$ for any $\alpha \in (0, 1)$.

Step 2 ($m = 2$). We want to repeat the above argument to show that $\psi \in M^{\frac{4-4\alpha}{1-3\alpha}, 2(1-\alpha)}(B_{\frac{1}{16}})$. In fact, since $(\eta_2|\psi|^3) \in M^{\frac{4}{3}, 2(1-\alpha)}(\mathbb{R}^n)$, (13) implies that $f_2^i \in M^{\frac{4(1-\alpha)}{1-3\alpha}, 2(1-\alpha)}(\mathbb{R}^n)$, and

$$\begin{aligned} \|f_2^i\|_{M^{\frac{4(1-\alpha)}{1-3\alpha}, 2(1-\alpha)}(B_{\frac{1}{8}})} &\leq \|f_2^i\|_{M^{\frac{4(1-\alpha)}{1-3\alpha}, 2(1-\alpha)}(\mathbb{R}^n)} \\ &\leq C\|\eta_2^3|\psi|^3\|_{M^{\frac{4}{3}, 2(1-\alpha)}(\mathbb{R}^n)} \\ &\leq C\|\psi\|_{M^{4,2(1-\alpha)}(B_{\frac{1}{4}})}. \end{aligned} \quad (21)$$

On the other hand, since g_2^i is a harmonic function on $B_{\frac{1}{8}}$, we have, by (21),

$$\begin{aligned}
& \|g_2^i\|_{M^{\frac{4(1-\alpha)}{1-3\alpha}, 2(1-\alpha)}(B_{\frac{1}{16}})} \\
& \leq C \|g_2^i\|_{M^{\frac{4(1-\alpha)}{1-3\alpha}, 2(1-\alpha)}(B_{\frac{1}{8}})} \\
& \leq C \left[\|f_2^i\|_{M^{\frac{4(1-\alpha)}{1-3\alpha}, 2(1-\alpha)}(B_{\frac{1}{8}})} + \|\psi^i\|_{M^{\frac{4(1-\alpha)}{1-3\alpha}, 2(1-\alpha)}(B_{\frac{1}{8}})} \right] \\
& \leq C \|\psi\|_{M^{\frac{4(1-\alpha)}{1-3\alpha}, 2(1-\alpha)}(B_{\frac{1}{8}})}. \tag{22}
\end{aligned}$$

Putting (21) and (22) into (9) yields that $\psi \in M^{\frac{4(1-\alpha)}{1-3\alpha}, 2(1-\alpha)}(B_{\frac{1}{16}})$ and

$$\|\psi\|_{M^{\frac{4(1-\alpha)}{1-3\alpha}, 2(1-\alpha)}(B_{\frac{1}{16}})} \leq C \|\psi\|_{M^{4,2}(B_{\frac{1}{4}})} \leq C \|\psi\|_{M^{4,2}(B_1)}. \tag{23}$$

Since

$$\lim_{\alpha \uparrow \frac{1}{3}} \frac{4(1-\alpha)}{1-3\alpha} = +\infty \text{ and } M^{\frac{4(1-\alpha)}{1-3\alpha}, 2(1-\alpha)}(B_{\frac{1}{16}}) \subseteq L^{\frac{4(1-\alpha)}{1-3\alpha}}(B_{\frac{1}{16}}),$$

it follows from that $\psi \in L^p(B_{\frac{1}{16}})$ for any $p > 4$, and

$$\|\psi\|_{L^p(B_{\frac{1}{16}})} \leq C(n, p) \|\psi\|_{M^{4,2}(B_1)}. \tag{24}$$

Since $|\delta\psi| \leq C|\psi|^3$, $W^{1,p}$ -estimate implies that $\psi \in W_{\text{loc}}^{1,p}(B_{\frac{1}{16}}, \mathbb{C}^{NL})$ for any $p > 4$. Hence, by the Sobolev embedding theorem, $\psi \in C^\mu(B_{\frac{1}{16}}, \mathbb{C}^{NL})$ for any $\mu \in (0, 1)$. By the Schauder estimate, this yields $\psi \in C^{1,\mu}(B_{\frac{1}{16}}, \mathbb{C}^{NL})$. Hence, by the bootstrap argument, we conclude $\psi \in C^\infty(B_{\frac{1}{16}}, \mathbb{C}^{NL})$ and the estimate (5) holds. \square

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